



• HAXPES Experiments



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- X-ray magnetic circular dichroism



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- Resonant Scattering



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Reading Assignment: Chapter 8.4



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Homework Assignment #06: Chapter 6: 1,6,7,8,9 due Tuesday, November 16, 2021

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- X-ray magnetic circular dichroism
- Resonant Scattering

Reading Assignment: Chapter 8.4

Homework Assignment #06: Chapter 6: 1,6,7,8,9 due Tuesday, November 16, 2021 Homework Assignment #07: Chapter 7: 2,3,9,10,11 due Tuesday, November 30, 2021





In-class student presentations on research topics

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In-class student presentations on research topics

• Choose a research article which features a synchrotron technique



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- Get it approved by instructor first!



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- Schedule a 20 minute time on Final Exam Day (Tuesday, Dec 5, 2021, 09:00-19:00)



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Final project - writing a General User Proposal



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- Must be different techique than your presentation!

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V

HAXPES is used to probe the thickness of a $CoFe_2O_4/La_{0.66}Sr_{0.34}MnO_3$ heterostructure by varying both angle of incidence and photon energy



The thickness of the CoFe $_2O_4$ overlayer measured as 6.5 \pm 0.5 nm by TEM was probed in two ways:

B. Pal, S. Mukherjee, and D.D. Sarma, "Probing complex heterostructures using hard x-ray photoelectron spectroscopy (HAXPES)," J. Electron Spect. Related Phenomena 200, 332-339 (2015).

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Both results give consistent results with proper normalization and also show the uniformity of the $CoFe_2O_4$ overlayer

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Energy dispersive measurements can provide depth profiling of spherical nanoparticles



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By fitting the S 2p and Se 3p photoemission line the structure is revealed to be CdSe at the core and ZnCdS in the outer shell

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The variation in intensity of the Se/S lines and the Zn/Cd lines suggest that the Se is primarily located in the 2 nm core of the 5 nm particles with the Cd

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Si nanoparticle anodes suffer from the accumulation of the SEI layer which reduces performance. The SEI is formed by electrochemical decomposition of the electrolyte at the anode surface.

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The SEI from three different electrolyte combinations were studied: ethylene carbonate (EC), fluoroethylene carbonate (FEC), and a combination. The first of which gives poorer capacity and cycling stability.

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HAXPES is used to determine the elemental distribution and compounds present as a function of depth in the cycled Si anode.

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By varying the incident photon energy, it is possible to probe the SEI as a function of depth.

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PHYS 570 - Fall 2021

November 09, 2021





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- pure FEC shows less change with cycling than EC containing electrolytes

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Using HAXPES data from Si, C, and F, a picture of SEI evolution dependence on electrolyte emerges

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as SEI grows, there is growth of Li_xSiO_y underneath as product of lithiation/delithiation

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EC – SEI contains LEDC-rich SEI which decomposes but continues be deposited with cycling

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FEC – SEI is mostly poly-FEC with LiF and $LiCO_3$ which remains stable with cycling

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EC – SEI contains LEDC-rich SEI which decomposes but continues be deposited with cycling

FEC – SEI is mostly poly-FEC with LiF and $LiCO_3$ which remains stable with cycling

The FEC acts to stabilize the SEI composition and prevent the change with depth that occurs with EC.

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V

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X-ray magnetic circular dichroism

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 $\Delta m = +1$ for "right-handed"

 $\Delta m = -1$ for "left-handed"

this measurement is sensitive to the internal/external magnetic fields which split the levels according to the Zeeman effect





 $m_{=+1}$

XMCD and electron sum rules

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The XMCD experiment requires a source capable of switching the polarization (quarter wave plate) or a sample whose magnetic splittings can be inverted by flipping an external magnetic field

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The absorption coefficient is first measured for both relative orientations of magnetic splitting and circular polarization

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field

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$$m_{orb} = -\frac{4q(10-n_{3d})}{r}$$

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$$p=\int_{L_3}(\mu^+-\mu^-)d{\cal E}$$



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 $\mu^+(\mathcal{E}) = \frac{1}{x} \ln \left(\frac{I_0^+}{I_c^+} \right)$ $\mu^{-}(\mathcal{E}) = \frac{1}{x} \ln \left(\frac{I_0^{-}}{I_0^{-}} \right)$

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 $q = \int_{L_3+L_2} (\mu^+ - \mu^-) d\mathcal{E}$

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$$\int_{L_3} (\mu^+ - \mu^-) d\mathcal{E}$$

= $\int_{L_3+L_2} (\mu^+ + \mu^-) d\mathcal{E}$

November 09, 2021







The Zintl compounds exhibit interesting magnetic properties including colossal magnetoresistance which can be of value for spintronics applications



"XMCD Characterization of the Ferromagnetic State of Yb₁₄MnSb₁₁," A.P. Holm, S.M. Kauzlarich, S.A. Morton, G.D. Waddill, W.E. Pickett, and J.G. Tobin, *J. Am. Chem. Soc.* **124**, 9894-9898 (2002).

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The Zintl compound, Yb_{14}MnSb_{11} is ferromagnetic below below 56K with a moment of $\sim 4\mu_B/{\rm formula}$ unit

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XMCD on a single crystal of $Yb_{14}MnSb_{11}$ can be used to understand the origin of the ferromagnetic moment

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The Sb edges show a tiny asymmetry that is in opposite sign compared to the Mn edges

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The Yb XMCD shows no asymmetry due to polarization

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Mn provides the bulk of the magnetic moment and appears to be in the divalent state. Sb provides a small antiferromagnetic component to the overall magnetic moment

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This will produce the resonant scattering term but not the XANES and EXAFS, which are purely quantum effects.

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Consider an electron under the influence of an oscillating electric field $\vec{E}_{in} = \hat{x}E_0e^{-i\omega t}$.

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The amplitude of the response has a resonance and dissipation

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The scattering factor can be rewritten

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$$f_s = \frac{\omega^2 + (-\omega_s^2 + i\omega\Gamma) - (-\omega_s^2 + i\omega\Gamma)}{(\omega^2 - \omega_s^2 + i\omega\Gamma)}$$



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The scattering factor can be rewritten

and since $\Gamma \ll \omega_s$

the second term being the dispersion correction whose real and imaginary components can be extracted

$$\begin{split} f_{s} &= \frac{\omega^{2} + \left(-\omega_{s}^{2} + i\omega\Gamma\right) - \left(-\omega_{s}^{2} + i\omega\Gamma\right)}{\left(\omega^{2} - \omega_{s}^{2} + i\omega\Gamma\right)} \\ &= 1 + \frac{\omega_{s}^{2} - i\omega\Gamma}{\left(\omega^{2} - \omega_{s}^{2} + i\omega\Gamma\right)} \approx 1 + \frac{\omega_{s}^{2}}{\left(\omega^{2} - \omega_{s}^{2} + i\omega\Gamma\right)} \end{split}$$

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Single oscillator dispersion terms

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Total cross-section



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this shows a frequency dependence with a peak at $\omega\approx\omega_s$

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when $\omega \gg \omega_s$, $\sigma_T \rightarrow \sigma_{free}$

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The Kramers-Kronig relations are derived using Cauchy's theorem to integrate a function with a pole



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The Kramers-Kronig relations can be rewritten by multiplying top and bottom by $(\omega' + \omega)$

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$$= \frac{1}{\pi} \mathcal{P} \int_{-\infty}^{+\infty} \frac{\omega' f''(\omega')}{(\omega'^2 - \omega^2)} d\omega' + \underbrace{\frac{\omega}{\pi} \mathcal{P} \int_{-\infty}^{+\infty} \frac{f''(\omega')}{(\omega'^2 - \omega^2)} d\omega'}_{-\infty} = \frac{2}{\pi} \mathcal{P} \int_{0}^{+\infty} \frac{\omega' f''(\omega')}{(\omega'^2 - \omega^2)} d\omega'$$

$$f''(\omega) = -\frac{1}{\pi} \mathcal{P} \int_{-\infty}^{+\infty} \frac{f'(\omega')}{(\omega' - \omega)} d\omega' = -\frac{1}{\pi} \mathcal{P} \int_{-\infty}^{+\infty} \frac{f'(\omega')}{(\omega' - \omega)} \frac{(\omega' + \omega)}{(\omega' + \omega)} d\omega'$$
$$= \frac{1}{\pi} \mathcal{P} \int_{-\infty}^{+\infty} \frac{\omega' f'(\omega')}{(\omega'^2 - \omega^2)} d\omega' - \frac{\omega}{\pi} \mathcal{P} \int_{-\infty}^{+\infty} \frac{f'(\omega')}{(\omega'^2 - \omega^2)} d\omega'$$

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Starting with the Kramers-Kronig relation for f^\prime

$$f'(\omega) = rac{2}{\pi} \mathcal{P} \int_0^\infty rac{\omega' f''(\omega')}{(\omega'^2 - \omega^2)} d\omega'$$



Starting with the Kramers-Kronig relation for f' and recalling that f'' is directly related to the absorption cross-section, σ_a

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$$\sigma_{a}(x) = \begin{cases} 0 & x \leq 1 \\ \sigma_{a}(1)x^{-3} & x \geq 1 \end{cases}$$





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if ω_K is the frequency of the K absorption edge and a substitution is made with $x = \omega'/\omega_K$ and $x_K = \omega/\omega_K$, f' becomes

assuming that σ_{s} is zero below ω_{K} and varies as ω'^{-3} above we have

$$\sigma_a(x) = \begin{cases} 0 & x \le 1 \\ \sigma_a(1)x^{-3} & x \ge 1 \end{cases}$$

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this can be evaluated for two $1 \ensuremath{\textit{s}}$ electrons

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$$f'(\omega) = -\frac{2\omega_K \sigma_a(1)}{4\pi r_0 c} \mathcal{P} \int_1^\infty \frac{1}{x(x^2 - x_K^2)} dx, \quad x_K = \frac{\omega}{\omega_K}$$



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For 2 1s electrons

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For 2 1s electrons and doing the integral gives

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at high energies $(x_K \to \infty)$ this dispersion correction vanishes as expected and at low energies $(x_K, q \to 0)$ the correction is -1.565, thereby partially quenching the scattering from the two 1s electrons


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More accurate calculations of the resonant corrections to the scattering factor can be made using a full quantum mechanical treatment

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The simple model, however, reproduces the main features of the Ar K-edge $% \left({{{\rm{T}}_{{\rm{T}}}}_{{\rm{T}}}} \right)$



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What is lacking, even in the more sophisticated calcuations, are the resonances near the absorption edges due to XANES, EXAFS and other localized resonance phenomena



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