

Today's outline - November 04, 2024





- Dumond diagrams & monochromators

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- Dumond diagrams & monochromators
- Photoelectric absorption

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- Dumond diagrams & monochromators
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- X-ray absorption spectroscopy

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Reading Assignment: Chapter 7.4

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- Photoelectric absorption
- X-ray absorption spectroscopy
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Reading Assignment: Chapter 7.4

Homework Assignment #06:

Chapter 6: 1,6,7,8,9

due Friday, November 15, 2024

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- Dumond diagrams & monochromators
- Photoelectric absorption
- X-ray absorption spectroscopy
- EXAFS theory

Reading Assignment: Chapter 7.4

Homework Assignment #06:

Chapter 6: 1,6,7,8,9

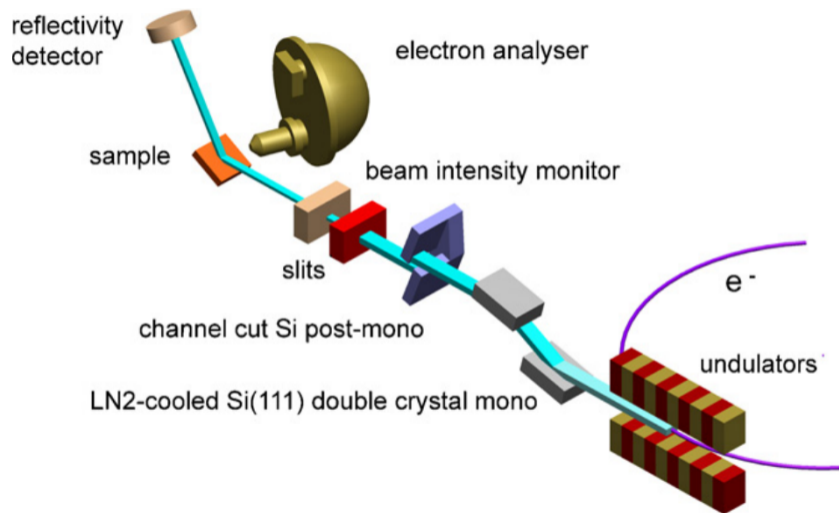
due Friday, November 15, 2024

Homework Assignment #07:

Chapter 7: 2,3,9,10,11

due Monday, November 25, 2024

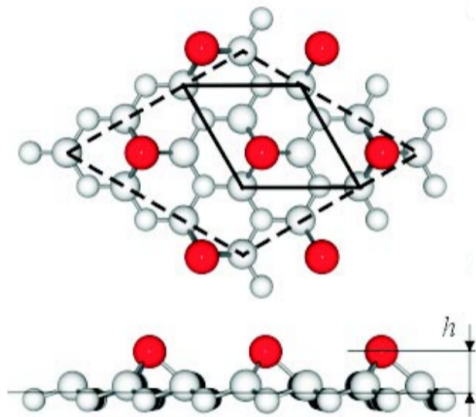
Beam line ID32 @ ESRF



Structure of Sn on Ge(111)



The low temperature 3×3 structure (dashed line) is well known but the room temperature $\sqrt{3} \times \sqrt{3}$ surface structure (solid line) is unresolved



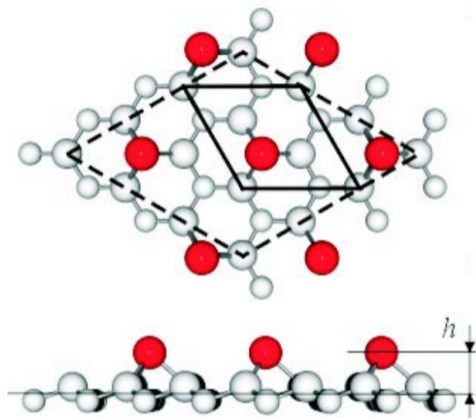
“Chemically resolved structure of the Sn/Ge(111) surface,” T.-L. Lee, S. Warren, B.C.C. Cowie, and J. Zengenhagen, *Phys. Rev. Lett.* **96**, 046103 (2006).

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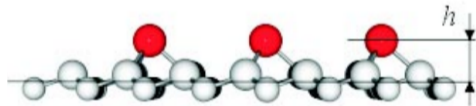
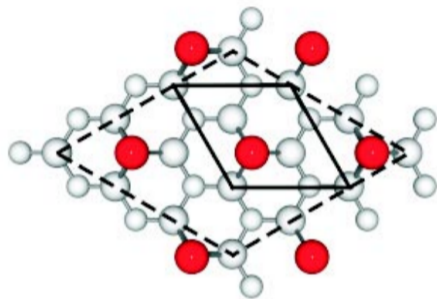
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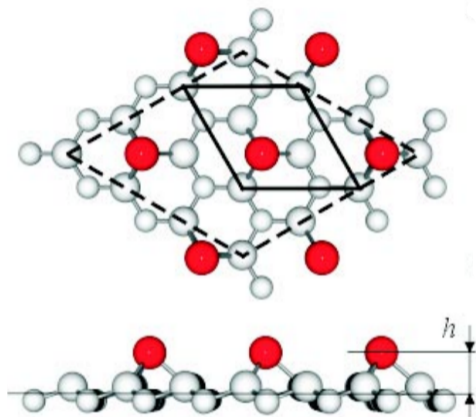


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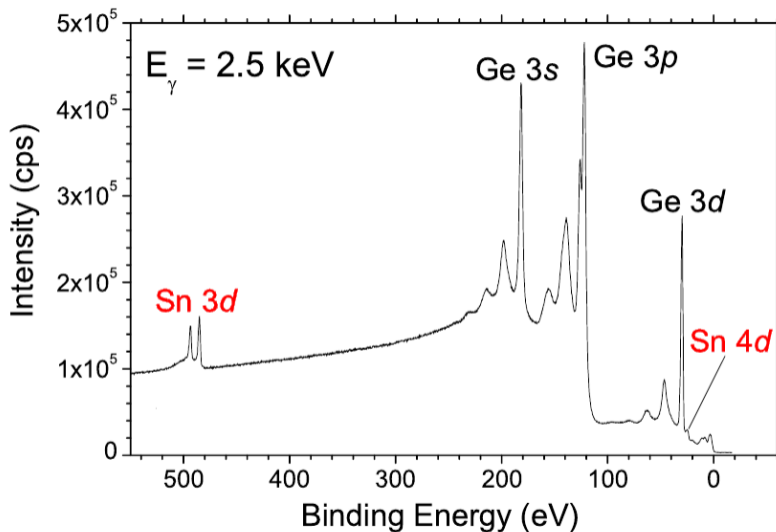
Below 0.2 ML, the well known 2×2 structure is measured as a reference

Above 0.2 ML, the $\sqrt{3} \times \sqrt{3}$ structure appears and then dominates



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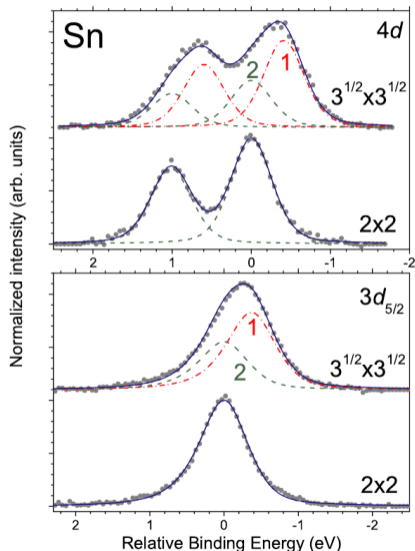
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With an incident energy of 2.5 keV, the 2×2 and $\sqrt{3} \times \sqrt{3}$ structures are measured in an off-Bragg condition

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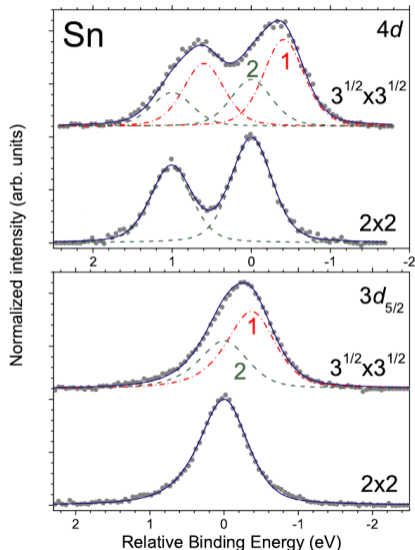


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The lines for both the Sn $3d_{5/2}$ and $4d$ peaks in the 2×2 phase are sharp, indicating a single chemical state

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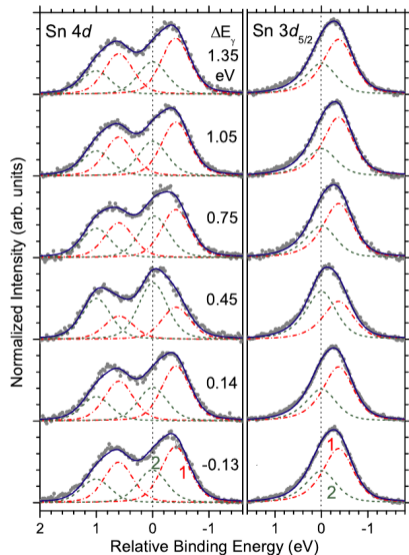
The $\sqrt{3} \times \sqrt{3}$ structure shows two distinct chemical shifts, with the majority component, l_1 , having a slightly lower binding energy than the minority component, l_2

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By varying the energy with a resolution of 500 meV, the standing wave is swept through the Sn layer



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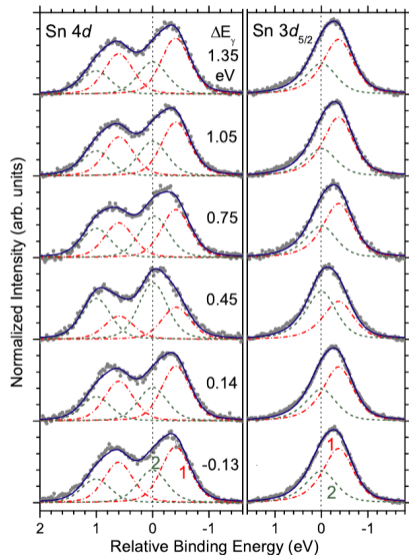
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As the energy is scanned around the center of the Ge(111) reflection, the fits using a mixture of Gaussian and Lorentzian line shapes show that the relative intensity, I_1/I_2 varies

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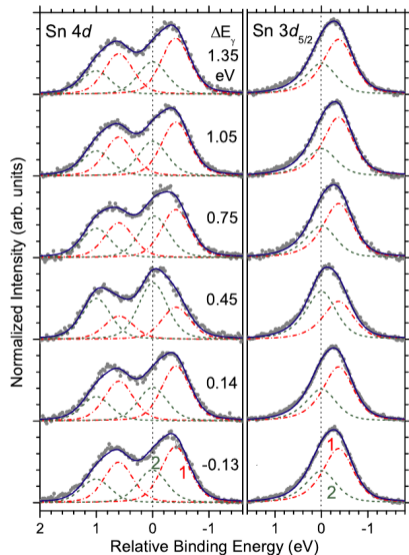


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At $\Delta E_\gamma = 0.45$ eV, the I_1/I_2 ratio almost completely inverts, showing that the two atom populations are at different heights above the surface

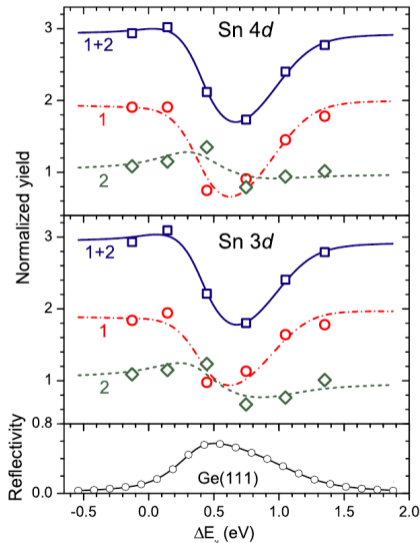
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The normalized peak intensities can be fitted to extract the relative positions of the two populations of atoms and their atomic ratio



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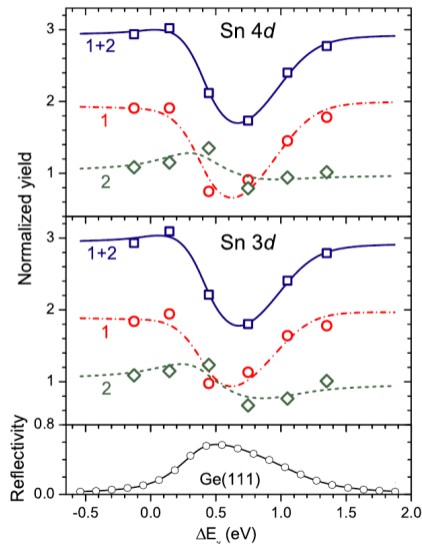
Structure of Sn on Ge(111)



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Population **1** is two times larger than population **2** and is located a height $\Delta h = 0.23 \text{ \AA}$ further from the Ge(111) surface

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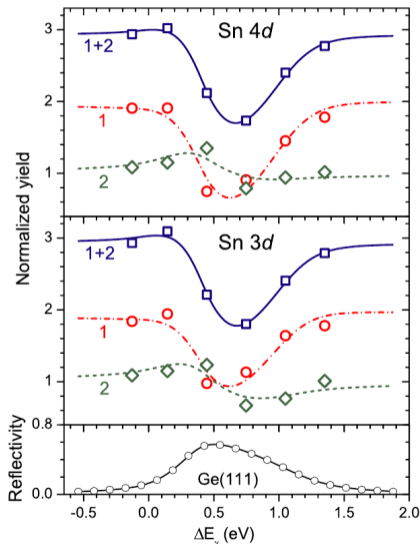


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Population **1** is two times larger than population **2** and is located a height $\Delta h = 0.23 \text{ \AA}$ further from the Ge(111) surface

Population **1** also has a lower binding energy, demonstrating that the binding energy is directly correlated to the height from the surface

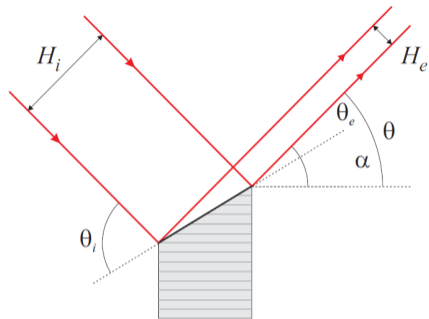
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Asymmetric geometry



When the diffracting planes are not precisely aligned with the surface of the crystal the asymmetry angle, α , is the important parameter

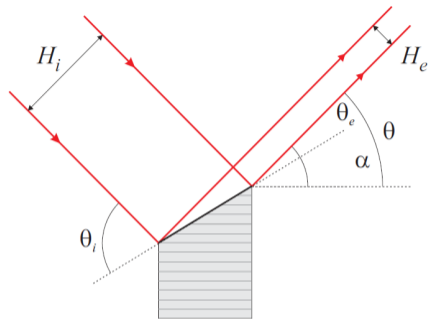


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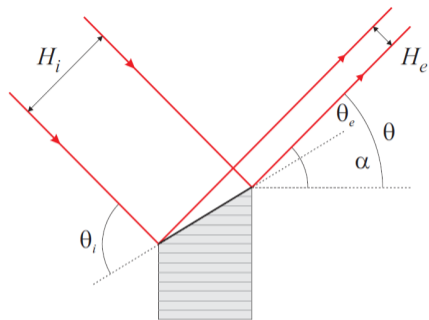
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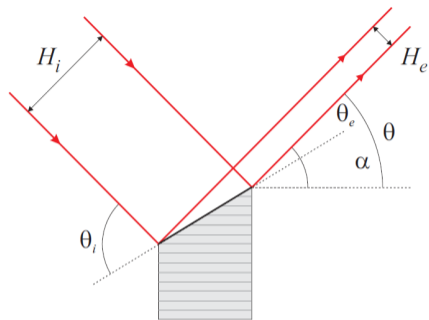


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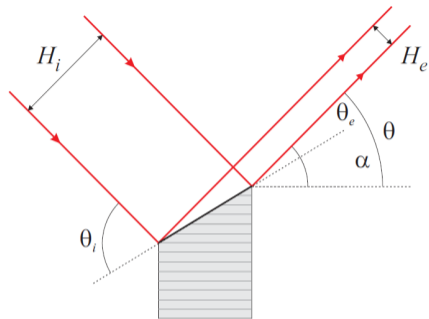


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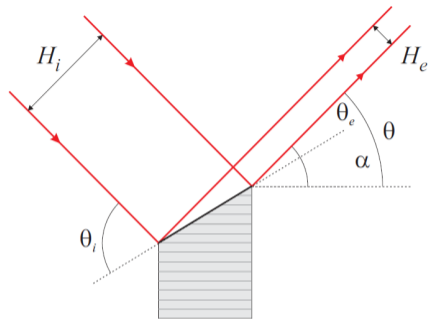


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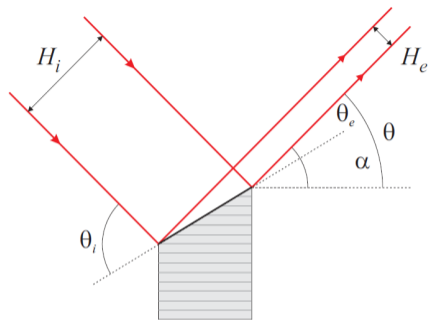
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according to Liouville's theorem, phase space is invariant so the divergence of the beam, $\delta\theta$, must also change



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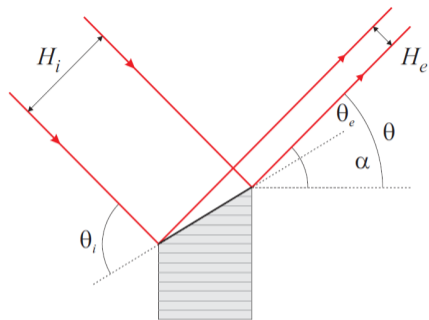
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$$\delta\theta_e = \sqrt{b}(\zeta_D \tan \theta)$$

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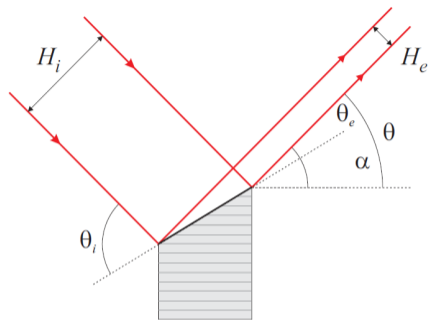
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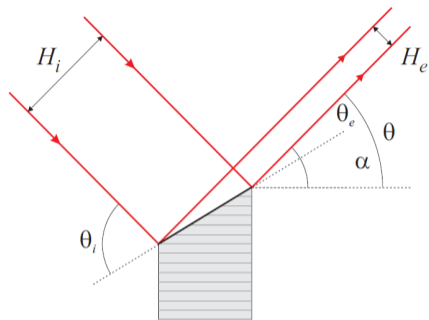
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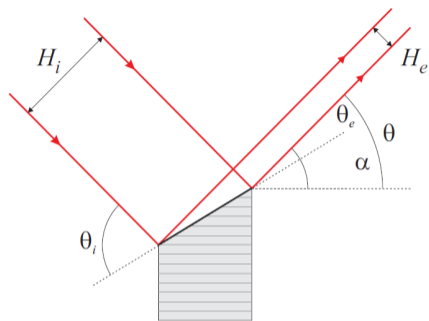
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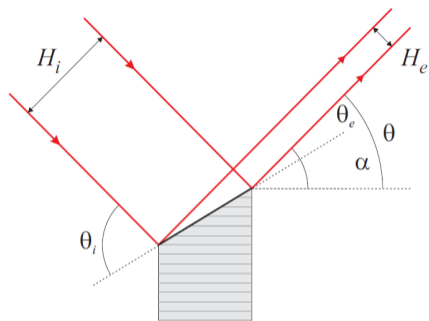
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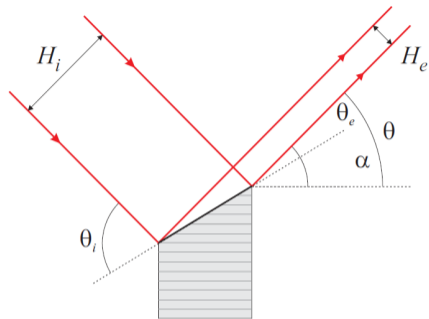
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Rocking curve measurements

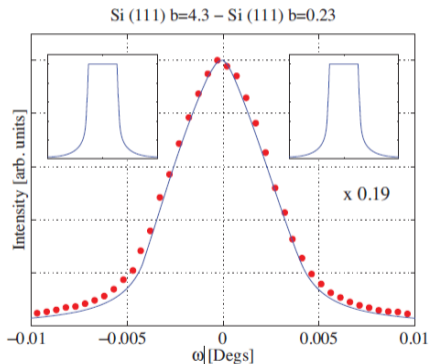


The measured “rocking” curve from a two crystal system is a convolution of the Darwin curves of both crystals.

Rocking curve measurements



The measured “rocking” curve from a two crystal system is a convolution of the Darwin curves of both crystals. When the two crystals have a matched asymmetry, we get a triangle.

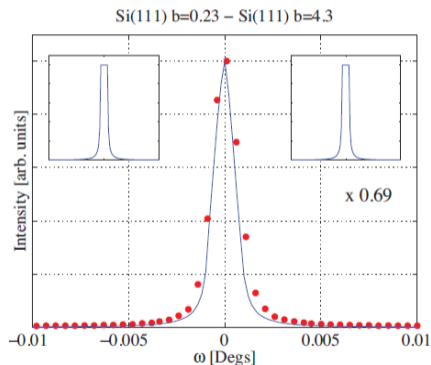
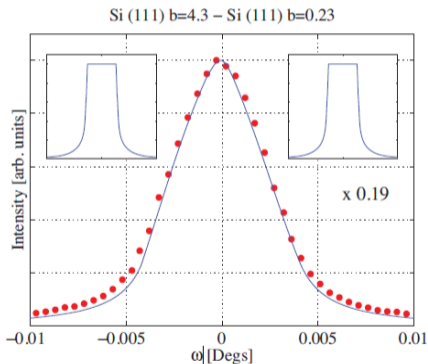


output divergence on left, input divergence on right

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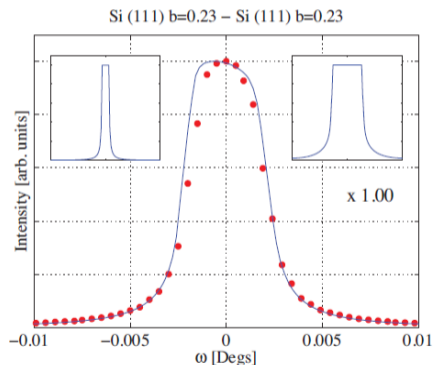
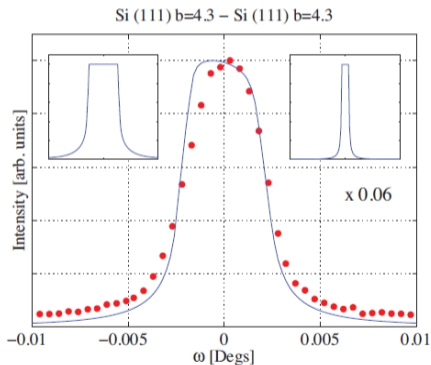


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Rocking curve measurements



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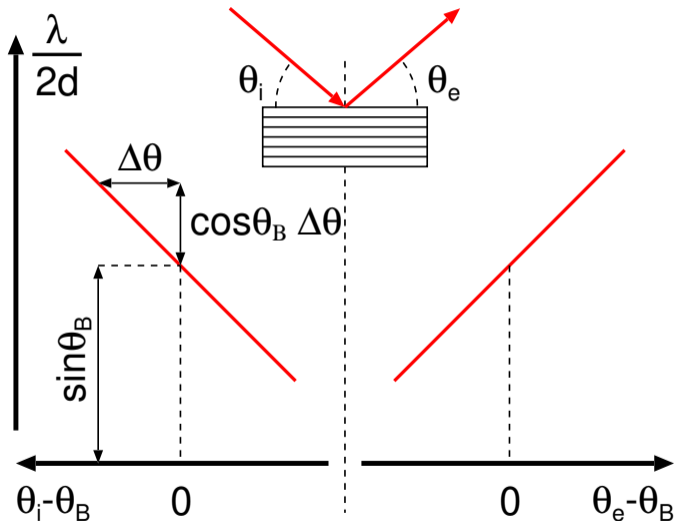


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Dumond diagram: no Darwin width



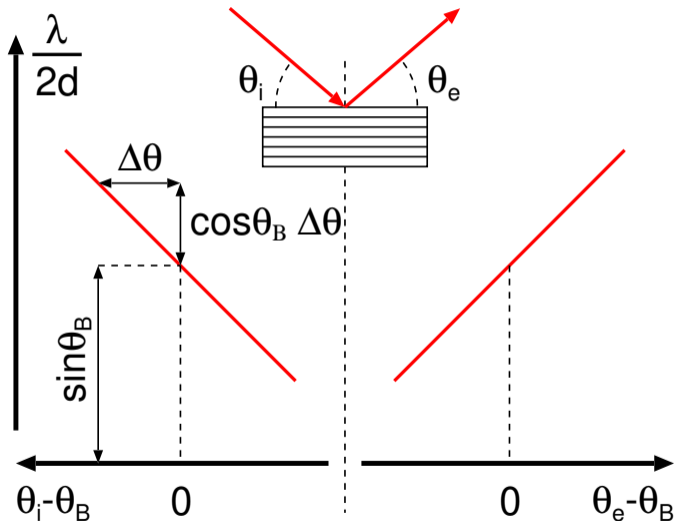
Transfer function of an optical element parametrized by angle and wavelength.



Dumond diagram: no Darwin width



Transfer function of an optical element parametrized by angle and wavelength. Here Darwin width is ignored.

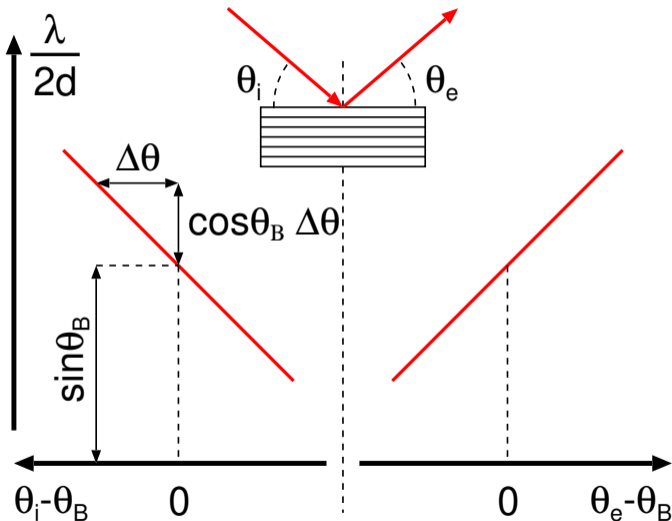


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for small angular deviations $\sin \theta$ is linear with a slope of $\cos \theta_B$



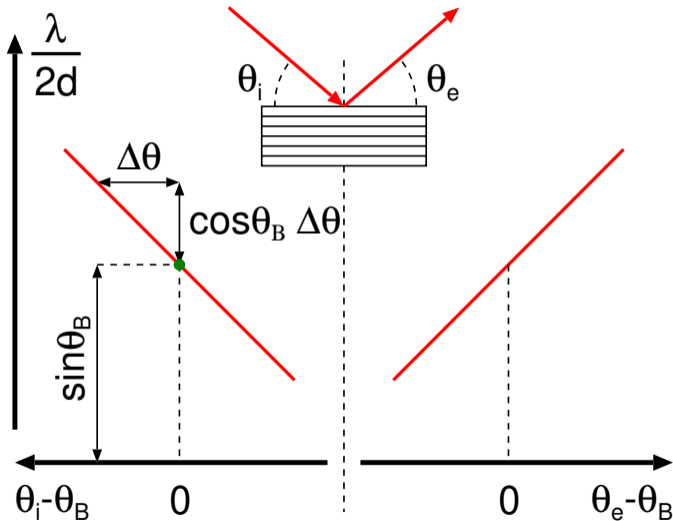
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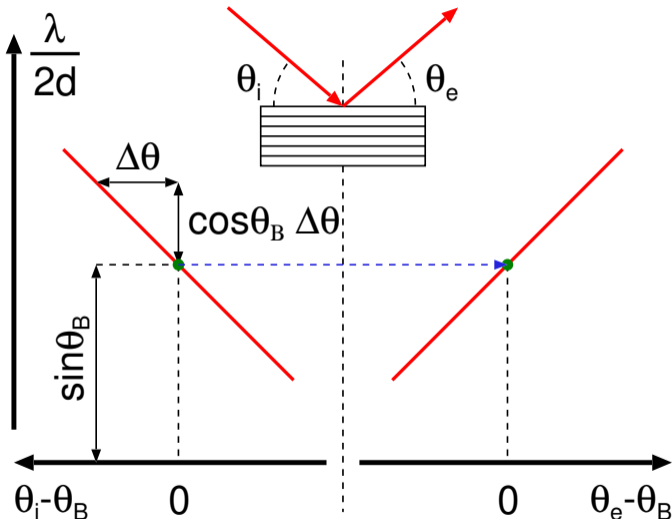


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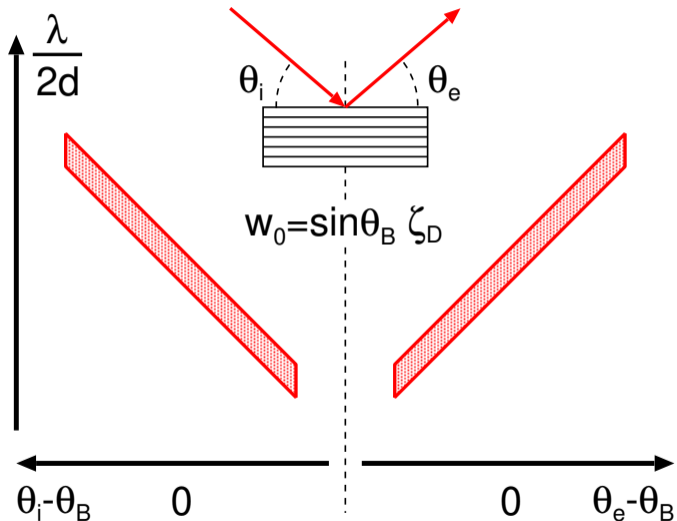
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a horizontal line transfers input to output beam characteristics

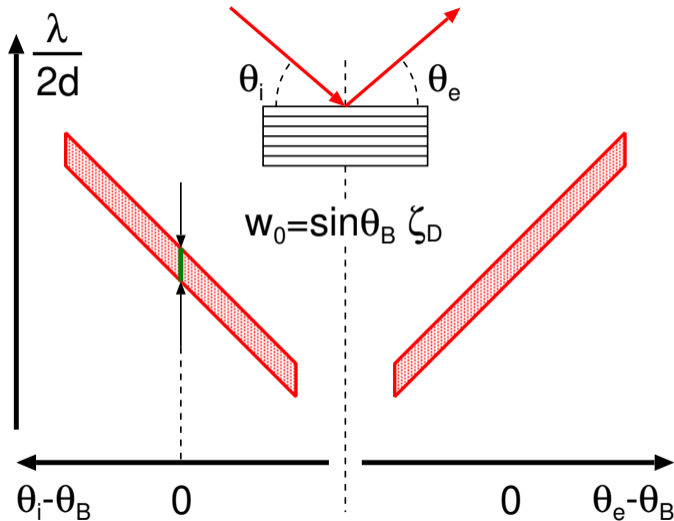


Dumond diagram: symmetric Bragg



If Darwin width is included, the Bragg condition is represented by a **box**

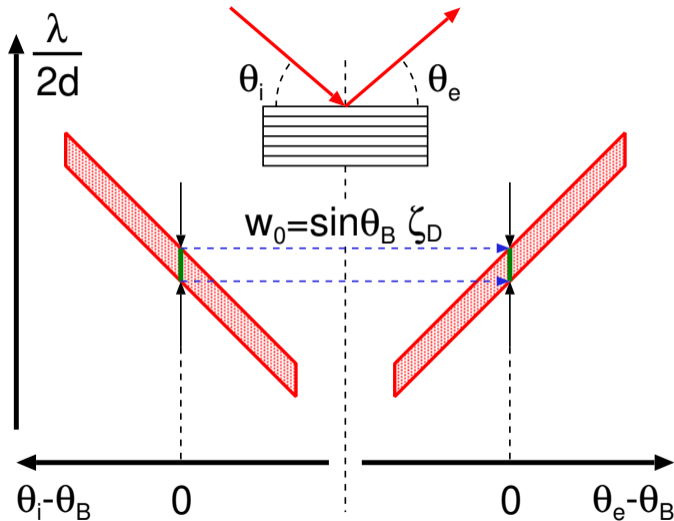
Dumond diagram: symmetric Bragg



If Darwin width is included, the Bragg condition is represented by a **box**

for a perfectly collimated (no angular divergence) input beam, a **bandwidth** of radiation is accepted by the crystal

Dumond diagram: symmetric Bragg

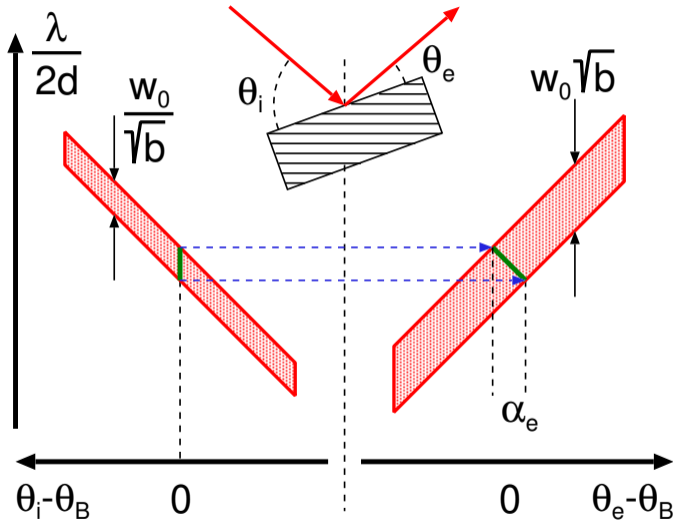


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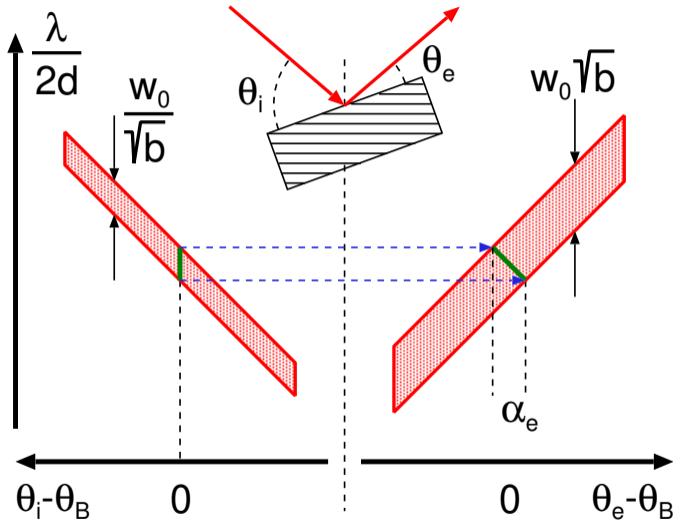
this input bandwidth is **transferred** to a similar output bandwidth which is also collimated

Dumond diagram: asymmetric Bragg



For an asymmetric crystal, the output beam is no longer collimated but acquires a divergence α_e

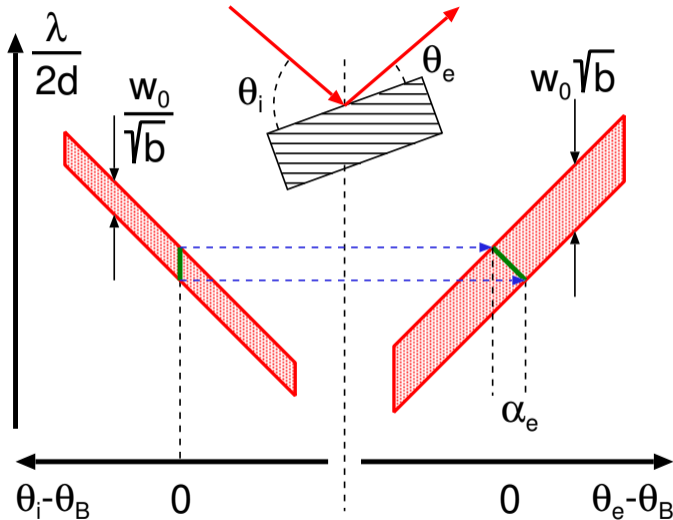
Dumond diagram: asymmetric Bragg



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a perfectly collimated input beam transfers to an output beam that has an angular divergence which depends on the asymmetry factor b

Dumond diagram: asymmetric Bragg

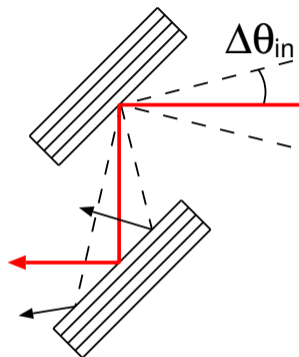


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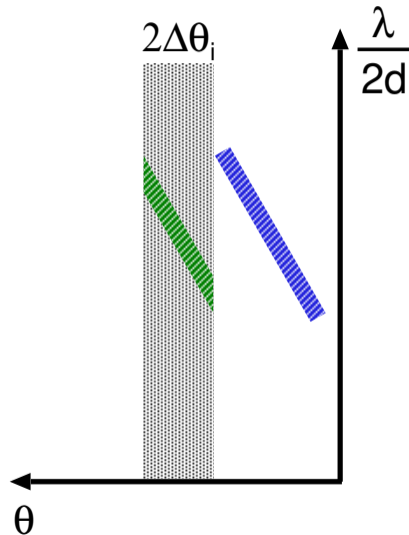
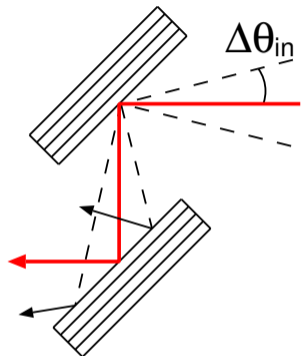
a perfectly collimated input beam transfers to an output beam that has an angular divergence which depends on the asymmetry factor b

this is in addition to a compression (in this case) of the beam height (Liouville's theorem!)

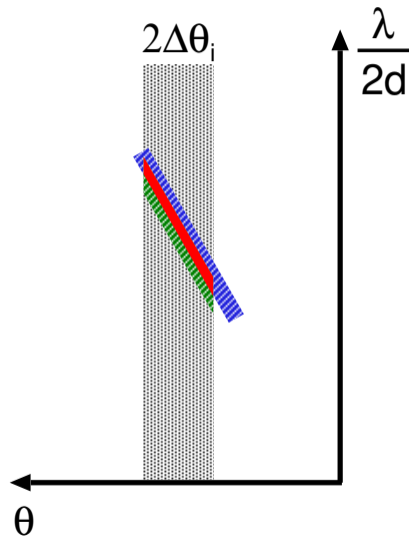
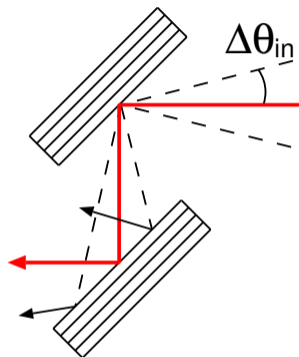
Double crystal monochromator: Non-dispersive



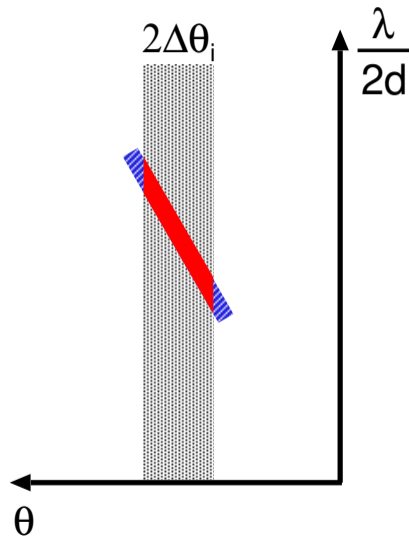
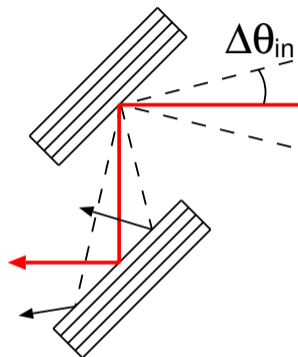
Double crystal monochromator: Non-dispersive



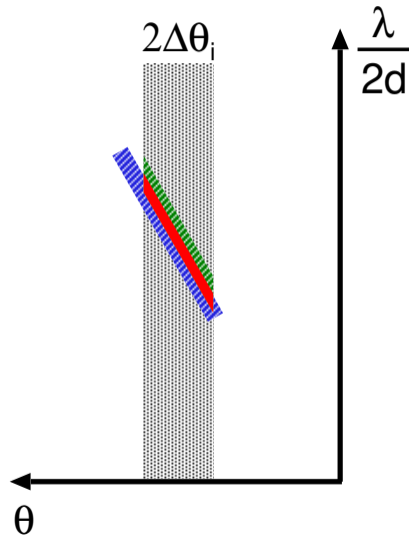
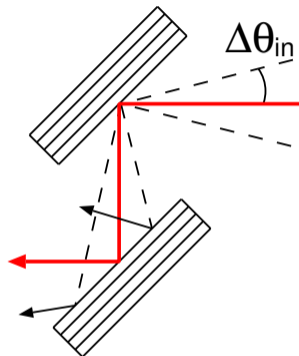
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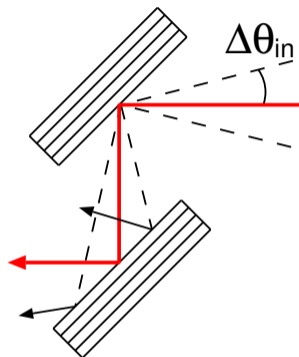
Double crystal monochromator: Non-dispersive



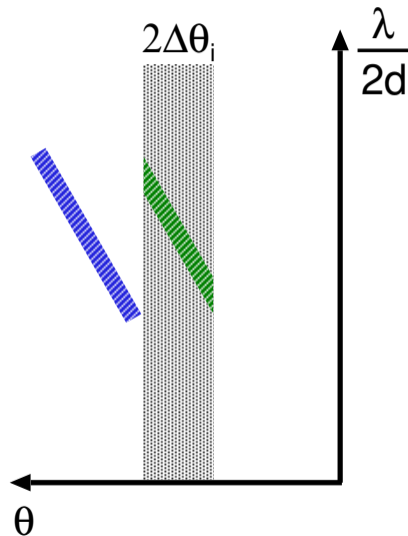
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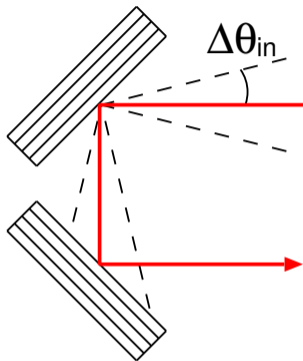
Double crystal monochromator: Non-dispersive



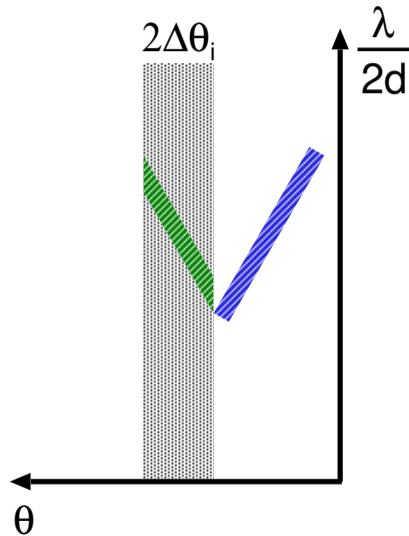
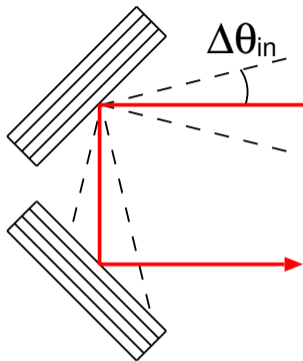
the transfer functions of the two crystals match and full bandwidth and divergence is preserved, giving a triangle intensity curve



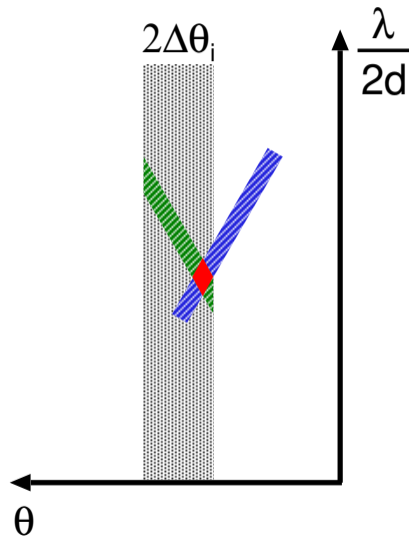
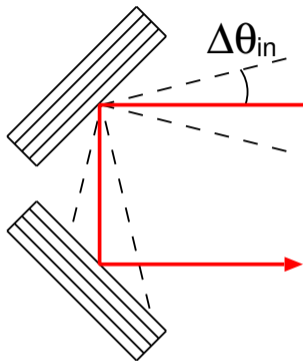
Double crystal monochromators: Dispersive



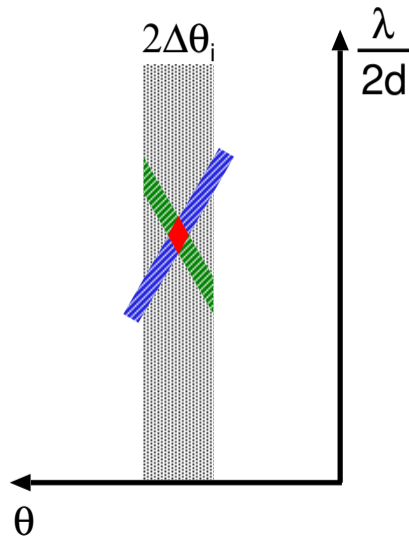
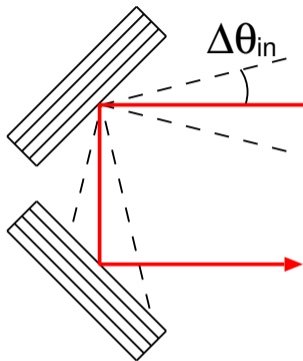
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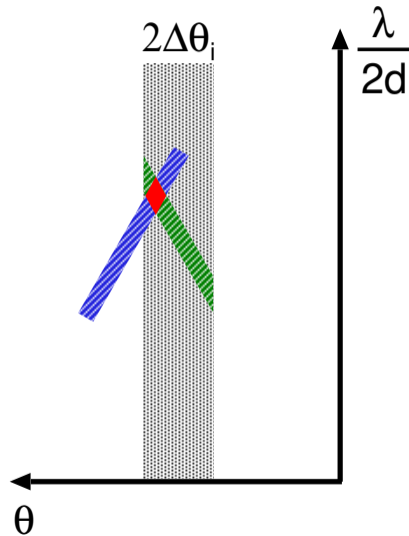
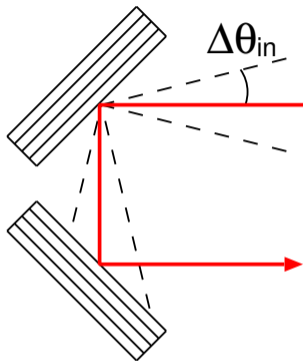
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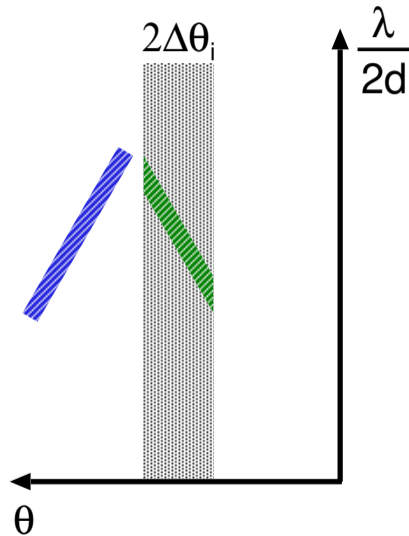
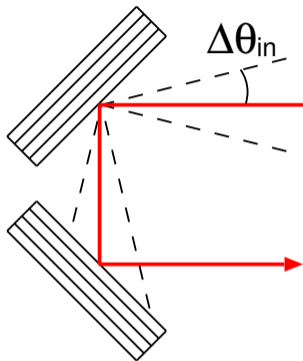
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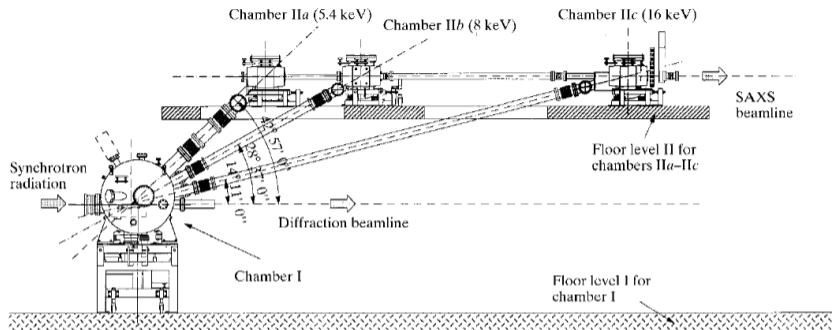


Double crystal monochromators: Dispersive



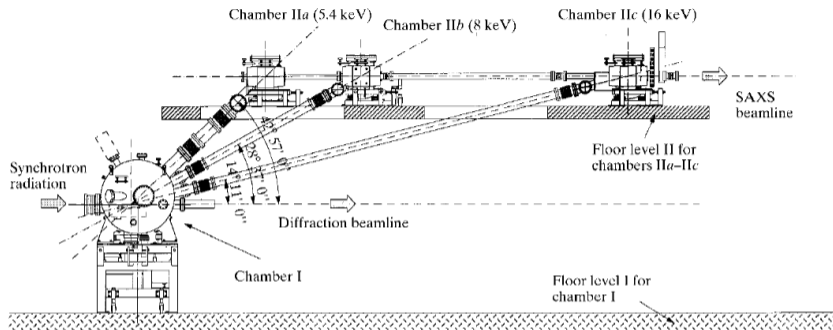
the transfer function matches only in a small energy band that varies with angle of the second crystal, mapping out the Darwin curve of the first crystal

Asymmetric monochromator at ELETTRA



"High-throughput asymmetric double-crystal monochromator of the SAXS beamline at ELETTRA," S. Bernstorff, H. Amentisch, and P. Laggner, *J. Synchrotron Rad.* **5**, 1215-1221 (1998).

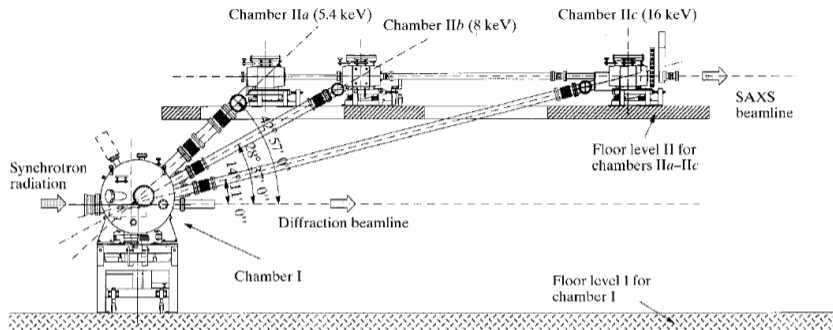
Asymmetric monochromator at ELETTRA



The SAXS beamline at ELETTRA has asymmetric cut crystals with 2° grazing incidence in order to spread the heat load

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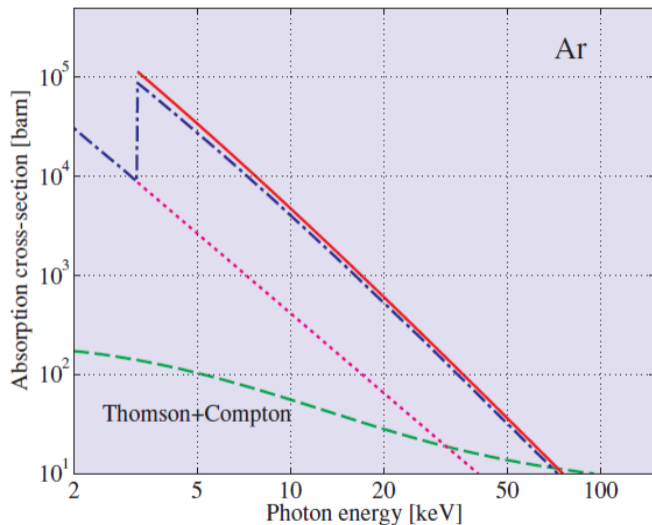


The SAXS beamline at ELETTRA has asymmetric cut crystals with 2° grazing incidence in order to spread the heat load

The three crystals are set for single energies of 5.6, 8.0, and 16 keV with a vertical displacement of 1.5 m and asymmetry parameter, b , of 0.053, 0.078, and 0.17, respectively

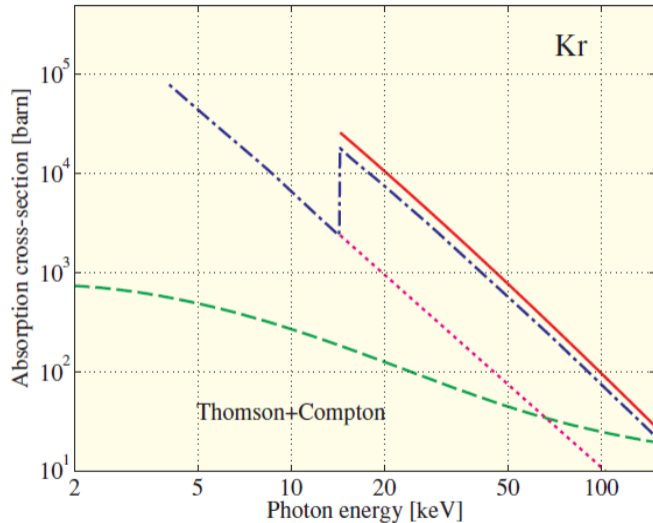
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Total cross section



The total cross-section for photon “absorption” includes elastic (or coherent) scattering, Compton (inelastic) scattering, and photoelectric absorption.

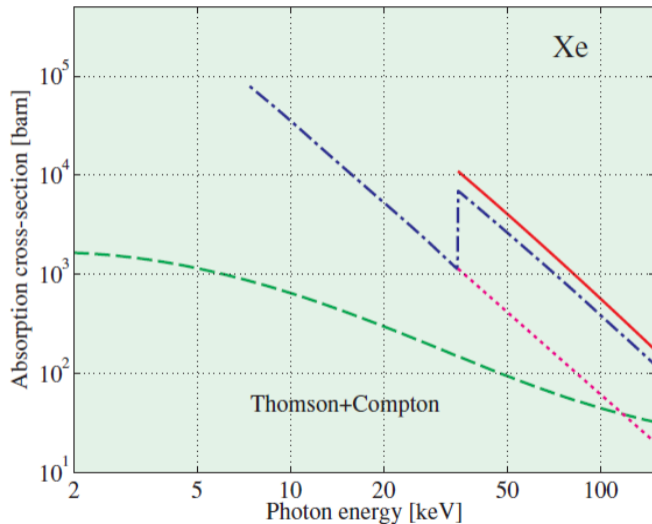
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Characteristic absorption jumps depend on the element

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Characteristic absorption jumps depend on the element

These quantities vary significantly over many decades but can easily be put on an equal footing.

Scaled absorption



$$T = \frac{I}{I_0} = e^{-\mu z}$$

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σ_a can be scaled for different elements
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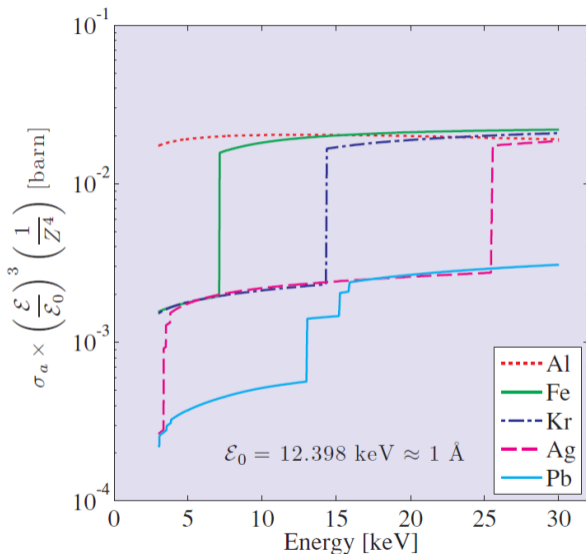
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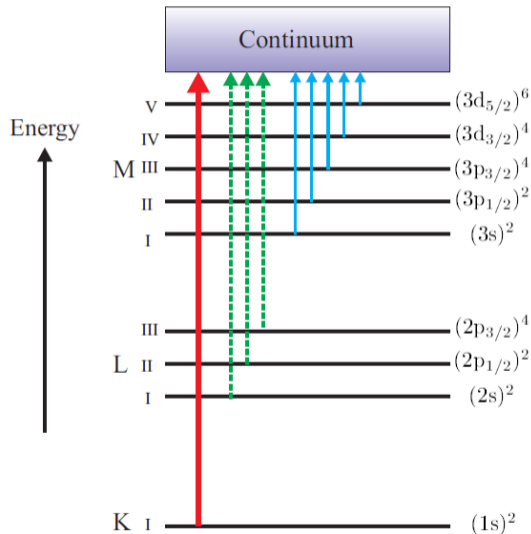
remarkably, all values lie on a common curve above the K edge and between the L and K edges and below the L edge



Absorption edge nomenclature



The states are labeled according to the principal, orbital angular momentum, and total angular momentum quantum numbers, n , l , and j , respectively

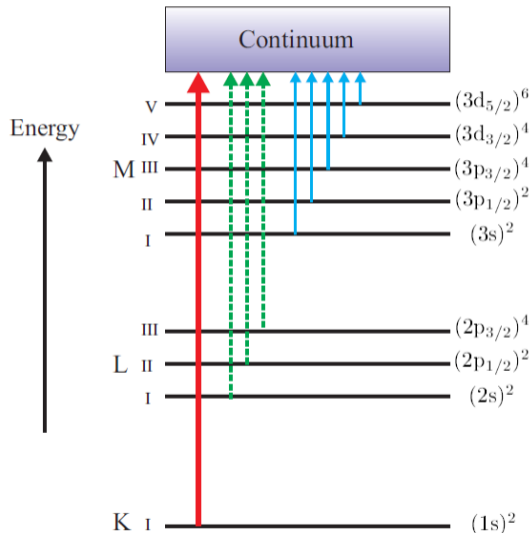


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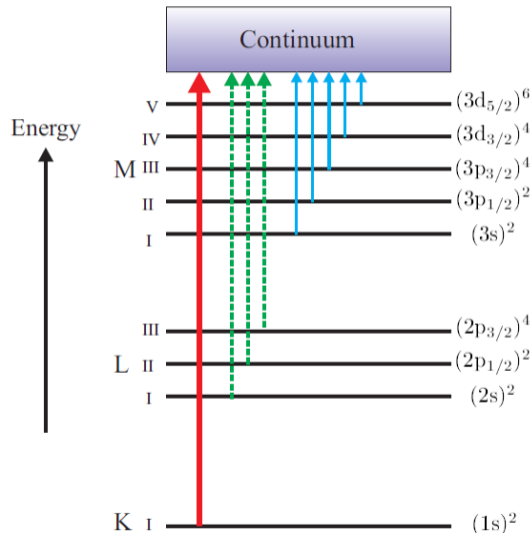
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Roman numerals increase from low to high values of l and j

