



• Writing a General User Proposal



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- Dynamical theory and the Darwin curve



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- Extinction and absorption



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- Perfect crystal integrated intensity



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Reading Assignment: Chapter 6.5; Chapter 7.1



1/31

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- Dynamical theory and the Darwin curve
- Extinction and absorption
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Reading Assignment: Chapter 6.5; Chapter 7.1

Homework Assignment #05: Chapter 5: 1,2,7,9,10 due Monday, October 28, 2024



- Writing a General User Proposal
- Dynamical theory and the Darwin curve
- Extinction and absorption
- Perfect crystal integrated intensity

Reading Assignment: Chapter 6.5; Chapter 7.1

Homework Assignment #05: Chapter 5: 1,2,7,9,10 due Monday, October 28, 2024 Homework Assignment #06: Chapter 6: 1,6,7,8,9 due Monday, November 11, 2024 Writing a General User Proposal



- 1. Log into the UPS site
- 2. Start an APS general user proposal
- 3. Add an Abstract
- 4. Choose a beam line
- 5. Answer the 6 important questions

A tutorial can be found on the course home page

 $http://csrri.iit.edu/{\sim}segre/phys570/24F/gu_proposal.html$

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World class

- State-of-the-art synchrotron radiation light sources at APS and NSLS-II offer continuous spectrum, high flux and brightness allowing scientists to probe the fundamental properties of matter.
- The free electron laser at LCLS generates ultra-bright, ultrafast, high coherence pulses, with the MeV-UED offering a powerful "electron camera" to study ultrafast atomic & molecular dynamics.

Learn more

- User facilities provide open access to specialized instrumentation to scientists from universities, national laboratories, and industry.
- For approved, peer-reviewed projects, instrument time is available without charge to researchers who intend to publish their results in the open ilkerature.
- Thousands of scientists conduct experiments at BES user facilities every year.

Get started

- Create a free ORCID profile or use your existing ORCID ID to register to use the proposal system.
- Submit a proposal to request experimental time or submit a request against a proposal that has already been awarded time.
- Contact User Program staff with any questions they are there to help!

U.S. DEPARTMENT OF ENERGY OFFICE OF SCIENCE X-RAY LIGHT SOURCES

Participating Facilities

This tool is currently being used to support the proposal submission and review processes for the following facilities



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October 23, 2024

Select the APS



	Carlo Segre Email segre@it.edu			User A	sknowledgement					
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MY PROPOSALS										
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Start a proposal for and APS call

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F	eature Beamlines	Contact Info	Beamlines			
and Da	Title 🔺		Types 🔺	Proposal Cycles 🔺	Deadline 🔺	Proposal Call Status
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	2025-1 Partner User	Proposals (PUP)	Partner Proposals	APS: 2025-1	10/25/2024 21:59:59	SUBMIT A PROPOSA
s.anl.gov/	2025-1 eBERlight Ge	eneral User	General User - Regular	APS: 2025-1	10/25/2024 21:59:59	SUBMIT A PROPOSA
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	2025-1 CAT Member	r Proposals	CAT Member	APS: 2025-1	04/17/2025 21:59:59	SUBMIT A PROPOSA

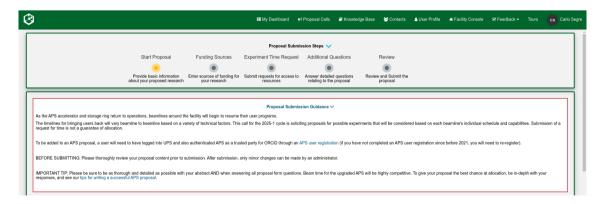
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Enter basic information and an abstract





Enter basic information and an abstract



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2025-1 Standard General User Proposals Proposal Type	Ÿ	Co-Principal Inv	estigator (PI)						
General User - Regular	¥	Co-proposers							
Primary Area of Research Materials science									
Additional Area(s) of Research									
* Keywords									
× x-ray absorption spectroscopy (XAS)									
PTease suggest the most appropriate review panel for your proposal. Spectroscopy-Chem/Catalysis	н т								
Abstract *Abstract									
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Enter basic information and an abstract



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Register & log into the APS Portal

Experiment Time Request * Proposal			
0 1009536	ж	•	Argonne
*Run Cycle			ETR Number
0 APS: 2025-1	ж	*	1033603
1st Choice Resource			2nd Choice Resource
0 10-ID-B	×	*	· · · · · · · · · · · · · · · · · · ·
1st Choice Instrument			2nd Choice Instrument
0 10-ID-B X-ray absorption fine structure	×	*	
1st Choice Technique			2nd Choice Technique
X-ray absorption spectroscopy (XAS)	×	•	
Shifts Requested This ETR			
8			
Minimum Useful Shifts This ETR			
8			
*Lifetime Shifts Requested			
24			

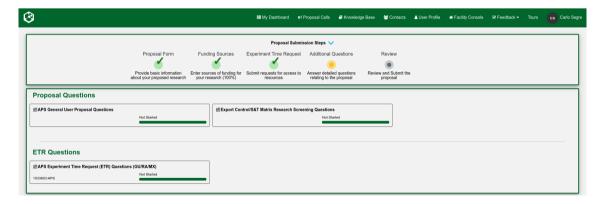
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1009536 APS General User Proposal Questions
General
•What mode(s) of access would you consider for this work? (Note: not all beamlines support all modes of access, choose all th apply.)
Remote
Mail-in
On-site
* Will the data collected be considered proprietary (e.g., work that will not be made available in the open literature)?
○ yes
O no
*Have you spoken to a beamline staff member?
O yes
O no
*Is this research required for a student's thesis?
O yes
O no
Is this proposal related to another proposal?
○ yes
O no

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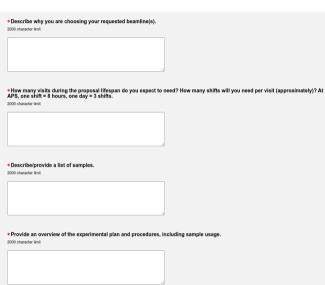


Did you previously receive experiment time at APS for this research.	arch?
⊖ yes	
O no	
• Will you be requesting beam time at APS sector 35, the Dynami	c Compression Sector (DCS)?
⊖ yes	
O no	
Technical	~
lecinical	•
• What is the scientific or technical purpose and importance of th	e proposed research?
*What is the scientific or technical purpose and importance of th 2000 character limit	e proposed research?
	e proposed research?
2000 character limit	e proposed research?
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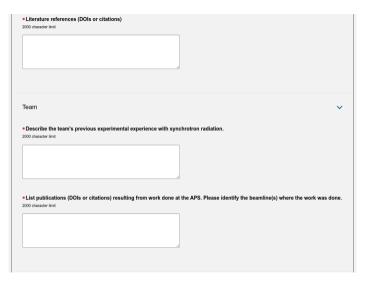
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Safe	ty	~
*Doe	es this research involve the use of radioactive samples/materials, sealed sources, or x-ray generating devices?	
⊖ y	es	
0 n	0	
*Doe	es this research involve the use of any of the following (pick all that apply):	
. e	xplosives or energetic materials	
a	new class 3 or class 4 laser that has not been approved by the Argonne Laser Safety Officer	
🗆 n	anoparticles (one or more dimensions is 100 nm or less), including thin films, powder, and solutions	
S	amples/materials that require a BSL-2 (biosafety level) facility	
🗆 h	uman subjects or human tissues, body fluids, or cells in culture	
🗆 р	lant pathogens, soil microbes, animals, insects, or insect/animal tissues, body fluids, matter, cells in culture	
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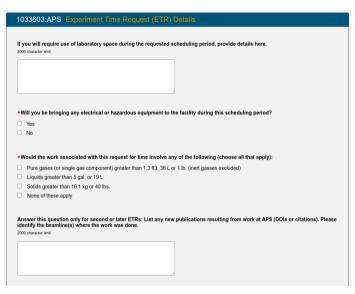
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eactors, nuclear grade graphite, uranium enrichment)?	cial
*	
imens/samples), technical data, or software to the user facility that require restri	cted
*	
al and emerging research areas and technologies. Note: If no or unsure, you sho nsible for screening research for the DOE S&T Risk Matrix. The User Facility mus	uld
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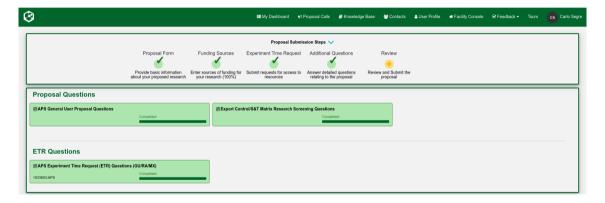


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Preferred experiment dates	for this request, enter date span(s) in format MM/DD/YYYY.	
Unacceptable experiment da	ites for this request, enter date span(s) in format MM/DD/YYYY.	
CANCEL		SAVE





What is the scientific or technical purpose and importance of the proposed research?

A.F

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Why do you need the APS for this research?

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Describe why you are choosing your requested beamline(s).

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Describe/provide a list of samples.

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Describe/provide a list of samples.

Provide an overview of the experimental plan and procedures, including sample usage.

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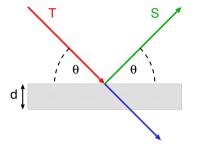
Describe the team's previous experimental experience with synchrotron radiation.

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Darwin approach review





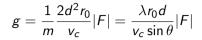
$$g = \frac{1}{m} \frac{2d^2 r_0}{v_c} |F| = \frac{\lambda r_0 d}{v_c \sin \theta} |F|$$

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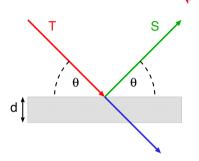
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October 23, 2024

Darwin approach review



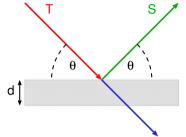
since $v_c \sim d^3$ then $g \sim r_0/d pprox 10^{-5}$



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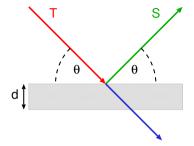




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the transmitted wave is equal in amplitude to the incident wave but gains a phase shift as it passes through the layer

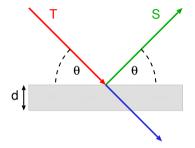
$$T' = (1 - ig_0)T$$

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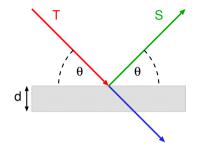
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since $v_c \sim d^3$ then $g \sim r_0/d \approx 10^{-5}$

from Chapter 3

$$g_0 = \frac{\lambda \rho_{at} f^0(0) r_0 d}{\sin \theta}$$





$$T' = (1 - ig_0)T pprox e^{-ig_0}T$$

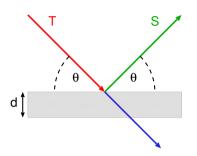
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$$g_0 = \frac{\lambda \rho_{at} f^0(0) r_0 d}{\sin \theta} = \frac{\lambda |F_0| r_0 d}{v_c \sin \theta}$$

where $|F_0| = \rho_{at} f^0(0) v_c$ is the unit cell structure factor in the forward direction at $Q = \theta = 0$



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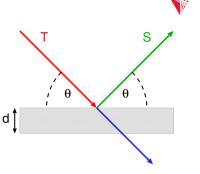
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this can be rewritten in terms of g as



$$T' = (1 - ig_0)T pprox e^{-ig_0}T$$

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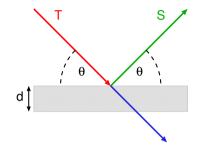
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$$g_0 = g \frac{|F_0|}{|F|}$$





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$$T' = (1 - ig_0)T pprox e^{-ig_0}T$$

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Now extend this model to N layers to get the kinematical scattering approximation as long as the total scattering is weak, $Ng \ll 1$.



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Proceed by adding reflectivity from each layer with the usual phase factor



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$$r_{N}(Q) = -ig\sum_{j=0}^{N-1}e^{iQdj}e^{-ig_{0}j}e^{-ig_{0}j}$$



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$$r_{N}(Q) = -ig\sum_{j=0}^{N-1}e^{iQdj}e^{-ig_{0}j}e^{-ig_{0}j}$$

where the x-rays pass through each layer twice



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$$r_N(Q) = -ig \sum_{j=0}^{N-1} e^{iQdj} e^{-ig_0j} e^{-ig_0j} = -ig \sum_{j=0}^{N-1} e^{i(Qd-2g_0)j}$$

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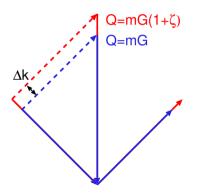
where the x-rays pass through each layer twice

these N unit cell layers will give a reciprocal lattice with points at multiples of $G=2\pi/d$



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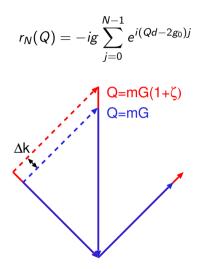
these *N* unit cell layers will give a reciprocal lattice with points at multiples of $G = 2\pi/d$ we are interested in small deviations from the Bragg condition:

$$\zeta = \frac{\Delta Q}{Q} = \frac{\Delta k}{k} = \frac{\Delta \mathcal{E}}{\mathcal{E}} = \frac{\Delta \lambda}{\lambda}$$

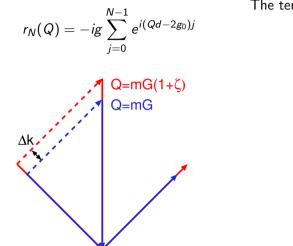
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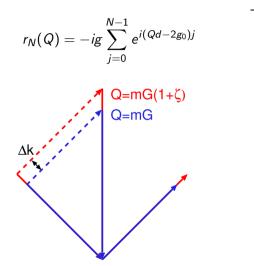




The term in the phase factor now becomes

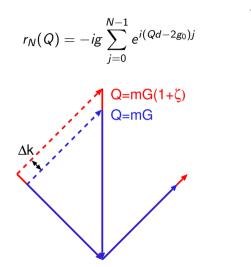
 $Qd - 2g_0$





$$Qd-2g_0=mG(1+\zeta)\frac{2\pi}{G}-2g_0$$





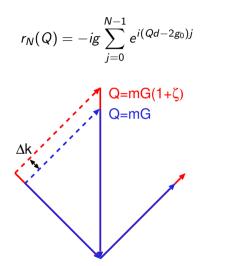
The term in the phase factor now becomes

$$egin{aligned} \mathcal{Q}d-2g_0&=mG(1+\zeta)rac{2\pi}{G}-2g_0\ &=2\pi(m+m\zeta-rac{g_0}{\pi}) \end{aligned}$$

(

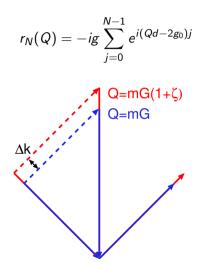


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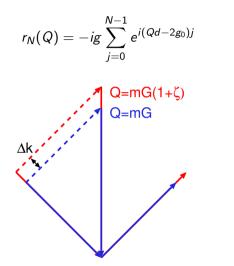
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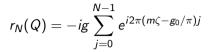




$$Qd - 2g_0 = mG(1 + \zeta)\frac{2\pi}{G} - 2g_0$$

= $2\pi(m + m\zeta - \frac{g_0}{\pi})$
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= $-ig \sum_{j=0}^{N-1} e^{i2\pi mj} e^{i2\pi(m\zeta - g_0/\pi)j}$
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$$\zeta_0 = \frac{g_0}{\pi} = \frac{2d^2|F_0|}{\pi m v_c} r_0$$



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As the crystal becomes infinite ($N \to \infty)$ this kinematical approximation breaks down because $gN \sim 1$

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Diffraction in the kinematical limit

V

It is useful to look at how the intensity of the reflection varies in the kinematical limit

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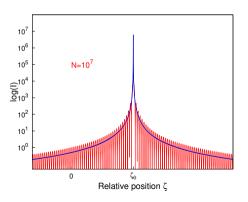
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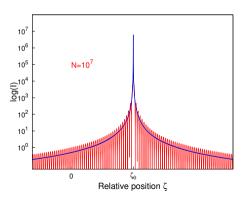
V

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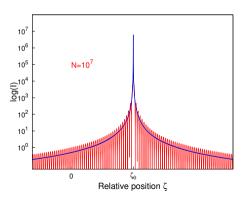
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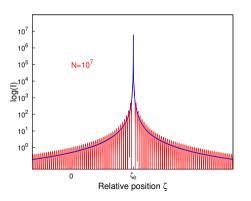
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26/31

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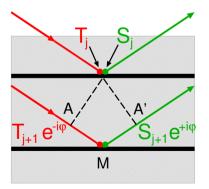
The kinematical limit clearly breaks down near ζ_0 so we need a dynamical diffraction theory

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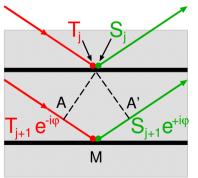
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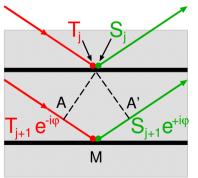
As the wavefields pass through an atomic plane, they experience an abrupt change with a small amount, -ig, of the wave being reflected and a phase shift, $(1 - ig_0)$, being added to the transmitted wave





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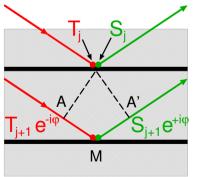
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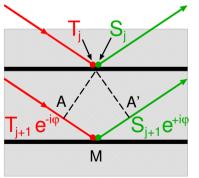


At the Bragg condition, the wave from the $j + 1^{th}$ plane must be in phase with the one from the j^{th} plane, or



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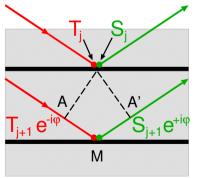


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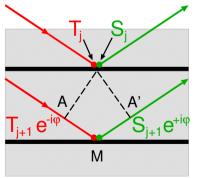
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V

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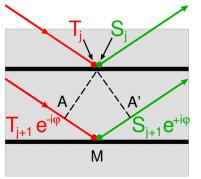
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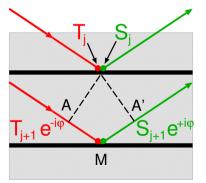
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October 23, 2024

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Let T_j and S_j be the fields just above layer j.

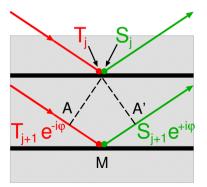




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Let T_j and S_j be the fields just above layer j.

at point M, just above the $j + 1^{th}$ layer, we have the scattered field S_{j+1} and at point A' it is $S_{j+1}e^{i\phi}$

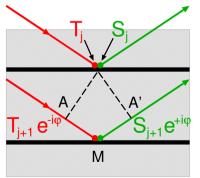




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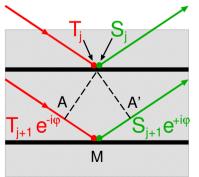
 $S_j= (1-ig_0)S_{j+1}e^{i\phi}$



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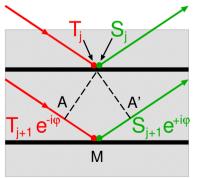
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similarly we can write an equation for T_{j+1} just below the j^{th} plane

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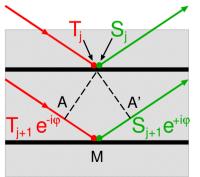


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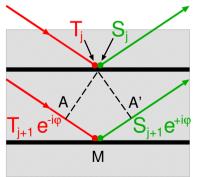
$$S_{j} = -ig T_{j+1} + (1 - ig_{0})S_{j+1}e^{i\phi}$$
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these coupled equations must be solved for an infinite stack of atomic layers

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Rearranging the equation for T_i (top right)



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$$S_j = -ig T_{j+1} + (1 - ig_0) S_{j+1} e^{i\phi}, \quad (1 - ig_0) T_j = T_{j+1}$$

Rearranging the equation for T_j (top right) shifting up by one plane: $j + 1 \rightarrow j$ and $j \rightarrow j - 1$

now substitute into the equation for S_i above

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 $-i\phi$ $i\pi c$ $i\phi$



$$S_{j} = -ig T_{j+1} + (1 - ig_{0})S_{j+1}e^{i\phi}, \quad (1 - ig_{0})T_{j} = T_{j+1}e^{-i\phi} + ig S_{j+1}e^{i\phi}$$

Rearranging the equation for T_j (top right) shifting up by one plane: $j+1 \rightarrow j$ and $j \rightarrow j-1$

$$igS_{j+1} = (1 - ig_0)T_je^{-i\phi} - T_{j+1}e^{-i2\phi}$$

$$igS_j = (1 - ig_0)T_{j-1}e^{-i\phi} - T_je^{-i2\phi}$$

now substitute into the equation for S_j above

$$(1 - ig_0)T_{j-1}e^{-i\phi} - T_je^{-i2\phi} = g^2T_j + (1 - ig_0)\left[(1 - ig_0)T_j - T_{j+1}e^{-i\phi}\right]$$



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$$(1 - ig_0)e^{-i\phi}[T_{j+1} + T_{j-1}] = \left[g^2 + (1 - ig_0)^2 + e^{-i2\phi}\right]T_j$$



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$$(1 - ig_0)e^{-i\phi}[T_{j+1} + T_{j-1}] = \left[g^2 + (1 - ig_0)^2 + e^{-i2\phi}\right]T_j$$

the fields T_j and T_{j+1} are out of phase by nearly $m\pi$ (top right equation) since g and g_0 are very small and the T wave field must attenuate as it penetrates deeper into the crystal so our trial solution is

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$$S_{j} = -ig T_{j+1} + (1 - ig_{0})S_{j+1}e^{i\phi}, \quad (1 - ig_{0})T_{j} = T_{j+1}e^{-i\phi} + ig S_{j+1}e^{i\phi}$$

Rearranging the equation for T_j (top right) shifting up by one plane: $j + 1 \rightarrow j$ and $j \rightarrow j - 1$

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$$T_{j+1} = e^{-\eta} e^{im\pi} T_j$$

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$$(1 - ig_0)e^{-i\phi} \left[T_{j+1} + T_{j-1} \right] = \left[g^2 + (1 - ig_0)^2 + e^{-i2\phi} \right] T_j$$

,



$$(1 - ig_0)e^{-i\phi} \left[T_{j+1} + T_{j-1} \right] = \left[g^2 + (1 - ig_0)^2 + e^{-i2\phi} \right] T_j$$

With the trial solution

,



$$(1 - ig_0)e^{-i\phi}\left[T_{j+1} + T_{j-1}\right] = \left[g^2 + (1 - ig_0)^2 + e^{-i2\phi}\right]T_j$$

With the trial solution
$$T_{j+1} = e^{-\eta}e^{im\pi}T_j,$$

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$$(1 - ig_0)e^{-i\phi} [T_{j+1} + T_{j-1}] = \left[g^2 + (1 - ig_0)^2 + e^{-i2\phi}\right] T_j$$

With the trial solution $T_{j+1} = e^{-\eta}e^{im\pi}T_j, \quad T_{j-1} = e^{\eta}e^{-im\pi}T_j$

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$$(1 - ig_0)e^{-i\phi} [T_{j+1} + T_{j-1}] = \left[g^2 + (1 - ig_0)^2 + e^{-i2\phi}\right] T_j$$

With the trial solution $T_{j+1} = e^{-\eta}e^{im\pi}T_j, \quad T_{j-1} = e^{\eta}e^{-im\pi}T_j$

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$$(1 - ig_0)e^{-i\phi}\left[e^{-\eta}e^{im\pi}T_j + e^{\eta}e^{-im\pi}T_j\right] = \left[g^2 + (1 - ig_0)^2 + e^{-i2\phi}\right]T_j$$



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and substituting this solution into the defining equation for ${\cal T}$ and noting that $\phi\equiv m\pi+\Delta$

$$(1 - ig_0)e^{-i\phi}\left[e^{-\eta}e^{im\pi}T_j + e^{\eta}e^{-im\pi}T_j\right] = \left[g^2 + (1 - ig_0)^2 + e^{-i2\phi}\right]T_j$$



$$(1 - ig_0)e^{-i\phi}[T_{j+1} + T_{j-1}] = \left[g^2 + (1 - ig_0)^2 + e^{-i2\phi}\right]T_j$$

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$$(1 - ig_0)e^{-i\phi} \left[e^{-\eta}e^{im\pi}T_j + e^{\eta}e^{-im\pi}T_j \right] = \left[g^2 + (1 - ig_0)^2 + e^{-i2\phi} \right] T_j$$

$$(1 - ig_0)e^{-im\pi}e^{-i\Delta} \left[e^{-\eta}e^{im\pi} + e^{\eta}e^{-im\pi} \right] = g^2 + (1 - ig_0)^2 + e^{-i2m\pi}e^{-i2\Delta}$$



$$(1 - ig_0)e^{-i\phi} [T_{j+1} + T_{j-1}] = \left[g^2 + (1 - ig_0)^2 + e^{-i2\phi}\right] T_j$$

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$$(1 - ig_0)e^{-i\phi} [T_{j+1} + T_{j-1}] = \left[g^2 + (1 - ig_0)^2 + e^{-i2\phi}\right] T_j$$

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$$(1 - ig_0)e^{-i\Delta} \left[e^{-\eta} + e^{\eta} \right] = g^2 + (1 - ig_0)^2 + e^{-i2\Delta}$$

assuming that g, g_0 , and Δ are very small quantities, we can expand



$$(1 - ig_0)e^{-i\phi} [T_{j+1} + T_{j-1}] = \left[g^2 + (1 - ig_0)^2 + e^{-i2\phi}\right] T_j$$

With the trial solution $T_{j+1} = e^{-\eta}e^{im\pi}T_j, \quad T_{j-1} = e^{\eta}e^{-im\pi}T_j$

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$$(1 - ig_0)e^{-i\phi} \left[e^{-\eta}e^{im\pi}T_j + e^{\eta}e^{-im\pi}T_j \right] = \left[g^2 + (1 - ig_0)^2 + e^{-i2\phi} \right] T_j$$

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$$(1 - ig_0)e^{-i\Delta} \left[e^{-\eta} + e^{\eta} \right] = g^2 + (1 - ig_0)^2 + e^{-i2\Delta}$$

assuming that g, $g_{0},$ and Δ are very small quantities, we can expand

$$(1 - ig_0)(1 - i\Delta - rac{\Delta^2}{2})\left[(1 - \eta + rac{\eta^2}{2}) + (1 + \eta + rac{\eta^2}{2})
ight] \ pprox g^2 + (1 - 2ig_0 - g_0^2) + (1 - i2\Delta - 2\Delta^2)$$

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$$egin{aligned} &(1-ig_0)(1-i\Delta-rac{\Delta^2}{2})\left[(1-\eta+rac{\eta^2}{2})+(1+\eta+rac{\eta^2}{2})
ight] \ &pprox g^2+(1-2ig_0-g_0^2)+(1-i2\Delta-2\Delta^2) \end{aligned}$$



$$(1 - ig_0)(1 - i\Delta - \frac{\Delta^2}{2}) \left[(1 - \eta + \frac{\eta^2}{2}) + (1 + \eta + \frac{\eta^2}{2})
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$$(1 - ig_0)(1 - i\Delta - \frac{\Delta^2}{2}) \left[(1 - \eta + \frac{\eta^2}{2}) + (1 + \eta + \frac{\eta^2}{2})
ight] \ pprox g^2 + (1 - 2ig_0 - g_0^2) + (1 - i2\Delta - 2\Delta^2)$$

$$(1 - ig_0 - i\Delta - g_0\Delta - \frac{\Delta^2}{2})(2 + \eta^2) \approx g^2 + 2 - 2ig_0 - 2i\Delta - g_0^2 - 2\Delta^2$$



$$(1 - ig_0)(1 - i\Delta - \frac{\Delta^2}{2}) \left[(1 - \eta + \frac{\eta^2}{2}) + (1 + \eta + \frac{\eta^2}{2})
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 $2 - 2ig_0 - 2i\Delta - 2g_0\Delta - \Delta^2 + \eta^2 \approx g^2 + 2 - 2ig_0 - 2i\Delta - g_0^2 - 2\Delta^2$



$$(1 - ig_0)(1 - i\Delta - \frac{\Delta^2}{2}) \left[(1 - \eta + \frac{\eta^2}{2}) + (1 + \eta + \frac{\eta^2}{2})
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Cancelling and expanding all products keeping only second order terms

$$(1 - ig_0 - i\Delta - g_0\Delta - \frac{\Delta^2}{2})(2 + \eta^2) \approx g^2 + 2 - 2ig_0 - 2i\Delta - g_0^2 - 2\Delta^2$$

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The solution for the attenuation factor of the transmitted field is thus



$$egin{aligned} &(1-ig_0)(1-i\Delta-rac{\Delta^2}{2})\left[(1-\eta+rac{\eta^2}{2})+(1+\eta+rac{\eta^2}{2})
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$$i\eta = \pm \sqrt{(\Delta - g_0) - g^2}$$

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The solution for the attenuation factor of the transmitted field is thus

$$i\eta=\pm\sqrt{(\Delta-g_0)-g^2}$$

with fields

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$$(1 - ig_0)(1 - i\Delta - \frac{\Delta^2}{2}) \left[(1 - \eta + \frac{\eta^2}{2}) + (1 + \eta + \frac{\eta^2}{2})
ight] \ \approx g^2 + (1 - 2ig_0 - g_0^2) + (1 - i2\Delta - 2\Delta^2)$$

Cancelling and expanding all products keeping only second order terms

$$(1 - ig_0 - i\Delta - g_0\Delta - \frac{\Delta^2}{2})(2 + \eta^2) \approx g^2 + 2 - 2ig_0 - 2i\Delta - g_0^2 - 2\Delta^2$$

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$$T_{j+1} = e^{-\eta} e^{im\pi} T_j,$$

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$$(1 - ig_0)(1 - i\Delta - \frac{\Delta^2}{2}) \left[(1 - \eta + \frac{\eta^2}{2}) + (1 + \eta + \frac{\eta^2}{2})
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Cancelling and expanding all products keeping only second order terms

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The solution for the attenuation factor of the transmitted field is thus

$$i\eta = \pm \sqrt{(\Delta - g_0) - g^2}$$

with fields

$$T_{j+1} = e^{-\eta} e^{im\pi} T_j, \quad S_{j+1} = e^{-\eta} e^{im\pi} S_j$$

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