



• Molecule scattering



- Molecule scattering
- Liquid scattering



- Molecule scattering
- Liquid scattering
- Small angle x-ray scattering



- Molecule scattering
- Liquid scattering
- Small angle x-ray scattering
- Calculating R<sub>g</sub>



- Molecule scattering
- Liquid scattering
- Small angle x-ray scattering
- Calculating R<sub>g</sub>
- Porod analysis



1/20

- Molecule scattering
- Liquid scattering
- Small angle x-ray scattering
- Calculating R<sub>g</sub>
- Porod analysis

Reading Assignment: Chapter 4.5; Chapter 5.1



1/20

- Molecule scattering
- Liquid scattering
- Small angle x-ray scattering
- Calculating R<sub>g</sub>
- Porod analysis

Reading Assignment: Chapter 4.5; Chapter 5.1

Homework Assignment #03: Chapter 3: 1,3,4,6,8 due Friday, October 05, 2024

V

- Molecule scattering
- Liquid scattering
- Small angle x-ray scattering
- Calculating R<sub>g</sub>
- Porod analysis

Reading Assignment: Chapter 4.5; Chapter 5.1

Homework Assignment #03: Chapter 3: 1,3,4,6,8 due Friday, October 05, 2024 Homework Assignment #04: Chapter 4: 2,4,6,7,10 due Monday, October 14, 2024



From the atomic form factor, we would like to abstract to the next level of complexity, a molecule (we will leave crystals for Chapter 5).

From the atomic form factor, we would like to abstract to the next level of complexity, a molecule (we will leave crystals for Chapter 5).

$$F^{mol}(ec{Q}) = \sum_j f_j(ec{Q}) e^{iec{Q}\cdotec{r}_j}$$

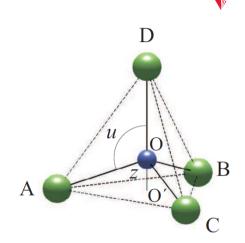


2 / 20

From the atomic form factor, we would like to abstract to the next level of complexity, a molecule (we will leave crystals for Chapter 5).

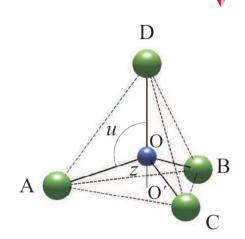
$${\cal F}^{mol}(ec Q) = \sum_j f_j(ec Q) e^{iec Q\cdot ec r_j}$$

As an example take the  $\mathsf{CF}_4$  molecule



From the atomic form factor, we would like to abstract to the next level of complexity, a molecule (we will leave crystals for Chapter 5).

$${\cal F}^{mol}(ec Q) = \sum_j f_j(ec Q) e^{iec Q\cdot ec r_j}$$

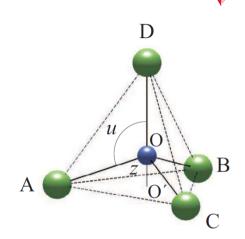


From the atomic form factor, we would like to abstract to the next level of complexity, a molecule (we will leave crystals for Chapter 5).

$${\cal F}^{mol}(ec Q) = \sum_j f_j(ec Q) e^{iec Q\cdot ec r_j}$$

As an example take the CF<sub>4</sub> molecule We have the following relationships:

 $\overline{OA}$ 

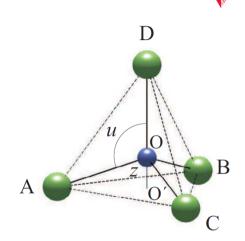


From the atomic form factor, we would like to abstract to the next level of complexity, a molecule (we will leave crystals for Chapter 5).

$${\cal F}^{mol}(ec Q) = \sum_j f_j(ec Q) e^{iec Q\cdot ec r_j}$$

As an example take the CF<sub>4</sub> molecule We have the following relationships:

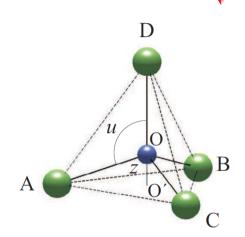
 $|\overline{OA}| = |\overline{OB}|$ 



From the atomic form factor, we would like to abstract to the next level of complexity, a molecule (we will leave crystals for Chapter 5).

$${\cal F}^{mol}(ec Q) = \sum_j f_j(ec Q) e^{iec Q\cdot ec r_j}$$

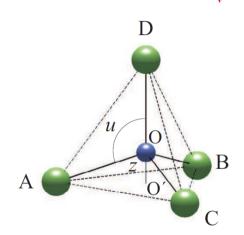
$$|\overline{OA}| = |\overline{OB}| = |\overline{OC}|$$



From the atomic form factor, we would like to abstract to the next level of complexity, a molecule (we will leave crystals for Chapter 5).

$${\cal F}^{mol}(ec Q) = \sum_j f_j(ec Q) e^{iec Q\cdot ec r_j}$$

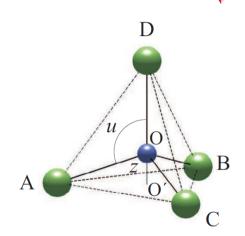
$$|\overline{OA}| = |\overline{OB}| = |\overline{OC}| = |\overline{OD}|$$



From the atomic form factor, we would like to abstract to the next level of complexity, a molecule (we will leave crystals for Chapter 5).

$${\cal F}^{mol}(ec Q) = \sum_j f_j(ec Q) e^{iec Q\cdot ec r_j}$$

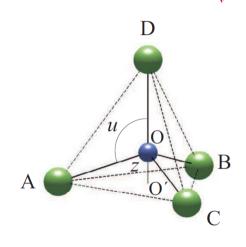
$$|\overline{\textit{OA}}| = |\overline{\textit{OB}}| = |\overline{\textit{OC}}| = |\overline{\textit{OD}}| = 1$$



From the atomic form factor, we would like to abstract to the next level of complexity, a molecule (we will leave crystals for Chapter 5).

$${\cal F}^{mol}(ec Q) = \sum_j f_j(ec Q) e^{iec Q\cdot ec r_j}$$

$$|\overline{\textit{OA}}| = |\overline{\textit{OB}}| = |\overline{\textit{OC}}| = |\overline{\textit{OD}}| = 1$$

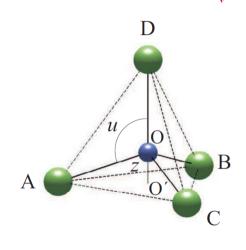


$$\overline{OA} \cdot \overline{OD} = 1 \cdot 1 \cdot \cos u$$

From the atomic form factor, we would like to abstract to the next level of complexity, a molecule (we will leave crystals for Chapter 5).

$${\cal F}^{mol}(ec Q) = \sum_j f_j(ec Q) e^{iec Q\cdot ec r_j}$$

$$|\overline{\textit{OA}}| = |\overline{\textit{OB}}| = |\overline{\textit{OC}}| = |\overline{\textit{OD}}| = 1$$

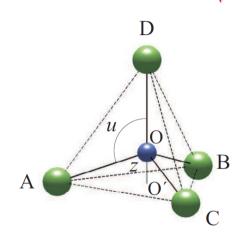


$$\overline{OA} \cdot \overline{OD} = 1 \cdot 1 \cdot \cos u = -z$$

From the atomic form factor, we would like to abstract to the next level of complexity, a molecule (we will leave crystals for Chapter 5).

$${\cal F}^{mol}(ec Q) = \sum_j f_j(ec Q) e^{iec Q\cdot ec r_j}$$

$$|\overline{\textit{OA}}| = |\overline{\textit{OB}}| = |\overline{\textit{OC}}| = |\overline{\textit{OD}}| = 1$$

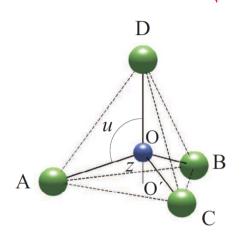


$$\overline{OA} \cdot \overline{OD} = 1 \cdot 1 \cdot \cos u = -z$$
$$= \overline{OA} \cdot \overline{OB}$$

From the atomic form factor, we would like to abstract to the next level of complexity, a molecule (we will leave crystals for Chapter 5).

$${\cal F}^{mol}(ec Q) = \sum_j f_j(ec Q) e^{iec Q\cdot ec r_j}$$

$$\begin{split} |\overline{OA}| &= |\overline{OB}| = |\overline{OC}| = |\overline{OD}| = 1\\ \overline{OA} &= \overline{OO'} + \overline{O'A} \end{split}$$



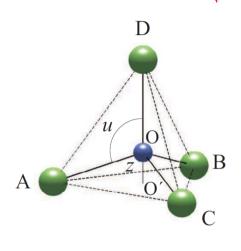
$$\overline{OA} \cdot \overline{OD} = 1 \cdot 1 \cdot \cos u = -z$$
$$= \overline{OA} \cdot \overline{OB}$$

From the atomic form factor, we would like to abstract to the next level of complexity, a molecule (we will leave crystals for Chapter 5).

$${\cal F}^{mol}(ec Q) = \sum_j f_j(ec Q) e^{iec Q\cdot ec r_j}$$

As an example take the CF<sub>4</sub> molecule We have the following relationships:

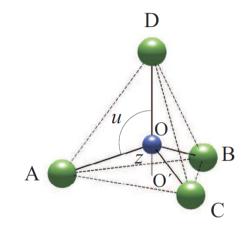
$$\begin{split} |\overline{OA}| &= |\overline{OB}| = |\overline{OC}| = |\overline{OD}| = 1\\ \overline{OA} &= \overline{OO'} + \overline{O'A}\\ \overline{OB} &= \overline{OO'} + \overline{O'B} \end{split}$$



$$\overline{OA} \cdot \overline{OD} = 1 \cdot 1 \cdot \cos u = -z$$
$$= \overline{OA} \cdot \overline{OB}$$

Carlo Segre (Illinois Tech)





$$-z = (\overline{OO'} + \overline{O'A}) \cdot (\overline{OO'} + \overline{O'B})$$

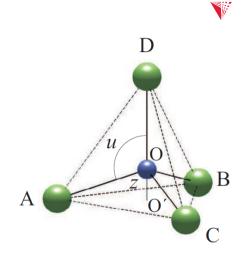
Carlo Segre (Illinois Tech)

PHYS 570 - Fall 2024

September 30, 2024 3 /

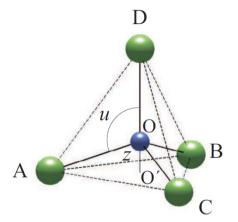
3 / 20

$$-z = (\overline{OO'} + \overline{O'A}) \cdot (\overline{OO'} + \overline{O'B})$$
$$= z^2 + 0 + 0 + \overline{O'A} \cdot \overline{O'B}$$









$$-z = (\overline{OO'} + \overline{O'A}) \cdot (\overline{OO'} + \overline{O'B})$$
  
=  $z^2 + 0 + 0 + \overline{O'A} \cdot \overline{O'B} = z^2 + (O'A)^2 \cos(120^\circ)$ 

Carlo Segre (Illinois Tech)

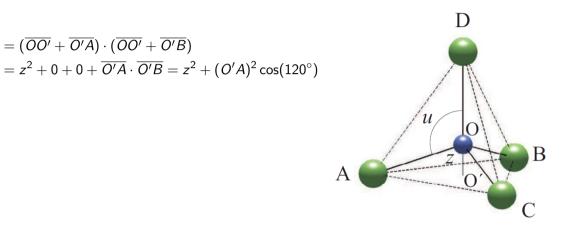
PHYS 570 - Fall 2024

September 30, 2024 3 /

3 / 20

 $-z = (\overline{OO'} + \overline{O'A}) \cdot (\overline{OO'} + \overline{O'B})$ 





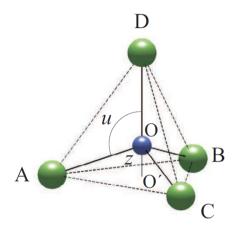
but from the triangle OO'A

Carlo Segre (Illinois Tech)

 $-z = (\overline{OO'} + \overline{O'A}) \cdot (\overline{OO'} + \overline{O'B})$ 

 $= z^2 + 0 + 0 + \overline{O'A} \cdot \overline{O'B} = z^2 + (O'A)^2 \cos(120^\circ)$ 





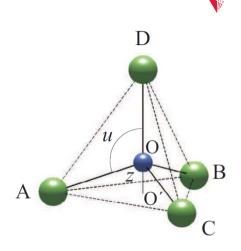
but from the triangle OO'A

$$(O'A)^2 = 1 - z^2$$

September 30, 2024 3 / 20

Carlo Segre (Illinois Tech)

$$-z = (\overline{OO'} + \overline{O'A}) \cdot (\overline{OO'} + \overline{O'B})$$
  
=  $z^2 + 0 + 0 + \overline{O'A} \cdot \overline{O'B} = z^2 + (O'A)^2 \cos(120^\circ)$   
=  $z^2 + (1 - z^2) \cos(120^\circ)$ 

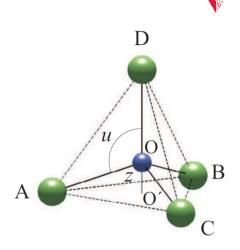


but from the triangle OO'A

$$(O'A)^2 = 1 - z^2$$

Carlo Segre (Illinois Tech)

$$-z = (\overline{OO'} + \overline{O'A}) \cdot (\overline{OO'} + \overline{O'B})$$
  
=  $z^2 + 0 + 0 + \overline{O'A} \cdot \overline{O'B} = z^2 + (O'A)^2 \cos(120^\circ)$   
=  $z^2 + (1 - z^2) \cos(120^\circ) = z^2 - \frac{1}{2}(1 - z^2)$ 

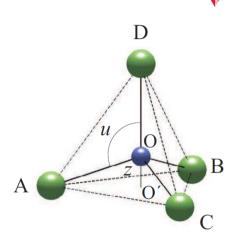


but from the triangle OO'A

$$(O'A)^2 = 1 - z^2$$

$$-z = (\overline{OO'} + \overline{O'A}) \cdot (\overline{OO'} + \overline{O'B})$$
  
=  $z^2 + 0 + 0 + \overline{O'A} \cdot \overline{O'B} = z^2 + (O'A)^2 \cos(120^\circ)$   
=  $z^2 + (1 - z^2) \cos(120^\circ) = z^2 - \frac{1}{2}(1 - z^2)$ 

$$0 = 3z^2 + 2z - 1$$



but from the triangle OO'A

$$(O'A)^2 = 1 - z^2$$

Carlo Segre (Illinois Tech)

$$-z = (\overline{OO'} + \overline{O'A}) \cdot (\overline{OO'} + \overline{O'B})$$
  
=  $z^2 + 0 + 0 + \overline{O'A} \cdot \overline{O'B} = z^2 + (O'A)^2 \cos(120^\circ)$   
=  $z^2 + (1 - z^2) \cos(120^\circ) = z^2 - \frac{1}{2}(1 - z^2)$   
 $0 = 3z^2 + 2z - 1$   
1

$$z=\frac{1}{3}$$

D U R -----А

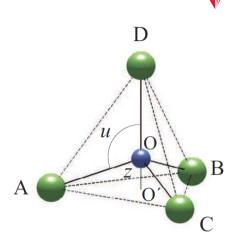
but from the triangle OO'A

$$(O'A)^2 = 1 - z^2$$



$$-z = (\overline{OO'} + \overline{O'A}) \cdot (\overline{OO'} + \overline{O'B})$$
  
=  $z^2 + 0 + 0 + \overline{O'A} \cdot \overline{O'B} = z^2 + (O'A)^2 \cos(120^\circ)$   
=  $z^2 + (1 - z^2) \cos(120^\circ) = z^2 - \frac{1}{2}(1 - z^2)$   
$$0 = 3z^2 + 2z - 1$$
  
$$z = \frac{1}{3}$$

$$u = \cos^{-1}(-z) = 109.5^{\circ}$$



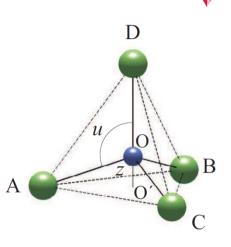
but from the triangle OO'A

$$(O'A)^2 = 1 - z^2$$

$$-z = (\overline{OO'} + \overline{O'A}) \cdot (\overline{OO'} + \overline{O'B})$$
  
=  $z^2 + 0 + 0 + \overline{O'A} \cdot \overline{O'B} = z^2 + (O'A)^2 \cos(120^\circ)$   
=  $z^2 + (1 - z^2) \cos(120^\circ) = z^2 - \frac{1}{2}(1 - z^2)$   
$$0 = 3z^2 + 2z - 1$$
  
$$z = \frac{1}{3}$$

$$u = \cos^{-1}(-z) = 109.5^{\circ}$$

$$F_{\pm}^{mol} = f^{C}(Q) + f^{F}(Q) \left[ 3e^{\pm iQR/3} + e^{\pm iQR} \right]$$



but from the triangle OO'A

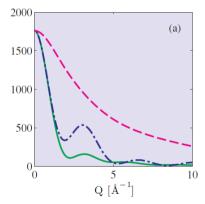
$$(O'A)^2 = 1 - z^2$$

Carlo Segre (Illinois Tech)



$$F_{\pm}^{mol} = f^{C}(Q) + f^{F}(Q) \left[ 3e^{\mp iQR/3} + e^{\pm iQR} \right]$$

$$F_{\pm}^{mol} = f^{C}(Q) + f^{F}(Q) \left[ 3e^{\pm iQR/3} + e^{\pm iQR} \right]$$
  
$$F_{\pm}^{mol}|^{2} = |f^{C}|^{2} + 10|f^{F}|^{2} + 6|f^{F}|^{2} \cos(QR/3) + 2f^{C}f^{F} \left[ 3\cos(QR/3) + \cos(QR) \right]$$



The plot shows the structure factor of  $CF_4$ ,

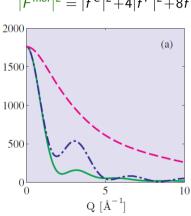
Carlo Segre (Illinois Tech)





$$F_{\pm}^{mol} = f^{C}(Q) + f^{F}(Q) \left[ 3e^{\pm iQR/3} + e^{\pm iQR} \right]$$
$$|F_{\pm}^{mol}|^{2} = |f^{C}|^{2} + 10|f^{F}|^{2} + 6|f^{F}|^{2} \cos(QR/3) + 2f^{C}f^{F} \left[ 3\cos(QR/3) + \cos(QR) \right]$$
$$|F^{mol}|^{2} = |f^{C}|^{2} + 4|f^{F}|^{2} + 8f^{C}f^{F} \frac{\sin(QR)}{QR} + 12|f^{F}|^{2} \frac{\sin(Q\sqrt{8/3}R)}{Q\sqrt{8/3}R}$$

The plot shows the structure factor of  $CF_4$ , its orientationally averaged structure factor,





$$F_{\pm}^{mol} = f^{C}(Q) + f^{F}(Q) \left[ 3e^{\pm iQR/3} + e^{\pm iQR} \right]$$
$$|F_{\pm}^{mol}|^{2} = |f^{C}|^{2} + 10|f^{F}|^{2} + 6|f^{F}|^{2} \cos(QR/3) + 2f^{C}f^{F} \left[ 3\cos(QR/3) + \cos(QR) \right]$$
$$|F^{mol}|^{2} = |f^{C}|^{2} + 4|f^{F}|^{2} + 8f^{C}f^{F} \frac{\sin(QR)}{QR} + 12|f^{F}|^{2} \frac{\sin(Q\sqrt{8/3}R)}{Q\sqrt{8/3}R}$$

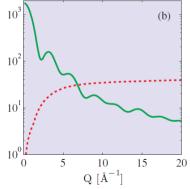
2000 (a) 1500 1000 500 0 í٥ 5 10 [Å -1Q

The plot shows the structure factor of  $CF_4$ , its orientationally averaged structure factor, and the form factor factor of Mo which has the same number of electrons as  $CF_4$ 

Carlo Segre (Illinois Tech)



$$F_{\pm}^{mol} = f^{C}(Q) + f^{F}(Q) \left[ 3e^{\pm iQR/3} + e^{\pm iQR} \right]$$
$$|F_{\pm}^{mol}|^{2} = |f^{C}|^{2} + 10|f^{F}|^{2} + 6|f^{F}|^{2} \cos(QR/3) + 2f^{C}f^{F} \left[ 3\cos(QR/3) + \cos(QR) \right]$$
$$|F^{mol}|^{2} = |f^{C}|^{2} + 4|f^{F}|^{2} + 8f^{C}f^{F} \frac{\sin(QR)}{QR} + 12|f^{F}|^{2} \frac{\sin(Q\sqrt{8/3}R)}{Q\sqrt{8/3}R}$$

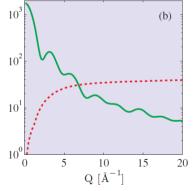


The plot shows the structure factor of  $CF_4$ , its orientationally averaged structure factor, and the form factor factor of Mo which has the same number of electrons as  $CF_4$ 

Carlo Segre (Illinois Tech)



$$F_{\pm}^{mol} = f^{C}(Q) + f^{F}(Q) \left[ 3e^{\pm iQR/3} + e^{\pm iQR} \right]$$
$$|F_{\pm}^{mol}|^{2} = |f^{C}|^{2} + 10|f^{F}|^{2} + 6|f^{F}|^{2} \cos(QR/3) + 2f^{C}f^{F} \left[ 3\cos(QR/3) + \cos(QR) \right]$$
$$|F^{mol}|^{2} = |f^{C}|^{2} + 4|f^{F}|^{2} + 8f^{C}f^{F} \frac{\sin(QR)}{QR} + 12|f^{F}|^{2} \frac{\sin(Q\sqrt{8/3}R)}{Q\sqrt{8/3}R}$$



The plot shows the structure factor of  $CF_4$ , its orientationally averaged structure factor, and the form factor factor of Mo which has the same number of electrons as  $CF_4$ 

The logarithmic plot shows the spherically averaged structure factor compared to the inelastic scattering for  $\mathsf{CF}_4$ 

Carlo Segre (Illinois Tech)

PHYS 570 - Fall 2024

September 30, 2024

4 / 20



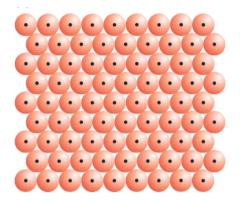
V

Ordered 2D crystal

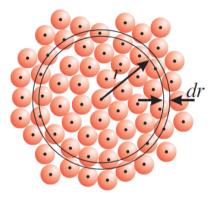
Amorphous solid or liquid



#### Ordered 2D crystal

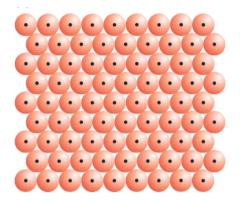


#### Amorphous solid or liquid

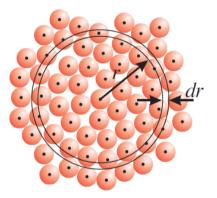




#### Ordered 2D crystal



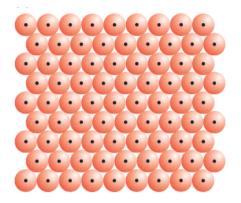
#### Amorphous solid or liquid



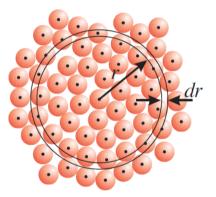
Take a circle (sphere) of radius r and thickness dr and count the number of atom centers lying within the ring.



#### Ordered 2D crystal



#### Amorphous solid or liquid

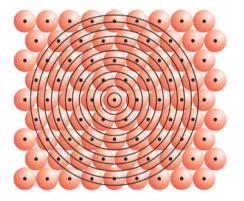


Take a circle (sphere) of radius r and thickness dr and count the number of atom centers lying within the ring. Then expand the ring radius by dr to map out the radial distribution function g(r)

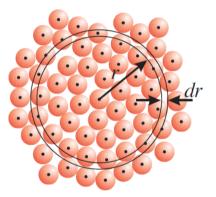
Carlo Segre (Illinois Tech)



#### Ordered 2D crystal



#### Amorphous solid or liquid

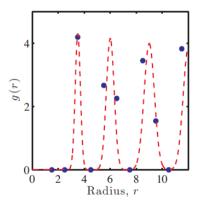


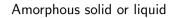
Take a circle (sphere) of radius r and thickness dr and count the number of atom centers lying within the ring. Then expand the ring radius by dr to map out the radial distribution function g(r)

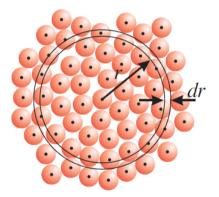
Carlo Segre (Illinois Tech)



Ordered 2D crystal





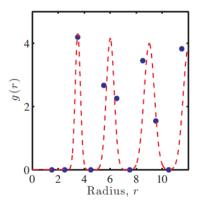


Take a circle (sphere) of radius r and thickness dr and count the number of atom centers lying within the ring. Then expand the ring radius by dr to map out the radial distribution function g(r)

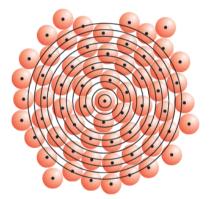
Carlo Segre (Illinois Tech)



Ordered 2D crystal



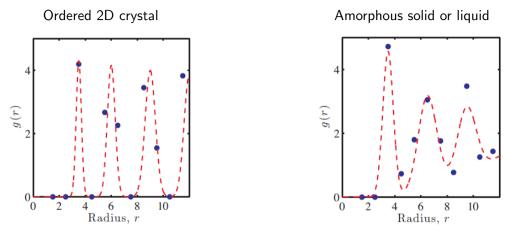
Amorphous solid or liquid



Take a circle (sphere) of radius r and thickness dr and count the number of atom centers lying within the ring. Then expand the ring radius by dr to map out the radial distribution function g(r)

Carlo Segre (Illinois Tech)





Take a circle (sphere) of radius r and thickness dr and count the number of atom centers lying within the ring. Then expand the ring radius by dr to map out the radial distribution function g(r)

Carlo Segre (Illinois Tech)





$$I(\vec{Q}) = f(\vec{Q})^2 \sum_n e^{i \vec{Q} \cdot \vec{r_n}} \sum_m e^{-i \vec{Q} \cdot \vec{r_m}}$$



$$I(\vec{Q}) = f(\vec{Q})^2 \sum_{n} e^{i\vec{Q}\cdot\vec{r_n}} \sum_{m} e^{-i\vec{Q}\cdot\vec{r_m}} = f(\vec{Q})^2 \sum_{n} \sum_{m} e^{i\vec{Q}\cdot(\vec{r_n}-\vec{r_m})}$$

$$I(\vec{Q}) = f(\vec{Q})^2 \sum_{n} e^{i\vec{Q}\cdot\vec{r_n}} \sum_{m} e^{-i\vec{Q}\cdot\vec{r_m}} = f(\vec{Q})^2 \sum_{n} \sum_{m} e^{i\vec{Q}\cdot(\vec{r_n}-\vec{r_m})}$$
$$= Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_{n} \sum_{m \neq n} e^{i\vec{Q}\cdot(\vec{r_n}-\vec{r_m})}$$





6/20

Consider a mono-atomic (-molecular) system where the total scattered intensity is given by

$$\begin{split} I(\vec{Q}) &= f(\vec{Q})^2 \sum_n e^{i\vec{Q}\cdot\vec{r_n}} \sum_m e^{-i\vec{Q}\cdot\vec{r_m}} = f(\vec{Q})^2 \sum_n \sum_m e^{i\vec{Q}\cdot(\vec{r_n}-\vec{r_m})} \\ &= Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \sum_{m \neq n} e^{i\vec{Q}\cdot(\vec{r_n}-\vec{r_m})} \end{split}$$

The sum over  $m \neq m$  is now replaced with an integral of the continuous atomic pair density function,  $\rho_n(\vec{r}_{nm})$ 

$$I(\vec{Q}) = Nf(\vec{Q})^{2} + f(\vec{Q})^{2} \sum_{n} \int_{V} \rho_{n}(\vec{r}_{nm}) e^{i\vec{Q}\cdot(\vec{r}_{n}-\vec{r}_{m})} dV_{m}$$

V

Consider a mono-atomic (-molecular) system where the total scattered intensity is given by

$$I(\vec{Q}) = f(\vec{Q})^2 \sum_{n} e^{i\vec{Q}\cdot\vec{r_n}} \sum_{m} e^{-i\vec{Q}\cdot\vec{r_m}} = f(\vec{Q})^2 \sum_{n} \sum_{m} e^{i\vec{Q}\cdot(\vec{r_n}-\vec{r_m})}$$
$$= Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_{n} \sum_{m \neq n} e^{i\vec{Q}\cdot(\vec{r_n}-\vec{r_m})}$$

The sum over  $m \neq m$  is now replaced with an integral of the continuous atomic pair density function,  $\rho_n(\vec{r}_{nm})$  and adding and subtracting the average atomic density  $\rho_{at}$ 

$$\begin{aligned} (\vec{Q}) &= Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \int_V \rho_n(\vec{r}_{nm}) e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} \, dV_m \\ &= Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \int_V [\rho_n(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} \, dV_m \\ &+ f(\vec{Q})^2 \rho_{at} \sum_n \int_V e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} \, dV_m \end{aligned}$$

Carlo Segre (Illinois Tech)

V

Consider a mono-atomic (-molecular) system where the total scattered intensity is given by

$$I(\vec{Q}) = f(\vec{Q})^2 \sum_{n} e^{i\vec{Q}\cdot\vec{r_n}} \sum_{m} e^{-i\vec{Q}\cdot\vec{r_m}} = f(\vec{Q})^2 \sum_{n} \sum_{m} e^{i\vec{Q}\cdot(\vec{r_n}-\vec{r_m})}$$
$$= Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_{n} \sum_{m \neq n} e^{i\vec{Q}\cdot(\vec{r_n}-\vec{r_m})}$$

The sum over  $m \neq m$  is now replaced with an integral of the continuous atomic pair density function,  $\rho_n(\vec{r}_{nm})$  and adding and subtracting the average atomic density  $\rho_{at}$ 

$$I(\vec{Q}) = Nf(\vec{Q})^{2} + f(\vec{Q})^{2} \sum_{n} \int_{V} \rho_{n}(\vec{r}_{nm}) e^{i\vec{Q}\cdot(\vec{r}_{n}-\vec{r}_{m})} dV_{m}$$
  
=  $Nf(\vec{Q})^{2} + f(\vec{Q})^{2} \sum_{n} \int_{V} [\rho_{n}(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q}\cdot(\vec{r}_{n}-\vec{r}_{m})} dV_{m}$   
+  $f(\vec{Q})^{2} \rho_{at} \sum_{n} \int_{V} e^{i\vec{Q}\cdot(\vec{r}_{n}-\vec{r}_{m})} dV_{m} = I^{SRO}(\vec{Q})$ 

Carlo Segre (Illinois Tech)

V

Consider a mono-atomic (-molecular) system where the total scattered intensity is given by

$$I(\vec{Q}) = f(\vec{Q})^2 \sum_{n} e^{i\vec{Q}\cdot\vec{r_n}} \sum_{m} e^{-i\vec{Q}\cdot\vec{r_m}} = f(\vec{Q})^2 \sum_{n} \sum_{m} e^{i\vec{Q}\cdot(\vec{r_n}-\vec{r_m})}$$
$$= Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_{n} \sum_{m \neq n} e^{i\vec{Q}\cdot(\vec{r_n}-\vec{r_m})}$$

The sum over  $m \neq m$  is now replaced with an integral of the continuous atomic pair density function,  $\rho_n(\vec{r}_{nm})$  and adding and subtracting the average atomic density  $\rho_{at}$ 

$$I(\vec{Q}) = Nf(\vec{Q})^{2} + f(\vec{Q})^{2} \sum_{n} \int_{V} \rho_{n}(\vec{r}_{nm}) e^{i\vec{Q}\cdot(\vec{r}_{n} - \vec{r}_{m})} dV_{m}$$
  
=  $Nf(\vec{Q})^{2} + f(\vec{Q})^{2} \sum_{n} \int_{V} [\rho_{n}(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q}\cdot(\vec{r}_{n} - \vec{r}_{m})} dV_{m}$   
+  $f(\vec{Q})^{2} \rho_{at} \sum_{n} \int_{V} e^{i\vec{Q}\cdot(\vec{r}_{n} - \vec{r}_{m})} dV_{m} = I^{SRO}(\vec{Q}) + I^{SAXS}(\vec{Q})$ 

Carlo Segre (Illinois Tech)



For the moment, let us ignore the SAXS term and focus on the short range order term.



For the moment, let us ignore the SAXS term and focus on the short range order term.

$$I^{SRO}(\vec{Q}) = Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_{n} \int_{V} [\rho_n(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} dV$$



For the moment, let us ignore the SAXS term and focus on the short range order term.

$$I^{SRO}(\vec{Q}) = Nf(\vec{Q})^{2} + f(\vec{Q})^{2} \sum_{n} \int_{V} [\rho_{n}(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q}\cdot(\vec{r}_{n} - \vec{r}_{m})} dV$$

When we average over all choices of origin in the liquid,  $\langle \rho_n(\vec{r}_{nm}) \rangle \rightarrow \rho(\vec{r})$  and the sum simplifies to N giving:



For the moment, let us ignore the SAXS term and focus on the short range order term.

$$I^{SRO}(\vec{Q}) = Nf(\vec{Q})^{2} + f(\vec{Q})^{2} \sum_{n} \int_{V} [\rho_{n}(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q}\cdot(\vec{r}_{n} - \vec{r}_{m})} dV$$

When we average over all choices of origin in the liquid,  $\langle \rho_n(\vec{r}_{nm}) \rangle \rightarrow \rho(\vec{r})$  and the sum simplifies to N giving:

$$I^{SRO}(ec{Q}) = Nf(ec{Q})^2 + Nf(ec{Q})^2 \int_V \left[
ho(ec{r}) - 
ho_{at}
ight] e^{iec{Q}\cdotec{r}} dV$$



For the moment, let us ignore the SAXS term and focus on the short range order term.

$$I^{SRO}(\vec{Q}) = Nf(\vec{Q})^{2} + f(\vec{Q})^{2} \sum_{n} \int_{V} [\rho_{n}(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q}\cdot(\vec{r}_{n} - \vec{r}_{m})} dV$$

When we average over all choices of origin in the liquid,  $\langle \rho_n(\vec{r}_{nm}) \rangle \rightarrow \rho(\vec{r})$  and the sum simplifies to N giving:

$$I^{SRO}(ec{Q}) = Nf(ec{Q})^2 + Nf(ec{Q})^2 \int_V \left[
ho(ec{r}) - 
ho_{at}
ight] e^{iec{Q}\cdotec{r}} dV$$

Performing an orientational average results in



For the moment, let us ignore the SAXS term and focus on the short range order term.

$$I^{SRO}(\vec{Q}) = Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \int_V [\rho_n(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} dV$$

When we average over all choices of origin in the liquid,  $\langle \rho_n(\vec{r}_{nm}) \rangle \rightarrow \rho(\vec{r})$  and the sum simplifies to N giving:

$$I^{SRO}(ec{Q}) = Nf(ec{Q})^2 + Nf(ec{Q})^2 \int_V \left[
ho(ec{r}) - 
ho_{at}
ight] e^{iec{Q}\cdotec{r}} dV$$

Performing an orientational average results in

$$I^{SRO}(\vec{Q}) = Nf(\vec{Q})^2 + Nf(\vec{Q})^2 \int_0^\infty 4\pi r^2 \left[\rho(r) - \rho_{at}\right] rac{\sin Qr}{Qr} dr$$

Carlo Segre (Illinois Tech)

# $\mathsf{S}(\mathsf{Q})$ - the liquid structure factor

V

A bit of rearrangement results in the expression for the liquid structure factor, S(Q).

A bit of rearrangement results in the expression for the liquid structure factor, S(Q).

$$S(Q) = \frac{I^{SRO}(\vec{Q})}{Nf(Q)^2} = 1 + \frac{4\pi}{Q} \int_0^\infty r \left[\rho(r) - \rho_{at}\right] \sin(Qr) dr$$





A bit of rearrangement results in the expression for the liquid structure factor, S(Q).

$$S(Q) = \frac{I^{SRO}(\vec{Q})}{Nf(Q)^2} = 1 + \frac{4\pi}{Q} \int_0^\infty r \left[\rho(r) - \rho_{at}\right] \sin(Qr) dr$$

When  $Q \to \infty$ , the short wavelength limit,  $1/Q \to 0$  eliminates all dependence on the interparticle correlations and  $S(Q) \to 1$ .







A bit of rearrangement results in the expression for the liquid structure factor, S(Q).

$$S(Q) = \frac{I^{SRO}(\bar{Q})}{Nf(Q)^2} = 1 + \frac{4\pi}{Q} \int_0^\infty r \left[\rho(r) - \rho_{at}\right] \sin(Qr) dr$$

When  $Q \to \infty$ , the short wavelength limit,  $1/Q \to 0$  eliminates all dependence on the interparticle correlations and  $S(Q) \to 1$ . When  $Q \rightarrow 0$ , i.e. the long wavelength limit, sin $(Qr)/Q \rightarrow r$  and S(Q) is dominated by the density fluctuations in the system



8/20

A bit of rearrangement results in the expression for the liquid structure factor, S(Q).

$$S(Q) = \frac{I^{SRO}(\vec{Q})}{Nf(Q)^2} = 1 + \frac{4\pi}{Q} \int_0^\infty r \left[\rho(r) - \rho_{at}\right] \sin(Qr) dr$$

When  $Q \to \infty$ , the short wavelength limit,  $1/Q \to 0$  eliminates all dependence on the interparticle correlations and  $S(Q) \to 1$ .

We can rewrite the structure factor equation

When  $Q \rightarrow 0$ , i.e. the long wavelength limit,  $\sin(Qr)/Q \rightarrow r$  and S(Q) is dominated by the density fluctuations in the system



A bit of rearrangement results in the expression for the liquid structure factor, S(Q).

$$S(Q) = \frac{I^{SRO}(\vec{Q})}{Nf(Q)^2} = 1 + \frac{4\pi}{Q} \int_0^\infty r \left[\rho(r) - \rho_{at}\right] \sin(Qr) dr$$

When  $Q \to \infty$ , the short wavelength limit,  $1/Q \to 0$  eliminates all dependence on the interparticle correlations and  $S(Q) \to 1$ .

We can rewrite the structure factor equation

$$Q\left[S(Q)-1\right] = \int_0^\infty 4\pi r \left[\rho(r)-\rho_{at}\right] \sin(Qr) dr$$

When  $Q \rightarrow 0$ , i.e. the long wavelength limit, sin $(Qr)/Q \rightarrow r$  and S(Q) is dominated by the density fluctuations in the system



A bit of rearrangement results in the expression for the liquid structure factor, S(Q).

$$S(Q) = \frac{I^{SRO}(\vec{Q})}{Nf(Q)^2} = 1 + \frac{4\pi}{Q} \int_0^\infty r \left[\rho(r) - \rho_{at}\right] \sin(Qr) dr$$

When  $Q \to \infty$ , the short wavelength limit,  $1/Q \to 0$  eliminates all dependence on the interparticle correlations and  $S(Q) \to 1$ . When  $Q \rightarrow 0$ , i.e. the long wavelength limit,  $\sin(Qr)/Q \rightarrow r$  and S(Q) is dominated by the density fluctuations in the system

We can rewrite the structure factor equation

$$Q\left[S(Q)-1\right] = \int_0^\infty 4\pi r \left[\rho(r)-\rho_{at}\right] \sin(Qr) dr = \int_0^\infty \mathcal{H}(r) \sin(Qr) dr$$



A bit of rearrangement results in the expression for the liquid structure factor, S(Q).

$$S(Q) = \frac{I^{SRO}(\vec{Q})}{Nf(Q)^2} = 1 + \frac{4\pi}{Q} \int_0^\infty r \left[\rho(r) - \rho_{at}\right] \sin(Qr) dr$$

When  $Q \to \infty$ , the short wavelength limit,  $1/Q \to 0$  eliminates all dependence on the interparticle correlations and  $S(Q) \to 1$ . When  $Q \rightarrow 0$ , i.e. the long wavelength limit,  $\sin(Qr)/Q \rightarrow r$  and S(Q) is dominated by the density fluctuations in the system

We can rewrite the structure factor equation

$$Q\left[S(Q)-1\right] = \int_0^\infty 4\pi r \left[\rho(r)-\rho_{at}\right] \sin(Qr) dr = \int_0^\infty \mathcal{H}(r) \sin(Qr) dr$$



A bit of rearrangement results in the expression for the liquid structure factor, S(Q).

$$S(Q) = \frac{I^{SRO}(\vec{Q})}{Nf(Q)^2} = 1 + \frac{4\pi}{Q} \int_0^\infty r \left[\rho(r) - \rho_{at}\right] \sin(Qr) dr$$

When  $Q \to \infty$ , the short wavelength limit,  $1/Q \to 0$  eliminates all dependence on the interparticle correlations and  $S(Q) \to 1$ . When  $Q \rightarrow 0$ , i.e. the long wavelength limit,  $\sin(Qr)/Q \rightarrow r$  and S(Q) is dominated by the density fluctuations in the system

We can rewrite the structure factor equation

$$Q\left[S(Q)-1\right] = \int_0^\infty 4\pi r \left[\rho(r)-\rho_{at}\right] \sin(Qr) dr = \int_0^\infty \mathcal{H}(r) \sin(Qr) dr$$

Which is the sine Fourier Transform of the deviation of the atomic density from its average,  $\mathcal{H}(r) = 4\pi r [g(r) - 1]$ 

Carlo Segre (Illinois Tech)



We can invert the Fourier Transform to obtain



We can invert the Fourier Transform to obtain

$$\mathcal{H}(r) = rac{2}{\pi} \int_0^\infty Q\left[S(Q) - 1
ight] \sin(Qr) dQ$$



9/20

We can invert the Fourier Transform to obtain

$$\mathcal{H}(r) = rac{2}{\pi} \int_0^\infty Q \left[ S(Q) - 1 
ight] \sin(Qr) dQ$$

and thus the radial distribution function can be obtained from the structure factor (an experimentally measureable quantity).



We can invert the Fourier Transform to obtain

$$\mathcal{H}(r) = rac{2}{\pi} \int_0^\infty Q \left[ S(Q) - 1 
ight] \sin(Qr) dQ$$

and thus the radial distribution function can be obtained from the structure factor (an experimentally measureable quantity).

$$g(r) = 1 + \frac{1}{2\pi^2 r \rho_{at}} \int_0^\infty Q\left[S(Q) - 1\right] \sin(Qr) dQ$$



We can invert the Fourier Transform to obtain

$$\mathcal{H}(r) = rac{2}{\pi} \int_0^\infty Q \left[ S(Q) - 1 
ight] \sin(Qr) dQ$$

and thus the radial distribution function can be obtained from the structure factor (an experimentally measureable quantity).

$$g(r) = 1 + rac{1}{2\pi^2 r 
ho_{at}} \int_0^\infty Q\left[S(Q) - 1
ight] \sin(Qr) dQ$$

This formalism holds for both non-crystalline solids and liquids, even though inelastic scattering dominates in the latter.

Carlo Segre (Illinois Tech)



9/20

We can invert the Fourier Transform to obtain

$$\mathcal{H}(r) = rac{2}{\pi} \int_0^\infty Q \left[ S(Q) - 1 
ight] \sin(Qr) dQ$$

and thus the radial distribution function can be obtained from the structure factor (an experimentally measureable quantity).

$$g(r) = 1 + rac{1}{2\pi^2 r 
ho_{at}} \int_0^\infty Q\left[S(Q) - 1
ight] \sin(Qr) dQ$$

This formalism holds for both non-crystalline solids and liquids, even though inelastic scattering dominates in the latter.

The relation between radial distribution function and structure factor can be extended to multi-component systems where  $g(r) \rightarrow g_{ij}(r)$  and  $S(Q) \rightarrow S_{ij}(Q)$ .

Carlo Segre (Illinois Tech)

Liquid Ni metal was suspended electrostatically and allowed to cool from its liquidus temperature of  $1450^{\circ}$ C.

"Difference in Icosahedral Short-Range Order in Early and Late Transition Metal Liquids," G.W. Lee et al. *Phys. Rev. Lett* **93**, 037802 (2004).

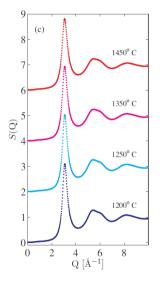






Liquid Ni metal was suspended electrostatically and allowed to cool from its liquidus temperature of  $1450^{\circ}$ C.

Measurement of the liquid structure factor shows supercooling with short range order.



"Difference in Icosahedral Short-Range Order in Early and Late Transition Metal Liquids," G.W. Lee et al. *Phys. Rev. Lett* **93**, 037802 (2004).

Carlo Segre (Illinois Tech)



Liquid Ni metal was suspended electrostatically and allowed to cool from its liquidus temperature of  $1450^{\circ}$ C.

Measurement of the liquid structure factor shows supercooling with short range order.

Integration of the radial distribution function indicates the presence of icosahedral clusters which inhibit crystallization.

0

8 (d)

"Difference in Icosahedral Short-Range Order in Early and Late Transition Metal Liquids," G.W. Lee et al. *Phys. Rev. Lett* **93**, 037802 (2004).

#### PHYS 570 - Fall 2024



10 / 20

Liquid Ni metal was suspended electrostatically and allowed to cool from its liquidus temperature of  $1450^{\circ}$ C.

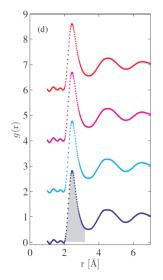
Measurement of the liquid structure factor shows supercooling with short range order.

Integration of the radial distribution function indicates the presence of icosahedral clusters which inhibit crystallization.



Details in the shape of the oscillations can be indicative of distortions in the icosahedra which depend on the metal species.

"Difference in Icosahedral Short-Range Order in Early and Late Transition Metal Liquids," G.W. Lee et al. *Phys. Rev. Lett* **93**, 037802 (2004).







Liquid scattering can be used to study dynamics

"Seeing real-space dynamics of liquid water through inelastic x-ray scattering," T. Iwashita et al. Sci. Adv. 3, e1603079 (2017).

Carlo Segre (Illinois Tech)



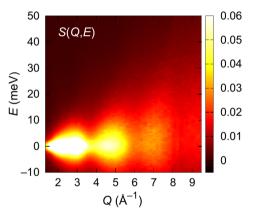
Liquid scattering can be used to study dynamics

In this article, the authors measured the liquid scattering as a function of both momentum, Q, and energy, E, transfer by using analyzers set for a specific energy (21.747 keV) but varying Q and then scanning the incident energy at fixed incident angle

"Seeing real-space dynamics of liquid water through inelastic x-ray scattering," T. Iwashita et al. Sci. Adv. 3, e1603079 (2017).

Liquid scattering can be used to study dynamics

In this article, the authors measured the liquid scattering as a function of both momentum, Q, and energy, E, transfer by using analyzers set for a specific energy (21.747 keV) but varying Q and then scanning the incident energy at fixed incident angle



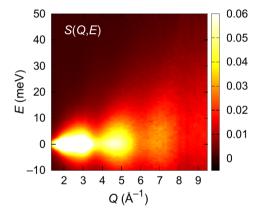
"Seeing real-space dynamics of liquid water through inelastic x-ray scattering," T. Iwashita et al. Sci. Adv. 3, e1603079 (2017).



Liquid scattering can be used to study dynamics

In this article, the authors measured the liquid scattering as a function of both momentum, Q, and energy, E, transfer by using analyzers set for a specific energy (21.747 keV) but varying Q and then scanning the incident energy at fixed incident angle

The Van Hoff function can be obtained by a double Fourier transform



"Seeing real-space dynamics of liquid water through inelastic x-ray scattering," T. Iwashita et al. Sci. Adv. 3, e1603079 (2017).

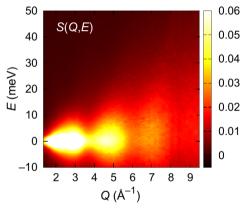




Liquid scattering can be used to study dynamics

In this article, the authors measured the liquid scattering as a function of both momentum, Q, and energy, E, transfer by using analyzers set for a specific energy (21.747 keV) but varying Q and then scanning the incident energy at fixed incident angle

The Van Hoff function can be obtained by a double Fourier transform

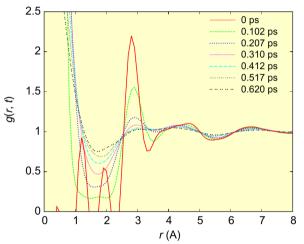


$$g(r,t)-1=rac{1}{2
ho\pi^2 r}\int\int e^{i\omega t}\sin(Qr)QS(Q,E)dQ\,dE$$

"Seeing real-space dynamics of liquid water through inelastic x-ray scattering," T. Iwashita et al. Sci. Adv. 3, e1603079 (2017).

Carlo Segre (Illinois Tech)



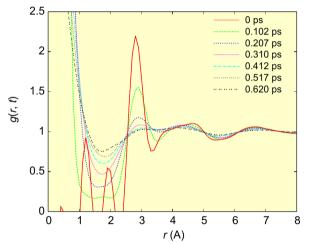


"Seeing real-space dynamics of liquid water through inelastic x-ray scattering," T. Iwashita et al. Sci. Adv. 3, e1603079 (2017).

Carlo Segre (Illinois Tech)







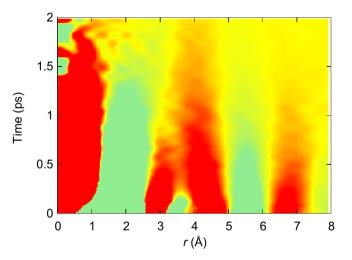
The peak below 1.5Å represents the self motion of the central atom while the data at longer distances represents the collective motions of two different atoms, in this case the oxygens

"Seeing real-space dynamics of liquid water through inelastic x-ray scattering," T. Iwashita et al. Sci. Adv. 3, e1603079 (2017).

Carlo Segre (Illinois Tech)

PHYS 570 - Fall 2024

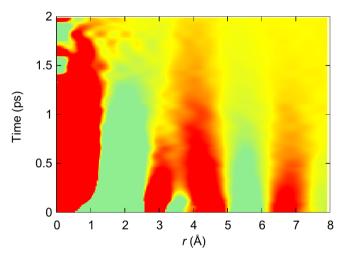
September 30, 2024 12 / 20



The peak below 1.5Å represents the self motion of the central atom while the data at longer distances represents the collective motions of two different atoms, in this case the oxygens

"Seeing real-space dynamics of liquid water through inelastic x-ray scattering," T. Iwashita et al. Sci. Adv. 3, e1603079 (2017).





The peak below 1.5Å represents the self motion of the central atom while the data at longer distances represents the collective motions of two different atoms, in this case the oxygens

The first and second peaks are highly coupled in space and time and merge within 0.8 ps. This behavior is different from liquid metals and leads to the viscosity of water.

"Seeing real-space dynamics of liquid water through inelastic x-ray scattering," T. Iwashita et al. Sci. Adv. 3, e1603079 (2017).







$$I(\vec{Q}) = Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \int_V [\rho_n(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} \, dV_m + f(\vec{Q})^2 \rho_{at} \sum_n \int_V e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} \, dV_m$$



$$I(\vec{Q}) = Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \int_V [\rho_n(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} \, dV_m + f(\vec{Q})^2 \rho_{at} \sum_n \int_V e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} \, dV_m$$



$$I(\vec{Q}) = Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \int_V [\rho_n(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} \, dV_m + f(\vec{Q})^2 \rho_{at} \sum_n \int_V e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} \, dV_m$$

$$I^{SAXS}(\vec{Q}) = f^2 \sum_{n} \int_{V} \rho_{at} e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} dV_m$$



$$I(\vec{Q}) = Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \int_V [\rho_n(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} \, dV_m + f(\vec{Q})^2 \rho_{at} \sum_n \int_V e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} \, dV_m$$

$$I^{SAXS}(\vec{Q}) = f^2 \sum_n \int_V \rho_{at} e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} dV_m = f^2 \sum_n e^{i\vec{Q}\cdot\vec{r}_n} \int_V \rho_{at} e^{-i\vec{Q}\cdot\vec{r}_m} dV_m$$



$$I(\vec{Q}) = Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \int_V [\rho_n(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} \, dV_m + f(\vec{Q})^2 \rho_{at} \sum_n \int_V e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} \, dV_m$$

Recall that there was an additional term in the scattering intensity which becomes important at small Q.

$$I^{SAXS}(\vec{Q}) = f^2 \sum_n \int_V \rho_{at} e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} dV_m = f^2 \sum_n e^{i\vec{Q}\cdot\vec{r}_n} \int_V \rho_{at} e^{-i\vec{Q}\cdot\vec{r}_m} dV_m$$
$$= f^2 \int_V \rho_{at} e^{i\vec{Q}\cdot\vec{r}_n} dV_n \int_V \rho_{at} e^{-i\vec{Q}\cdot\vec{r}_m} dV_m$$

Carlo Segre (Illinois Tech)

PHYS 570 - Fall 2024

September 30, 2024 13 / 20



$$I(\vec{Q}) = Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \int_V [\rho_n(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} \, dV_m + f(\vec{Q})^2 \rho_{at} \sum_n \int_V e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} \, dV_m$$

Recall that there was an additional term in the scattering intensity which becomes important at small Q.

$$I^{SAXS}(\vec{Q}) = f^2 \sum_{n} \int_{V} \rho_{at} e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} dV_m = f^2 \sum_{n} e^{i\vec{Q}\cdot\vec{r}_n} \int_{V} \rho_{at} e^{-i\vec{Q}\cdot\vec{r}_m} dV_m$$
$$= f^2 \int_{V} \rho_{at} e^{i\vec{Q}\cdot\vec{r}_n} dV_n \int_{V} \rho_{at} e^{-i\vec{Q}\cdot\vec{r}_m} dV_m = \left| \int_{V} \rho_{sl} e^{i\vec{Q}\cdot\vec{r}} dV \right|^2$$

Carlo Segre (Illinois Tech)



$$I(\vec{Q}) = Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \int_V [\rho_n(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} \, dV_m + f(\vec{Q})^2 \rho_{at} \sum_n \int_V e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} \, dV_m$$

Recall that there was an additional term in the scattering intensity which becomes important at small Q.

$$I^{SAXS}(\vec{Q}) = f^2 \sum_n \int_V \rho_{at} e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} dV_m = f^2 \sum_n e^{i\vec{Q}\cdot\vec{r}_n} \int_V \rho_{at} e^{-i\vec{Q}\cdot\vec{r}_m} dV_m$$
$$= f^2 \int_V \rho_{at} e^{i\vec{Q}\cdot\vec{r}_n} dV_n \int_V \rho_{at} e^{-i\vec{Q}\cdot\vec{r}_m} dV_m = \left| \int_V \rho_{sl} e^{i\vec{Q}\cdot\vec{r}} dV \right|^2$$

Where we have assumed sufficient averaging and introduced  $\rho_{sl} = f \rho_{at}$ .

Carlo Segre (Illinois Tech)



$$I(\vec{Q}) = Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \int_V [\rho_n(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} \, dV_m + f(\vec{Q})^2 \rho_{at} \sum_n \int_V e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} \, dV_m$$

Recall that there was an additional term in the scattering intensity which becomes important at small Q.

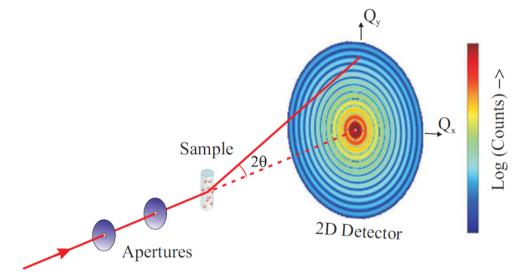
$$I^{SAXS}(\vec{Q}) = f^2 \sum_{n} \int_{V} \rho_{at} e^{i\vec{Q}\cdot(\vec{r}_n - \vec{r}_m)} dV_m = f^2 \sum_{n} e^{i\vec{Q}\cdot\vec{r}_n} \int_{V} \rho_{at} e^{-i\vec{Q}\cdot\vec{r}_m} dV_m$$
$$= f^2 \int_{V} \rho_{at} e^{i\vec{Q}\cdot\vec{r}_n} dV_n \int_{V} \rho_{at} e^{-i\vec{Q}\cdot\vec{r}_m} dV_m = \left| \int_{V} \rho_{sl} e^{i\vec{Q}\cdot\vec{r}} dV \right|^2$$

Where we have assumed sufficient averaging and introduced  $\rho_{sl} = f \rho_{at}$ . This final expression looks just like an atomic form factor but the charge density that we consider here is on a much longer length scale than an atom.

Carlo Segre (Illinois Tech)

# The SAXS experiment





V

15 / 20

The simplest case is for a dilute solution of non-interacting molecules.

$$I^{SAXS}(\vec{Q}) = \left| \int_{V_p} \rho_{sl} e^{i\vec{Q}\cdot\vec{r}} dV_p \right|^2$$



15 / 20

The simplest case is for a dilute solution of non-interacting molecules.

Assume that the scattering length density of each identical particle (molecule) is given by  $\rho_{sl,p}$  and the scattering length density of the solvent is  $\rho_{sl,0}$ .

$$I^{SAXS}(\vec{Q}) = \left| \int_{V_p} \rho_{sl} e^{i \vec{Q} \cdot \vec{r}} dV_p \right|^2$$

V

The simplest case is for a dilute solution of non-interacting molecules.

Assume that the scattering length density of each identical particle (molecule) is given by  $\rho_{sl,p}$  and the scattering length density of the solvent is  $\rho_{sl,0}$ .

$$I^{SAXS}(\vec{Q}) = \left| \int_{V_p} \rho_{sl} e^{i\vec{Q}\cdot\vec{r}} dV_p \right|^2 = (\rho_{sl,p} - \rho_{sl,0})^2 \left| \int_{V_p} e^{i\vec{Q}\cdot\vec{r}} dV_p \right|^2$$

The simplest case is for a dilute solution of non-interacting molecules.

Assume that the scattering length density of each identical particle (molecule) is given by  $\rho_{sl,p}$  and the scattering length density of the solvent is  $\rho_{sl,0}$ .

$$I^{SAXS}(\vec{Q}) = \left| \int_{V_{\rho}} \rho_{sl} e^{i\vec{Q}\cdot\vec{r}} dV_{\rho} \right|^2 = (\rho_{sl,\rho} - \rho_{sl,0})^2 \left| \int_{V_{\rho}} e^{i\vec{Q}\cdot\vec{r}} dV_{\rho} \right|^2$$

If we introduce the single-particle form factor  $\mathcal{F}(\vec{Q})$ :



The simplest case is for a dilute solution of non-interacting molecules.

Assume that the scattering length density of each identical particle (molecule) is given by  $\rho_{sl,p}$  and the scattering length density of the solvent is  $\rho_{sl,0}$ .

$$I^{SAXS}(\vec{Q}) = \left| \int_{V_{\rho}} \rho_{sl} e^{i\vec{Q}\cdot\vec{r}} dV_{\rho} \right|^2 = (\rho_{sl,\rho} - \rho_{sl,0})^2 \left| \int_{V_{\rho}} e^{i\vec{Q}\cdot\vec{r}} dV_{\rho} \right|^2$$

If we introduce the single-particle form factor  $\mathcal{F}(\vec{Q})$ :

$$\mathcal{F}(ec{Q}) = rac{1}{V_{
ho}}\int_{V_{
ho}}e^{iec{Q}\cdotec{r}}dV_{
ho}$$



The simplest case is for a dilute solution of non-interacting molecules.

Assume that the scattering length density of each identical particle (molecule) is given by  $\rho_{sl,p}$  and the scattering length density of the solvent is  $\rho_{sl,0}$ .

$$I^{SAXS}(\vec{Q}) = \left| \int_{V_{\rho}} \rho_{sl} e^{i\vec{Q}\cdot\vec{r}} dV_{\rho} \right|^2 = (\rho_{sl,\rho} - \rho_{sl,0})^2 \left| \int_{V_{\rho}} e^{i\vec{Q}\cdot\vec{r}} dV_{\rho} \right|^2$$

If we introduce the single-particle form factor  $\mathcal{F}(\vec{Q})$ :

$$\mathcal{F}(\vec{Q}) = \frac{1}{V_p} \int_{V_p} e^{i\vec{Q}\cdot\vec{r}} dV_p \qquad \qquad I^{SAXS}(\vec{Q}) = \Delta \rho^2 V_p^2 |\mathcal{F}(\vec{Q})|^2$$



15 / 20

The simplest case is for a dilute solution of non-interacting molecules.

Assume that the scattering length density of each identical particle (molecule) is given by  $\rho_{sl,p}$  and the scattering length density of the solvent is  $\rho_{sl,0}$ .

$$I^{SAXS}(\vec{Q}) = \left| \int_{V_{\rho}} \rho_{sl} e^{i\vec{Q}\cdot\vec{r}} dV_{\rho} \right|^2 = (\rho_{sl,\rho} - \rho_{sl,0})^2 \left| \int_{V_{\rho}} e^{i\vec{Q}\cdot\vec{r}} dV_{\rho} \right|^2$$

If we introduce the single-particle form factor  $\mathcal{F}(\vec{Q})$ :

$$\mathcal{F}(\vec{Q}) = \frac{1}{V_p} \int_{V_p} e^{i\vec{Q}\cdot\vec{r}} dV_p \qquad \qquad I^{SAXS}(\vec{Q}) = \Delta \rho^2 V_p^2 |\mathcal{F}(\vec{Q})|^2$$

Where  $\Delta \rho = (\rho_{sl,p} - \rho_{sl,0})$ , and the form factor depends on the morphology of the particle (size and shape).

Carlo Segre (Illinois Tech)



## Scattering from a sphere



There are only a few morphologies which can be computed exactly and the simplest is a constant density sphere of radius R.



There are only a few morphologies which can be computed exactly and the simplest is a constant density sphere of radius R.

$$\mathcal{F}(Q) = \frac{1}{V_{p}} \int_{0}^{R} \int_{0}^{2\pi} \int_{0}^{\pi} e^{iQr\cos\theta} r^{2}\sin\theta \,d\theta \,d\phi \,dr$$



16 / 20

There are only a few morphologies which can be computed exactly and the simplest is a constant density sphere of radius R.

$$\mathcal{F}(Q) = \frac{1}{V_p} \int_0^R \int_0^{2\pi} \int_0^{\pi} e^{iQr\cos\theta} r^2 \sin\theta \, d\theta \, d\phi \, dr$$
$$\mathcal{F}(Q) = \frac{1}{V_p} \int_0^R 4\pi \frac{\sin(Qr)}{Qr} r^2 \, dr$$



16 / 20

There are only a few morphologies which can be computed exactly and the simplest is a constant density sphere of radius R.

$$\mathcal{F}(Q) = \frac{1}{V_p} \int_0^R \int_0^{2\pi} \int_0^\pi e^{iQr\cos\theta} r^2 \sin\theta \, d\theta \, d\phi$$
$$\mathcal{F}(Q) = \frac{1}{V_p} \int_0^R 4\pi \frac{\sin(Qr)}{Qr} r^2 \, dr$$
$$= 3 \left[ \frac{\sin(QR) - QR\cos(QR)}{Q^3 R^3} \right]$$

dr



There are only a few morphologies which can be computed exactly and the simplest is a constant density sphere of radius R.

$$\mathcal{F}(Q) = \frac{1}{V_{p}} \int_{0}^{R} \int_{0}^{2\pi} \int_{0}^{\pi} e^{iQr\cos\theta} r^{2}\sin\theta \,d\theta \,d\phi \,dr$$
$$\mathcal{F}(Q) = \frac{1}{V_{p}} \int_{0}^{R} 4\pi \frac{\sin(Qr)}{Qr} r^{2} \,dr$$
$$= 3\left[\frac{\sin(QR) - QR\cos(QR)}{Q^{3}R^{3}}\right] \equiv \frac{3J_{1}(QR)}{QR}$$



There are only a few morphologies which can be computed exactly and the simplest is a constant density sphere of radius R.

$$\mathcal{F}(Q) = \frac{1}{V_{p}} \int_{0}^{R} \int_{0}^{2\pi} \int_{0}^{\pi} e^{iQr\cos\theta} r^{2}\sin\theta \,d\theta \,d\phi \,dr$$
$$\mathcal{F}(Q) = \frac{1}{V_{p}} \int_{0}^{R} 4\pi \frac{\sin(Qr)}{Qr} r^{2} \,dr$$
$$= 3\left[\frac{\sin(QR) - QR\cos(QR)}{Q^{3}R^{3}}\right] \equiv \frac{3J_{1}(QR)}{QR}$$

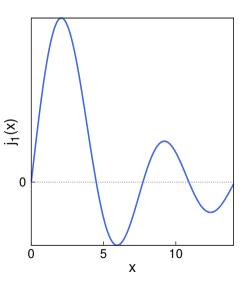
Where  $J_1(x)$  is the Bessel function of the first kind



There are only a few morphologies which can be computed exactly and the simplest is a constant density sphere of radius R.

$$\mathcal{F}(Q) = \frac{1}{V_p} \int_0^R \int_0^{2\pi} \int_0^{\pi} e^{iQr\cos\theta} r^2 \sin\theta \, d\theta \, d\phi \, dr$$
$$\mathcal{F}(Q) = \frac{1}{V_p} \int_0^R 4\pi \frac{\sin(Qr)}{Qr} r^2 \, dr$$
$$= 3 \left[ \frac{\sin(QR) - QR\cos(QR)}{Q^3 R^3} \right] \equiv \frac{3J_1(QR)}{QR}$$

Where  $J_1(x)$  is the Bessel function of the first kind



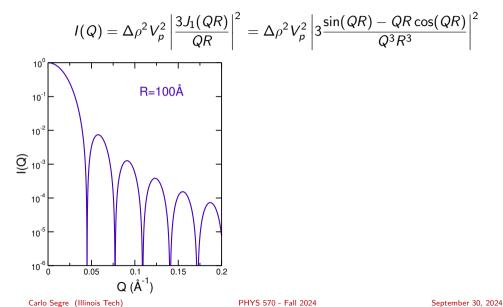


$$I(Q) = \Delta \rho^2 V_p^2 \left| \frac{3J_1(QR)}{QR} \right|^2$$



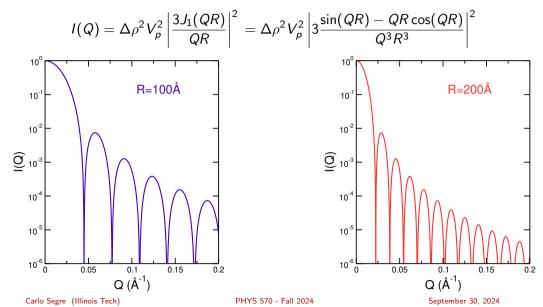
$$I(Q) = \Delta \rho^2 V_p^2 \left| \frac{3J_1(QR)}{QR} \right|^2 = \Delta \rho^2 V_p^2 \left| 3 \frac{\sin(QR) - QR\cos(QR)}{Q^3 R^3} \right|^2$$





17 / 20





17 / 20



$$I^{SAXS}(\vec{Q}) = \Delta \rho^2 V_p^2 |\mathcal{F}(\vec{Q})|^2,$$

Carlo Segre (Illinois Tech)



$$I^{SAXS}(\vec{Q}) = \Delta 
ho^2 V_p^2 |\mathcal{F}(\vec{Q})|^2,$$

$$\mathcal{F}(\vec{Q}) = 3\left[\frac{\sin(QR) - QR\cos(QR)}{Q^3R^3}\right]$$



$$I^{SAXS}(\vec{Q}) = \Delta \rho^2 V_{\rho}^2 |\mathcal{F}(\vec{Q})|^2, \qquad \qquad \mathcal{F}(\vec{Q}) = 3 \left[ rac{\sin(QR) - QR\cos(QR)}{Q^3 R^3} 
ight]$$

In the long wavelength limit QR 
ightarrow 0 we can approximate the scattering factor with the first terms of the sum



$$I^{SAXS}(\vec{Q}) = \Delta \rho^2 V_p^2 |\mathcal{F}(\vec{Q})|^2, \qquad \qquad \mathcal{F}(\vec{Q}) = 3 \left[ \frac{\sin(QR) - QR\cos(QR)}{Q^3 R^3} \right]$$

In the long wavelength limit QR 
ightarrow 0 we can approximate the scattering factor with the first terms of the sum

$$\mathcal{F}(Q) pprox rac{3}{Q^3 R^3} \left[ QR - rac{Q^3 R^3}{6} + rac{Q^5 R^5}{120} - \cdots - QR \left( 1 - rac{Q^2 R^2}{2} + rac{Q^4 R^4}{24} - \cdots 
ight) 
ight]$$



$$I^{SAXS}(\vec{Q}) = \Delta \rho^2 V_{\rho}^2 |\mathcal{F}(\vec{Q})|^2, \qquad \qquad \mathcal{F}(\vec{Q}) = 3 \left[ \frac{\sin(QR) - QR\cos(QR)}{Q^3 R^3} \right]$$

In the long wavelength limit QR 
ightarrow 0 we can approximate the scattering factor with the first terms of the sum

$$\mathcal{F}(Q) \approx \frac{3}{Q^3 R^3} \left[ QR - \frac{Q^3 R^3}{6} + \frac{Q^5 R^5}{120} - \dots - QR \left( 1 - \frac{Q^2 R^2}{2} + \frac{Q^4 R^4}{24} - \dots \right) \right]$$
  
this simplifies to  $\mathcal{F}(Q) \approx 1 - \frac{Q^2 R^2}{10}$  and

\_



$$I^{SAXS}(\vec{Q}) = \Delta \rho^2 V_{\rho}^2 |\mathcal{F}(\vec{Q})|^2, \qquad \qquad \mathcal{F}(\vec{Q}) = 3 \left[ \frac{\sin(QR) - QR\cos(QR)}{Q^3 R^3} \right]$$

In the long wavelength limit QR 
ightarrow 0 we can approximate the scattering factor with the first terms of the sum

$$\mathcal{F}(Q) \approx \frac{3}{Q^3 R^3} \left[ QR - \frac{Q^3 R^3}{6} + \frac{Q^5 R^5}{120} - \dots - QR \left( 1 - \frac{Q^2 R^2}{2} + \frac{Q^4 R^4}{24} - \dots \right) \right]$$
  
this simplifies to  $\mathcal{F}(Q) \approx 1 - \frac{Q^2 R^2}{10}$  and

$$I^{SAXS}(Q)pprox \Delta
ho^2 V_{
ho}^2 \left[1-rac{Q^2R^2}{10}
ight]^2 \, .$$

Carlo Segre (Illinois Tech)

\_



$$I^{SAXS}(\vec{Q}) = \Delta \rho^2 V_{\rho}^2 |\mathcal{F}(\vec{Q})|^2, \qquad \qquad \mathcal{F}(\vec{Q}) = 3 \left[ \frac{\sin(QR) - QR\cos(QR)}{Q^3 R^3} \right]$$

In the long wavelength limit QR 
ightarrow 0 we can approximate the scattering factor with the first terms of the sum

$$\mathcal{F}(Q) \approx \frac{3}{Q^3 R^3} \left[ QR - \frac{Q^3 R^3}{6} + \frac{Q^5 R^5}{120} - \dots - QR \left( 1 - \frac{Q^2 R^2}{2} + \frac{Q^4 R^4}{24} - \dots \right) \right]$$
  
this simplifies to  $\mathcal{F}(Q) \approx 1 - \frac{Q^2 R^2}{10}$  and

$$\mathcal{I}^{SAXS}(Q) pprox \Delta 
ho^2 V_
ho^2 \left[1 - rac{Q^2 R^2}{10}
ight]^2 pprox \Delta 
ho^2 V_
ho^2 \left[1 - rac{Q^2 R^2}{5}
ight]$$

Carlo Segre (Illinois Tech)



\_

$$I^{SAXS}(\vec{Q}) = \Delta \rho^2 V_{\rho}^2 |\mathcal{F}(\vec{Q})|^2, \qquad \qquad \mathcal{F}(\vec{Q}) = 3 \left[ \frac{\sin(QR) - QR\cos(QR)}{Q^3 R^3} \right]$$

In the long wavelength limit QR 
ightarrow 0 we can approximate the scattering factor with the first terms of the sum

$$\mathcal{F}(Q) \approx \frac{3}{Q^3 R^3} \left[ QR - \frac{Q^3 R^3}{6} + \frac{Q^5 R^5}{120} - \dots - QR \left( 1 - \frac{Q^2 R^2}{2} + \frac{Q^4 R^4}{24} - \dots \right) \right]$$
  
this simplifies to  $\mathcal{F}(Q) \approx 1 - \frac{Q^2 R^2}{10}$  and

$$I^{SAXS}(Q) \approx \Delta \rho^2 V_p^2 \left[ 1 - \frac{Q^2 R^2}{10} \right]^2 \approx \Delta \rho^2 V_p^2 \left[ 1 - \frac{Q^2 R^2}{5} \right] \approx \Delta \rho^2 V_p^2 e^{-Q^2 R^2/5}, \quad QR \ll 1$$

Carlo Segre (Illinois Tech)

PHYS 570 - Fall 2024



In the long wavelength limit ( $QR \rightarrow 0$ ), the form factor becomes



In the long wavelength limit ( $QR \rightarrow 0$ ), the form factor becomes

$$\mathcal{F}(Q)pprox 1-rac{Q^2R^2}{10}$$



In the long wavelength limit ( $QR \rightarrow 0$ ), the form factor becomes

$$egin{split} \mathcal{F}(Q) &pprox 1 - rac{Q^2 R^2}{10} \ I(Q) &pprox \Delta 
ho^2 V_p^2 e^{-Q^2 R^2/5} \end{split}$$



19 / 20

In the long wavelength limit ( $QR \rightarrow 0$ ), the form factor becomes

$$egin{aligned} \mathcal{F}(Q) &pprox 1 - rac{Q^2 R^2}{10} \ I(Q) &pprox \Delta 
ho^2 V_p^2 e^{-Q^2 R^2/5} \end{aligned}$$

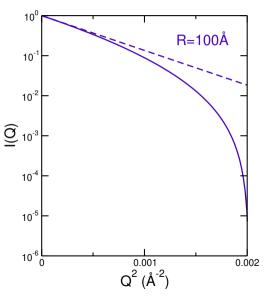
and the initial slope of the  $\log(I)$  vs  $Q^2$  plot is  $-R^2/5$ 



In the long wavelength limit (  $QR \rightarrow 0$  ), the form factor becomes

$$egin{split} \mathcal{F}(Q) &pprox 1 - rac{Q^2 R^2}{10} \ I(Q) &pprox \Delta 
ho^2 V_p^2 e^{-Q^2 R^2/5} \end{split}$$

and the initial slope of the  $\log(I)$  vs  $Q^2$  plot is  $-R^2/5$ 



Carlo Segre (Illinois Tech)

PHYS 570 - Fall 2024

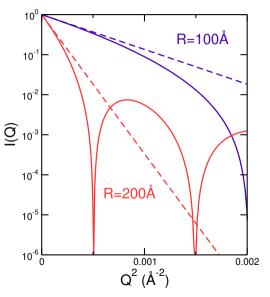
September 30, 2024 19 / 20



In the long wavelength limit ( QR 
ightarrow 0 ), the form factor becomes

$$\mathcal{F}(Q)pprox 1-rac{Q^2R^2}{10} \ I(Q)pprox \Delta
ho^2 V_
ho^2 e^{-Q^2R^2/5}$$

and the initial slope of the  $\log(I)$  vs  $Q^2$  plot is  $-R^2/5$ 



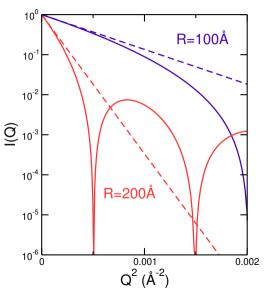


In the long wavelength limit ( QR 
ightarrow 0 ), the form factor becomes

$$egin{split} \mathcal{F}(Q) &pprox 1 - rac{Q^2 R^2}{10} \ I(Q) &pprox \Delta 
ho^2 V_
ho^2 e^{-Q^2 R^2/5} \end{split}$$

and the initial slope of the  $\log(I)$  vs  $Q^2$  plot is  $-R^2/5$ 

In terms of the radius of gyration,  $R_g$ , which for a sphere is given by  $\sqrt{\frac{3}{5}}R$ 





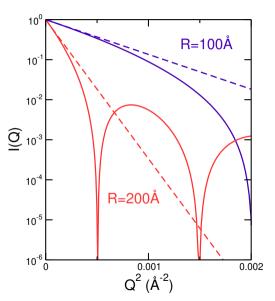
In the long wavelength limit ( QR 
ightarrow 0 ), the form factor becomes

$$egin{aligned} \mathcal{F}(Q) &pprox 1 - rac{Q^2 R^2}{10} \ & I(Q) &pprox \Delta 
ho^2 V_
ho^2 e^{-Q^2 R^2/5} \end{aligned}$$

and the initial slope of the  $\log(I)$  vs  $Q^2$  plot is  $-R^2/5$ 

In terms of the radius of gyration,  $R_g$ , which for a sphere is given by  $\sqrt{\frac{3}{5}}R$ 

$$I(Q) \approx \Delta \rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$



Carlo Segre (Illinois Tech)

PHYS 570 - Fall 2024

19 / 20



$$I(Q) \approx \Delta \rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$

Carlo Segre (Illinois Tech)

20 / 20



$$I(Q) \approx \Delta \rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$

The radius of gyration  $R_{\!g}$  is defined as the second moment of the volume occupied by the particle



$$I(Q) \approx \Delta \rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$

The radius of gyration  $R_g$  is defined as the second moment of the volume occupied by the particle

$$R_g^2 = \frac{1}{V_p} \int_{V_p} r^2 dV_p$$



$$I(Q) \approx \Delta \rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$

The radius of gyration  $R_{\!g}$  is defined as the second moment of the volume occupied by the particle

in terms of the scattering length density, it can be rewritten as

$$R_g^2 = \frac{1}{V_p} \int_{V_p} r^2 dV_p$$



$$I(Q) \approx \Delta \rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$

The radius of gyration  $R_{\!g}$  is defined as the second moment of the volume occupied by the particle

in terms of the scattering length density, it can be rewritten as

$$R_{g}^{2} = \frac{1}{V_{p}} \int_{V_{p}} r^{2} dV_{p} = \frac{\int_{V_{p}} \rho_{sl,p}(\vec{r}) r^{2} dV_{p}}{\int_{V_{p}} \rho_{sl,p}(\vec{r}) dV_{p}}$$



$$I(Q) \approx \Delta \rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$

The radius of gyration  $R_g$  is defined as the second moment of the volume occupied by the particle

in terms of the scattering length density, it can be rewritten as

$$R_{g}^{2} = \frac{1}{V_{p}} \int_{V_{p}} r^{2} dV_{p} = \frac{\int_{V_{p}} \rho_{sl,p}(\vec{r}) r^{2} dV_{p}}{\int_{V_{p}} \rho_{sl,p}(\vec{r}) dV_{p}}$$

after orientational averaging this expression can be used to extract  $R_g\,$  from experimental data using



20 / 20

$$I(Q) \approx \Delta \rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$

The radius of gyration  $R_g$  is defined as the second moment of the volume occupied by the particle

in terms of the scattering length density, it can be rewritten as

$$R_{g}^{2} = \frac{1}{V_{p}} \int_{V_{p}} r^{2} dV_{p} = \frac{\int_{V_{p}} \rho_{sl,p}(\vec{r}) r^{2} dV_{p}}{\int_{V_{p}} \rho_{sl,p}(\vec{r}) dV_{p}}$$

after orientational averaging this expression can be used to extract  $R_g\,$  from experimental data using

$$I_1^{SAXS}(Q) pprox \Delta 
ho^2 V_p^2 e^{-Q^2 R_g^2/3}$$

Carlo Segre (Illinois Tech)

PHYS 570 - Fall 2024



$$I(Q) \approx \Delta \rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$

The radius of gyration  $R_g$  is defined as the second moment of the volume occupied by the particle

in terms of the scattering length density, it can be rewritten as

$$R_{g}^{2} = \frac{1}{V_{p}} \int_{V_{p}} r^{2} dV_{p} = \frac{\int_{V_{p}} \rho_{sl,p}(\vec{r}) r^{2} dV_{p}}{\int_{V_{p}} \rho_{sl,p}(\vec{r}) dV_{p}}$$

after orientational averaging this expression can be used to extract  $R_g\,$  from experimental data using

$$I_1^{SAXS}(Q) \approx \Delta \rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$

this expression holds for uniform and non-uniform densities

Carlo Segre (Illinois Tech)

PHYS 570 - Fall 2024