

Today's outline - October 07, 2021





- Equivalence of Laue & Bragg conditions

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- Crystal structure factor

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- Lattices & space groups

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- Ewald sphere

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- XRayView demonstration

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Reading Assignment: Chapter 5.4

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Reading Assignment: Chapter 5.4

Homework Assignment #04:

Chapter 4: 2,4,6,7,10

due Tuesday, October 19, 2021



- Equivalence of Laue & Bragg conditions
- Crystal structure factor
- Lattices & space groups
- Ewald sphere
- XRayView demonstration

Reading Assignment: Chapter 5.4

Homework Assignment #04:

Chapter 4: 2,4,6,7,10

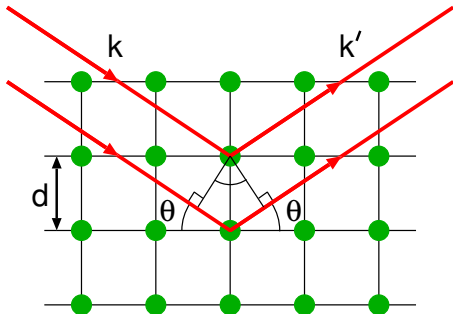
due Tuesday, October 19, 2021

Homework Assignment #05:

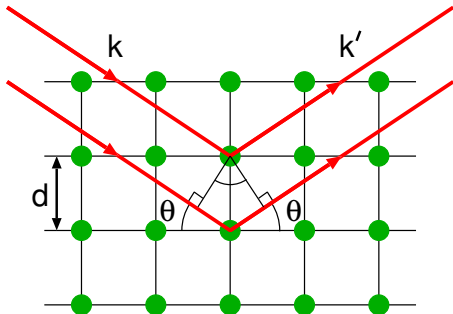
Chapter 5: 1,3,7,9,10

due Tuesday, November 02, 2021

Bragg condition

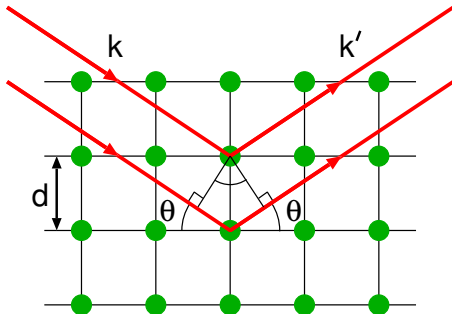


Bragg condition



The Bragg condition for diffraction is derived by assuming specular reflection from parallel planes separated by a distance d .

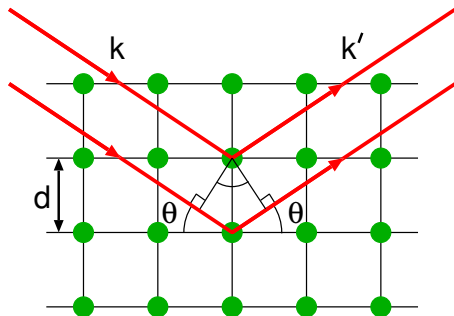
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Bragg condition

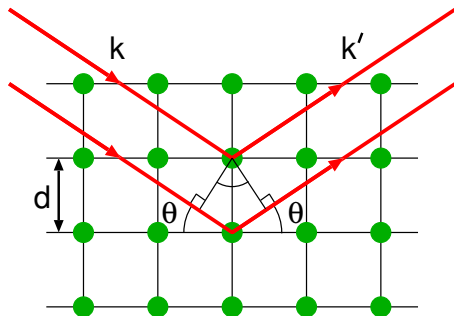


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If there is to be constructive interference, this additional distance must correspond to an integer number of wavelengths and we get the Bragg condition

Bragg condition



$$2d \sin \theta = \lambda$$

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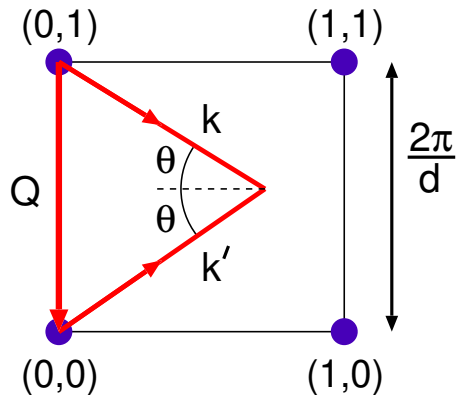
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Laue condition



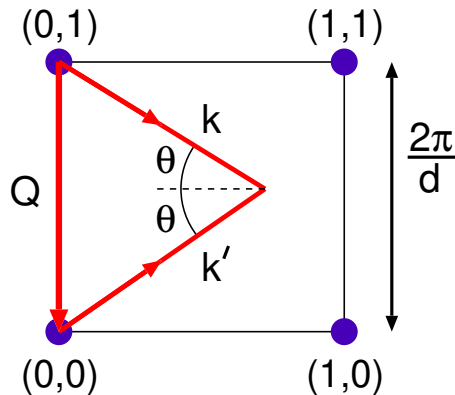
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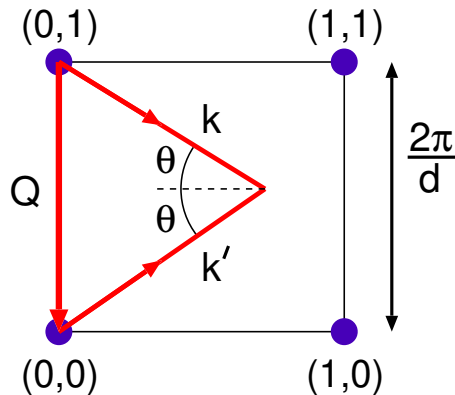


Laue condition



The Laue condition states that the scattering vector must be equal to a reciprocal lattice vector

$$\vec{Q} = \vec{G}_{hk}$$



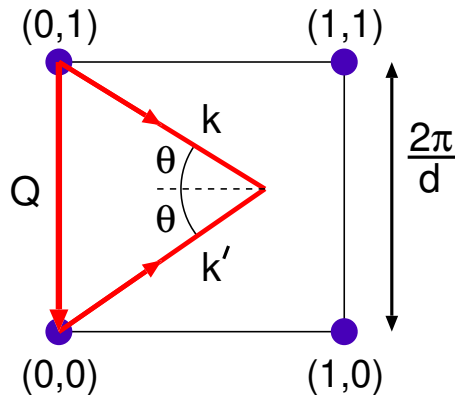
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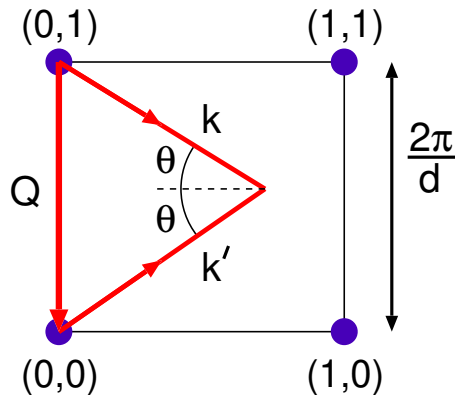
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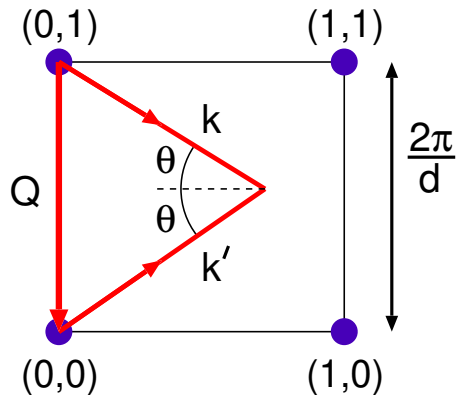


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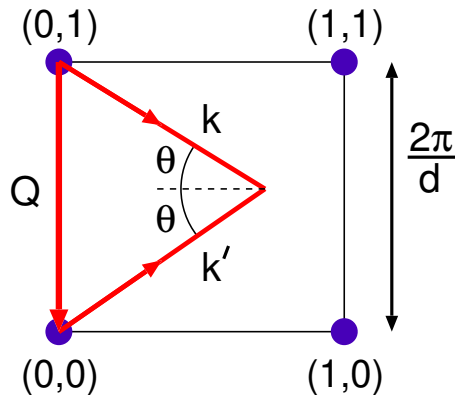


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Laue condition



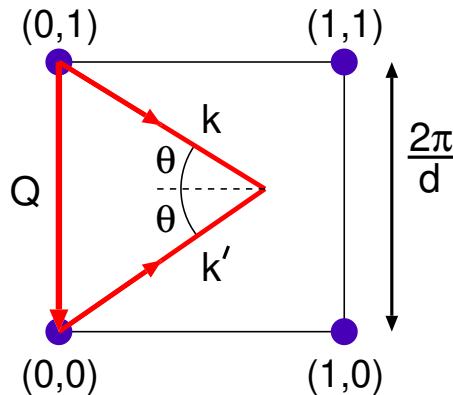
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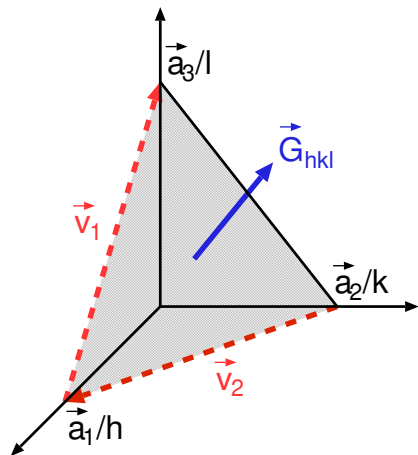
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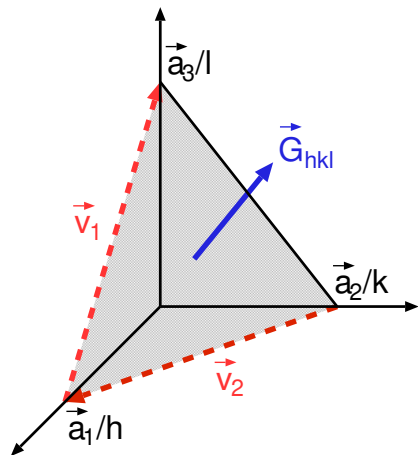
Thus the Bragg and Laue conditions are equivalent



General proof of Bragg-Laue equivalence

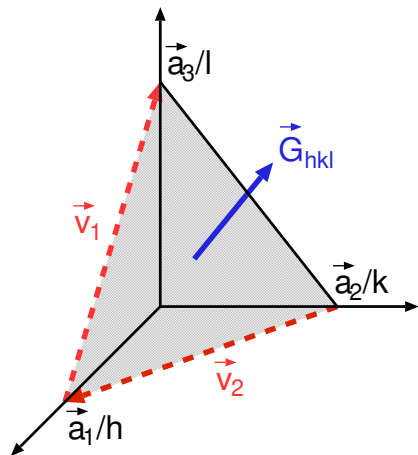


General proof of Bragg-Laue equivalence



Must show that for each point in reciprocal space, there exists a set of planes in the real space lattice such that:

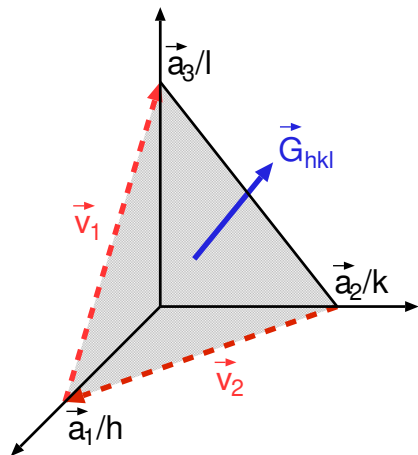
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Must show that for each point in reciprocal space, there exists a set of planes in the real space lattice such that:

\vec{G}_{hkl} is perpendicular to the planes with Miller indices (hkl)

General proof of Bragg-Laue equivalence

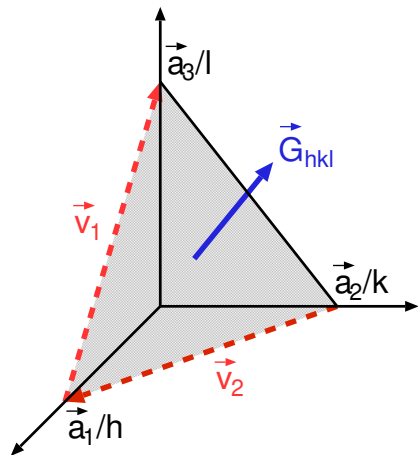


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$$|\vec{G}_{hkl}| = \frac{2\pi}{d_{hkl}}$$

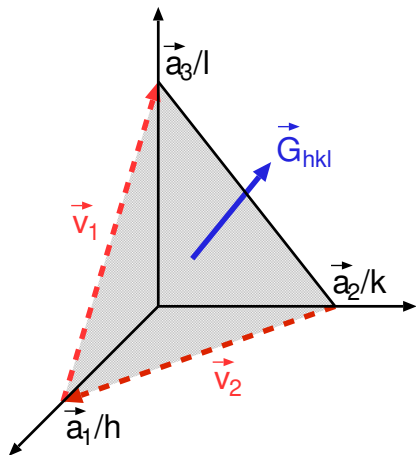
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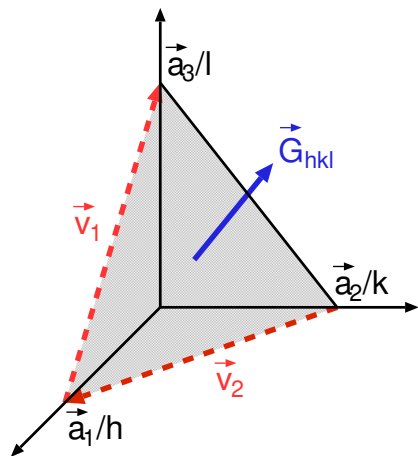
General proof of Bragg-Laue equivalence



The plane with Miller indices (hkl) intersects the three basis vectors of the lattice at a_1/h , a_2/k , and a_3/l



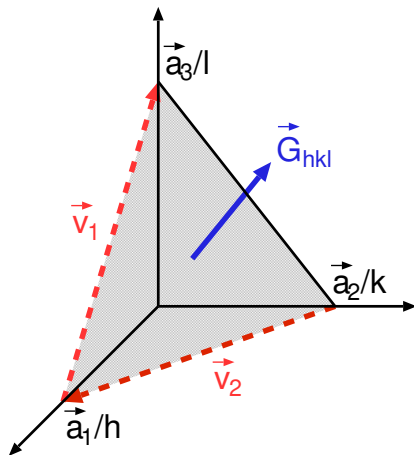
General proof of Bragg-Laue equivalence



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General proof of Bragg-Laue equivalence

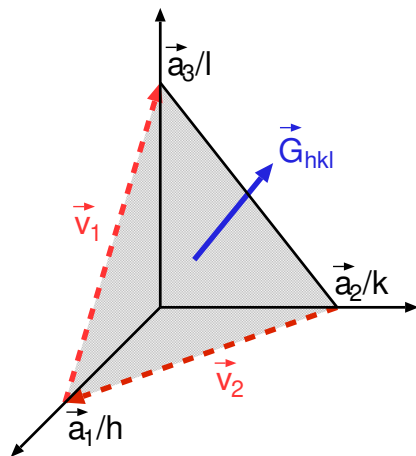


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General proof of Bragg-Laue equivalence

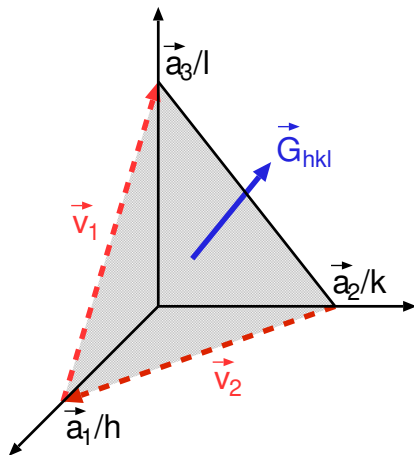


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General proof of Bragg-Laue equivalence



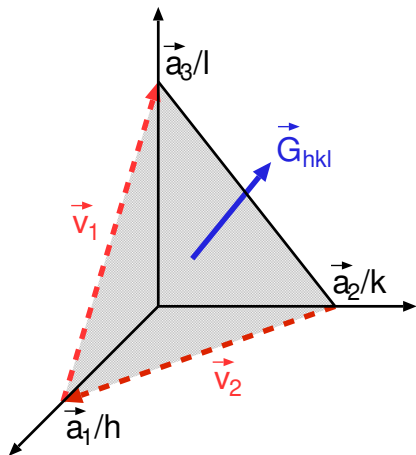
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$$\vec{v} = \epsilon_1 \vec{v}_1 + \epsilon_2 \vec{v}_2$$

General proof of Bragg-Laue equivalence



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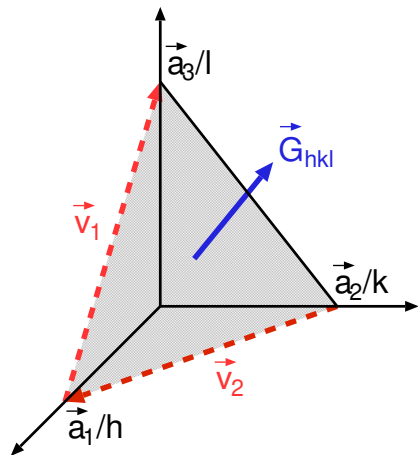
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$$\vec{G}_{hkl} \cdot \vec{v} = (h\vec{a}_1^* + k\vec{a}_2^* + l\vec{a}_3^*) \cdot \left((\epsilon_2 - \epsilon_1) \frac{\vec{a}_1}{h} - \epsilon_2 \frac{\vec{a}_2}{k} + \epsilon_1 \frac{\vec{a}_3}{l} \right)$$

General proof of Bragg-Laue equivalence



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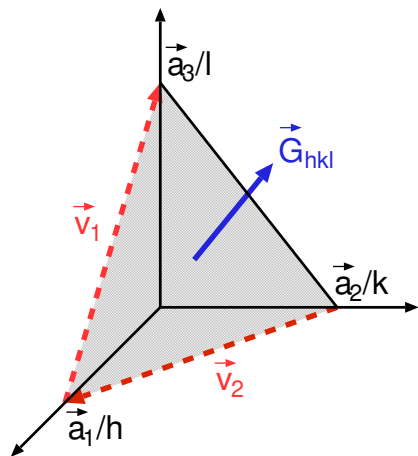
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General proof of Bragg-Laue equivalence



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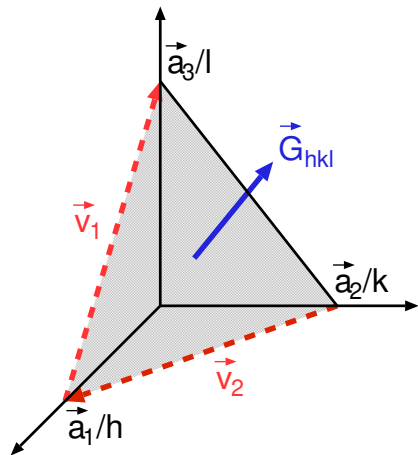
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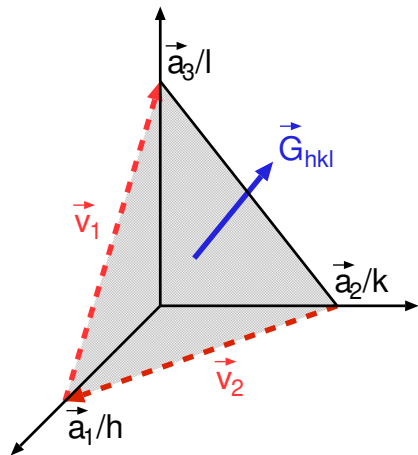
$$\begin{aligned} \vec{G}_{hkl} \cdot \vec{v} &= (h\vec{a}_1^* + k\vec{a}_2^* + l\vec{a}_3^*) \cdot \left((\epsilon_2 - \epsilon_1) \frac{\vec{a}_1}{h} - \epsilon_2 \frac{\vec{a}_2}{k} + \epsilon_1 \frac{\vec{a}_3}{l} \right) \\ &= 2\pi(\epsilon_2 - \epsilon_1 - \epsilon_2 + \epsilon_1) = 0 \end{aligned}$$

Thus \vec{G}_{hkl} is indeed normal to the plane with Miller indices (hkl)

General proof of Bragg-Laue equivalence

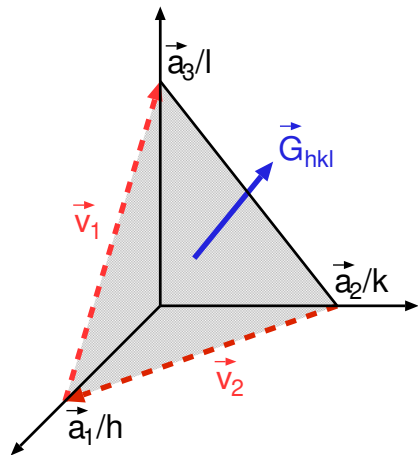


General proof of Bragg-Laue equivalence



The spacing between planes (hkl) is simply given by the distance from the origin to the plane along a normal vector

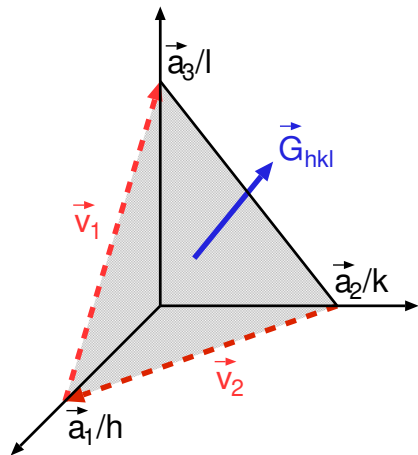
General proof of Bragg-Laue equivalence



The spacing between planes (hkl) is simply given by the distance from the origin to the plane along a normal vector

This can be computed as the projection of any vector which connects the origin to the plane onto the unit vector in the \vec{G}_{hkl} direction. In this case, we choose, \vec{a}_1/h

General proof of Bragg-Laue equivalence

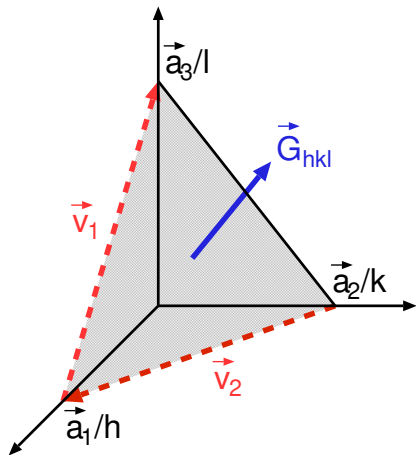


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General proof of Bragg-Laue equivalence



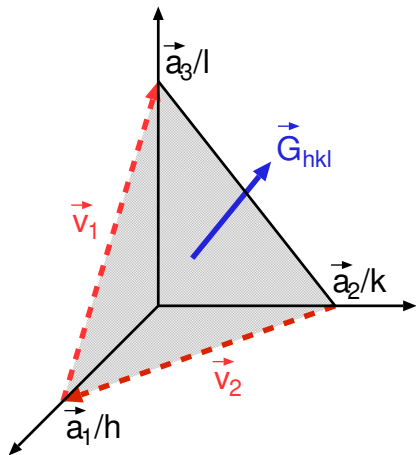
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$$\hat{G}_{hkl} \cdot \frac{\vec{a}_1}{h} = \frac{(h\vec{a}_1^* + k\vec{a}_2^* + l\vec{a}_3^*)}{|\vec{G}_{hkl}|} \cdot \frac{\vec{a}_1}{h}$$

General proof of Bragg-Laue equivalence



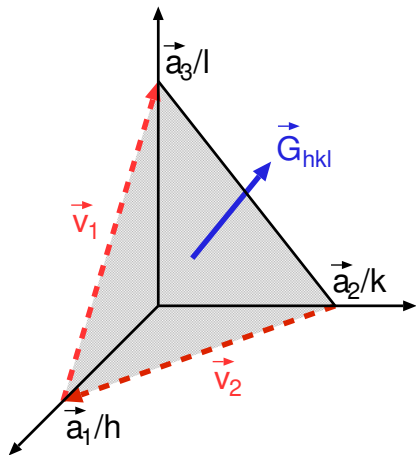
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General proof of Bragg-Laue equivalence



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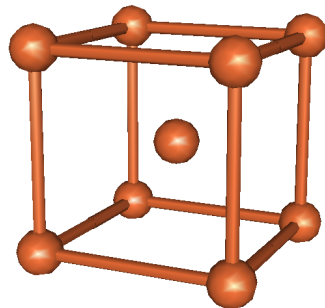
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BCC structure factor



In the body-centered cubic structure, there are 2 atoms in the conventional, cubic unit cell. These are located at

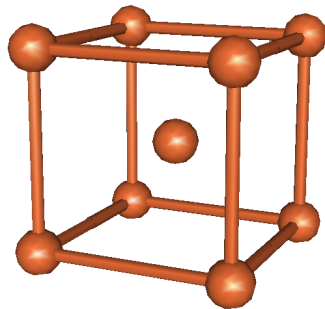


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$$\vec{r}_1 = 0, \quad \vec{r}_2 = \frac{1}{2}(\vec{a}_1 + \vec{a}_2 + \vec{a}_3)$$



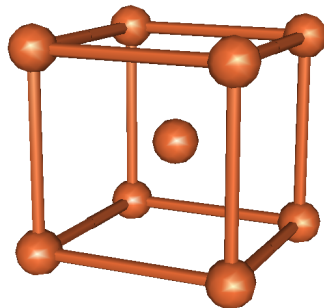
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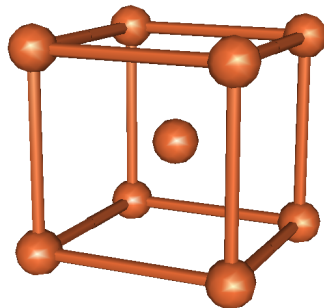


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$$F_{hkl}^{bcc} = f(\vec{G}) \sum_j e^{i\vec{G} \cdot \vec{r}_j}$$



BCC structure factor

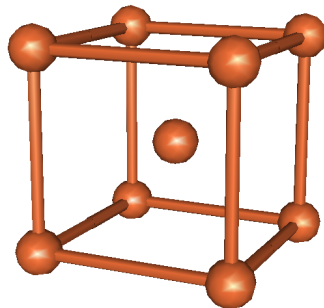


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$$\begin{aligned} F_{hkl}^{bcc} &= f(\vec{G}) \sum_j e^{i\vec{G} \cdot \vec{r}_j} \\ &= f(\vec{G}) \left(1 + e^{i\pi(h+k+l)} \right) \end{aligned}$$



BCC structure factor

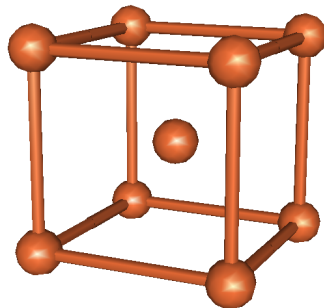


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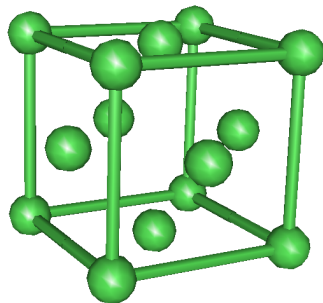
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FCC structure factor



In the face-centered cubic structure, there are 4 atoms in the conventional, cubic unit cell. These are located at

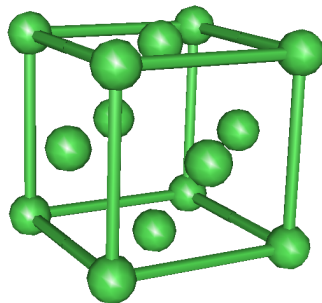


FCC structure factor



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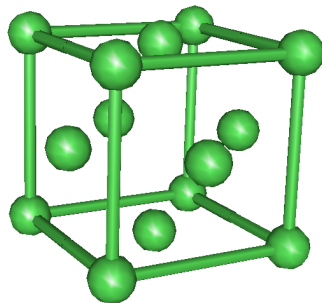
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the unit cell structure factor is thus



FCC structure factor

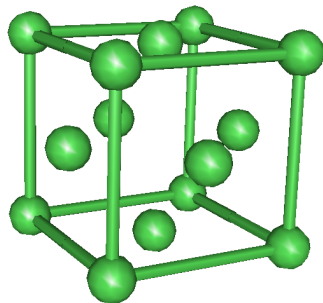


In the face-centered cubic structure, there are 4 atoms in the conventional, cubic unit cell. These are located at

$$\vec{r}_1 = 0, \quad \vec{r}_2 = \frac{1}{2}(\vec{a}_1 + \vec{a}_2), \quad \vec{r}_3 = \frac{1}{2}(\vec{a}_2 + \vec{a}_3), \quad \vec{r}_4 = \frac{1}{2}(\vec{a}_1 + \vec{a}_3)$$

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$$F_{hkl}^{fcc} = f(\vec{G}) \sum_j e^{i\vec{G} \cdot \vec{r}_j}$$



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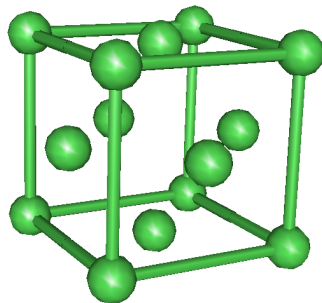


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$$\begin{aligned} F_{hkl}^{fcc} &= f(\vec{G}) \sum_j e^{i\vec{G} \cdot \vec{r}_j} \\ &= f(\vec{G}) \left(1 + e^{i\pi(h+k)} + e^{i\pi(k+l)} + e^{i\pi(h+l)} \right) \end{aligned}$$



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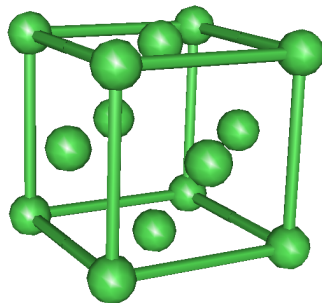


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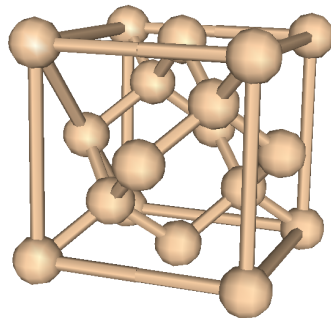
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Diamond structure



This is a face centered cubic structure with two atoms in the basis which leads to 8 atoms in the conventional unit cell. These are located at

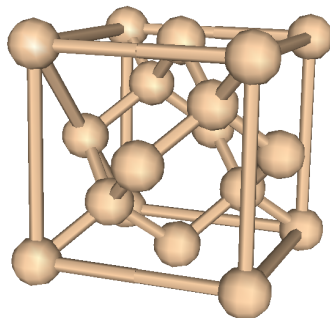


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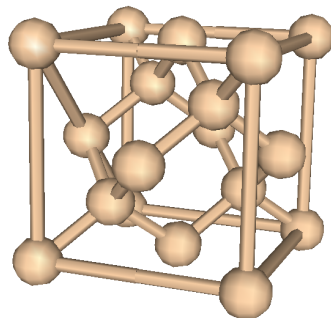
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Diamond structure

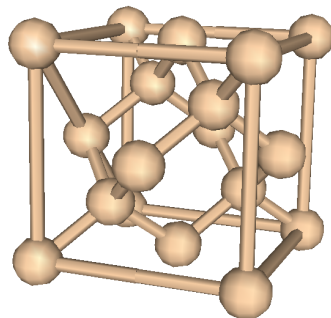


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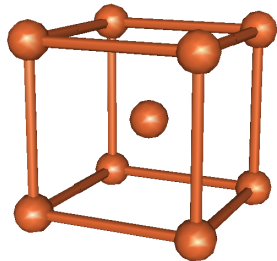
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This is non-zero when h, k, l all even and $h + k + l = 4n$ or h, k, l all odd

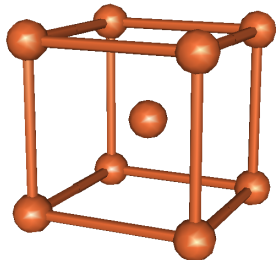


Heteroatomic structures

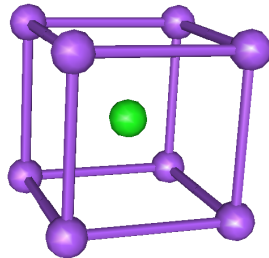


← bcc

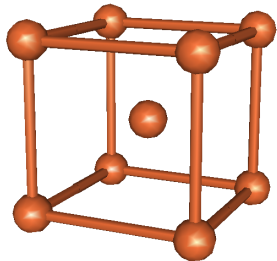
Heteroatomic structures



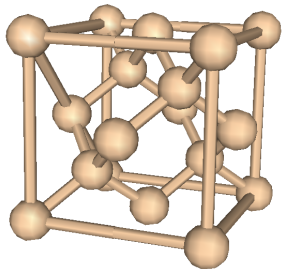
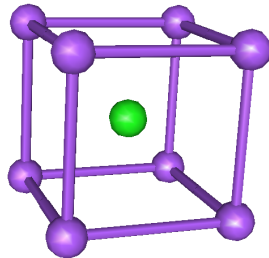
← bcc
sc →



Heteroatomic structures

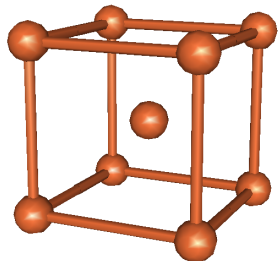


← bcc
sc →

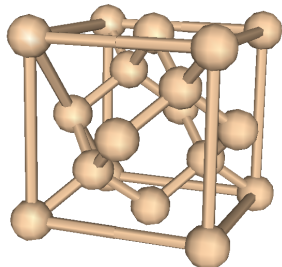
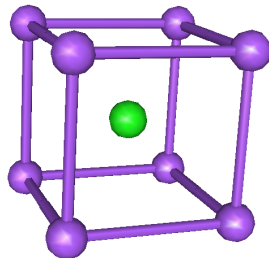


← diamond

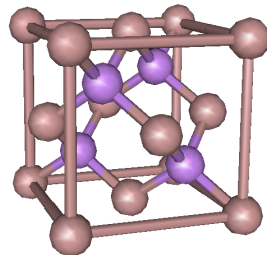
Heteroatomic structures



← bcc
sc →



← diamond
fcc →



The Ewald sphere



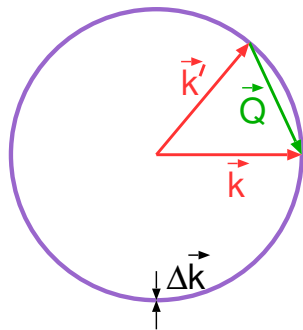
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The sphere radius is set by the length of the \vec{k} and \vec{k}' vectors which characterize the incident and scattered (where the detector is placed) x-rays and $\Delta\vec{k}$ being the bandwidth of the incident x-rays



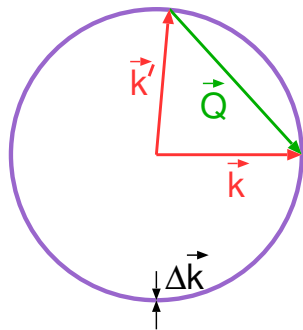
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As the detector moves, \vec{k}' rotates but the Ewald sphere remains constant.



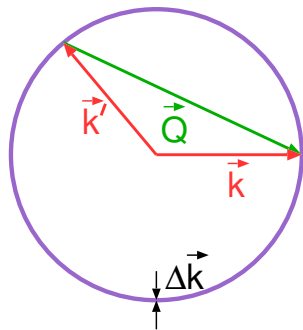
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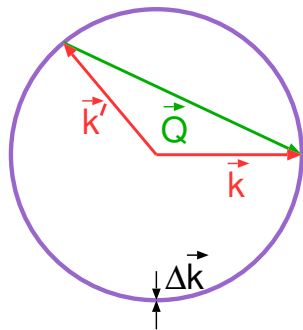


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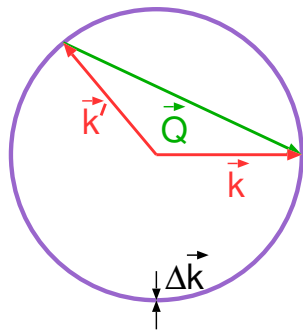
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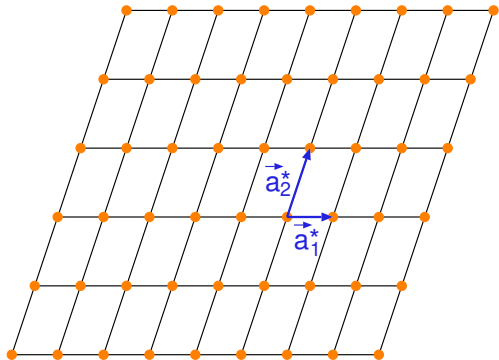
<http://www.phillipslab.org/software>



Ewald sphere & the reciprocal lattice



The reciprocal lattice is defined by the unit vectors \vec{a}_1^* and \vec{a}_2^* .

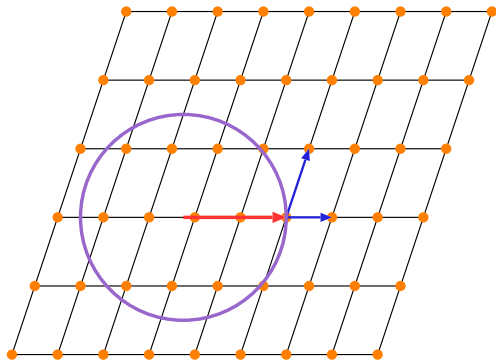


Ewald sphere & the reciprocal lattice



The reciprocal lattice is defined by the unit vectors \vec{a}_1^* and \vec{a}_2^* .

The key parameter is the relative orientation of the incident wave vector \vec{k}



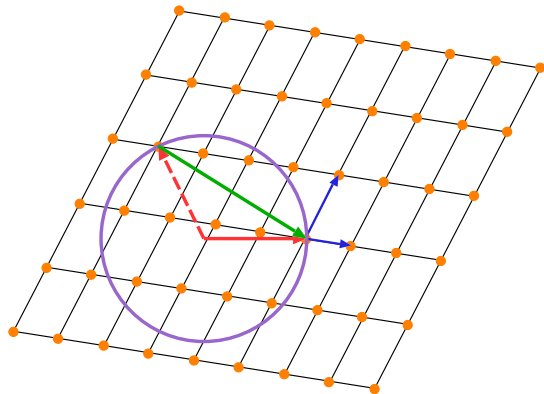
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As the crystal is rotated with respect to the incident beam, the reciprocal lattice also rotates



Ewald sphere & the reciprocal lattice

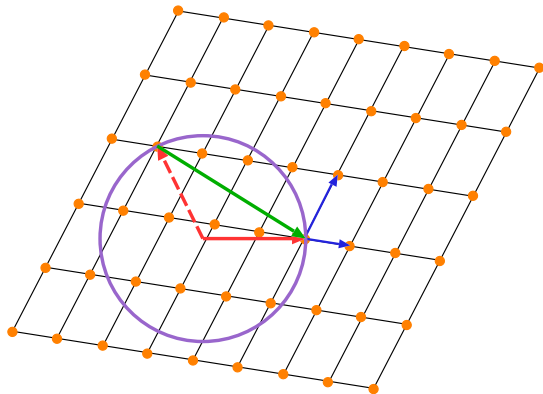


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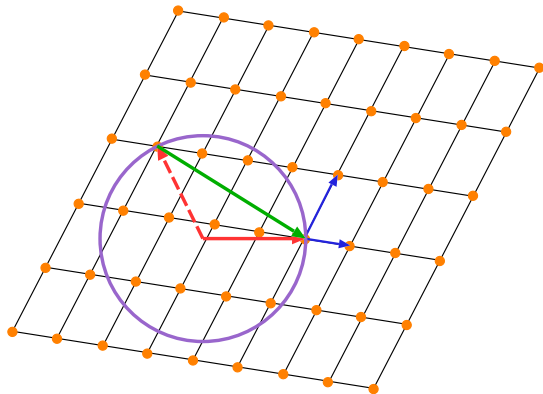


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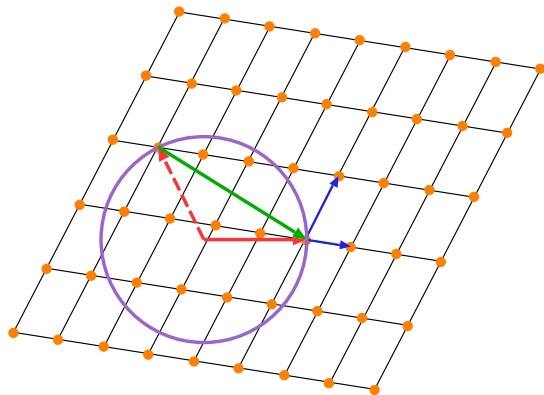


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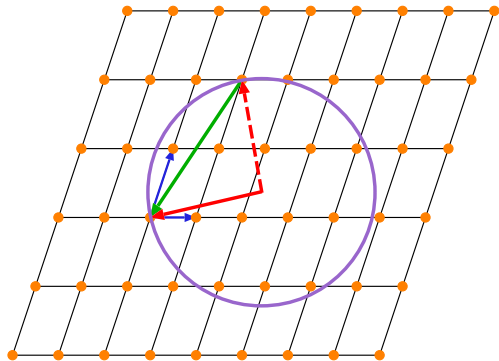


$$\vec{G}_{h|k} = h\vec{a}_1^* + k\vec{a}_2^*$$

Ewald construction



It is often more convenient to visualize the Ewald sphere by keeping the reciprocal lattice fixed and “rotating” the incident beam to visualize the scattering geometry.

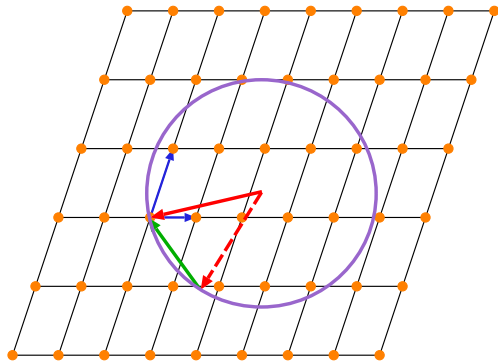


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In directions of \vec{k}' (detector position) where there is no reciprocal lattice point, there can be no diffraction peak.



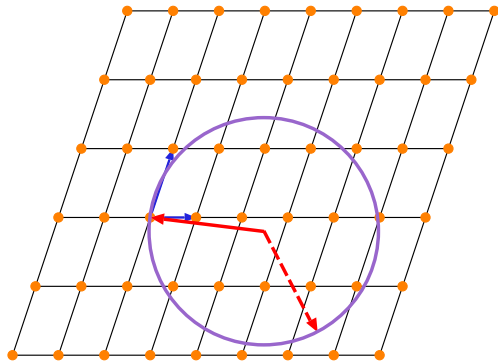
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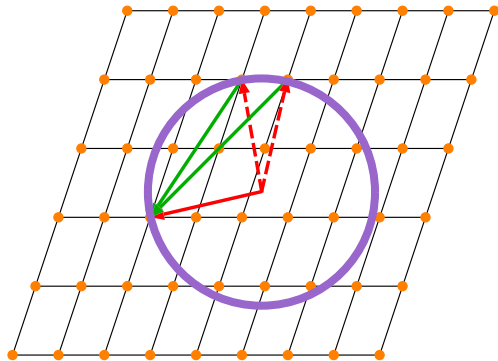
If the crystal is rotated slightly with respect to the incident beam, \vec{k} , there may be no Bragg reflections possible at all.



Polychromatic radiation



If $\Delta \vec{k}$ is large enough, there may be more than one reflection lying on the Ewald sphere.

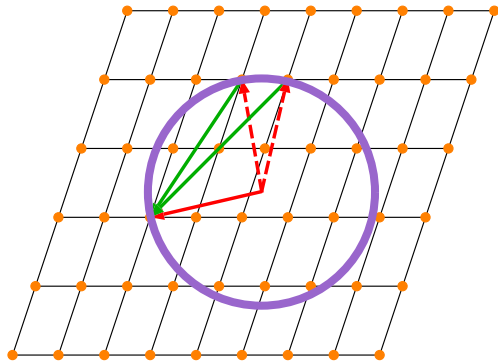


Polychromatic radiation



If $\Delta \vec{k}$ is large enough, there may be more than one reflection lying on the Ewald sphere.

With an area detector, there may then be multiple reflections appearing for a particular orientation (very common with protein crystals where the unit cell is very large).



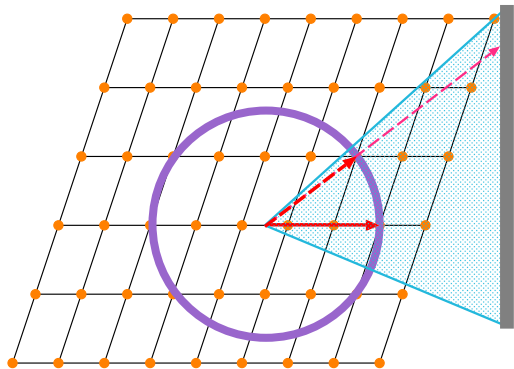
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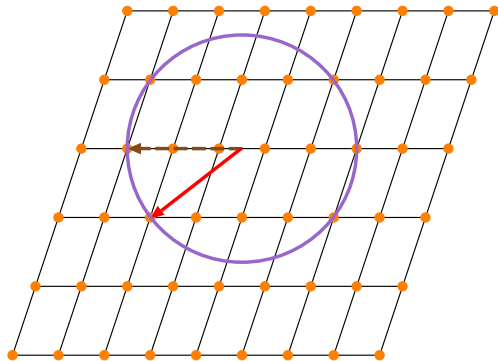
In protein crystallography, the area detector is in a fixed location with respect to the incident beam and the crystal is rotated on a spindle so that as Laue conditions are met, spots are produced on the detector at the diffraction angle



Multiple scattering



If more than one reciprocal lattice point is on the Ewald sphere, scattering can occur internal to the crystal.

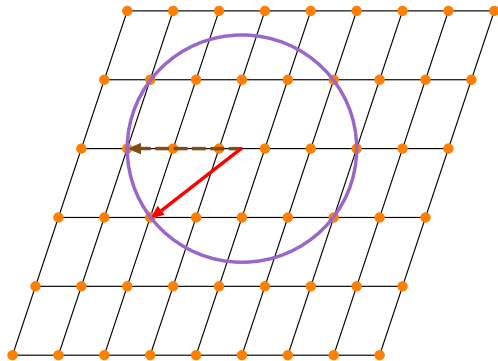


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The x-rays are first scattered along \vec{k}_{int}

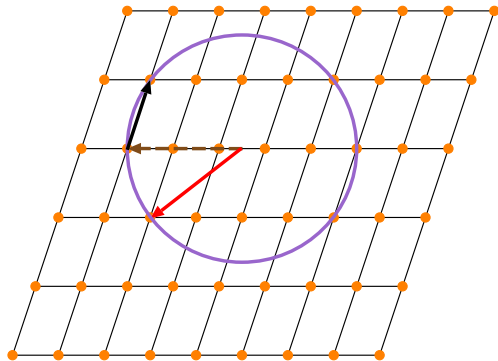


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The x-rays are first scattered along \vec{k}_{int} then along the reciprocal lattice vector which connects the two points on the Ewald sphere, \vec{G}

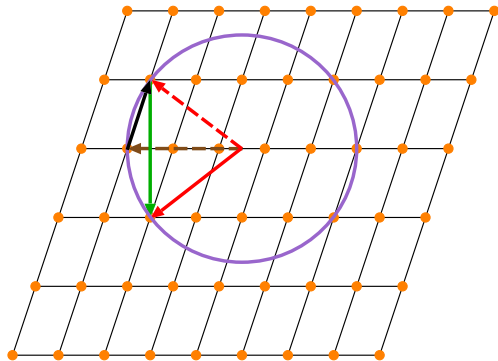


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The x-rays are first scattered along \vec{k}_{int} then along the reciprocal lattice vector which connects the two points on the Ewald sphere, \vec{G} and to the detector at \vec{k}' .



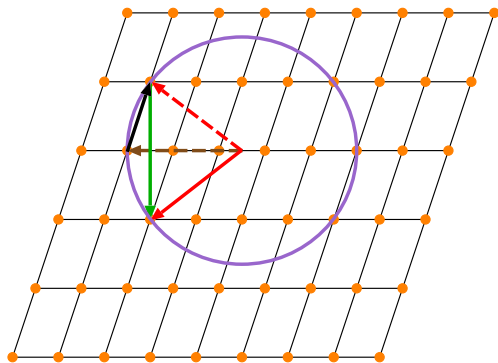
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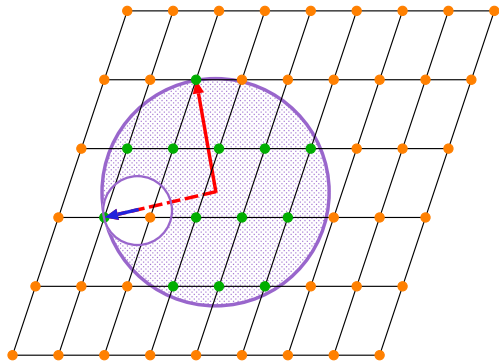
This is the cause of monochromator glitches which sometimes remove intensity but can also add intensity to the reflection the detector is set to measure.



Laue diffraction



The Laue diffraction technique uses a wide range of radiation from \vec{k}_{min} to \vec{k}_{max}

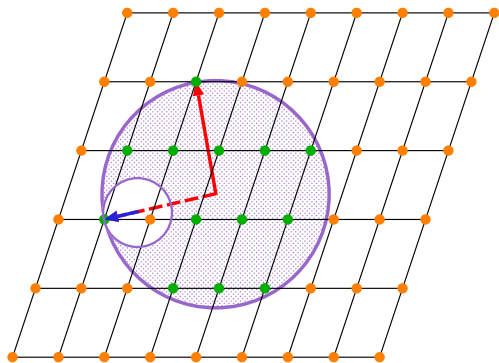


Laue diffraction



The Laue diffraction technique uses a wide range of radiation from \vec{k}_{min} to \vec{k}_{max}

These define two Ewald spheres and a volume between them such that any **reciprocal lattice point** which lies in the volume will meet the Laue condition for reflection.



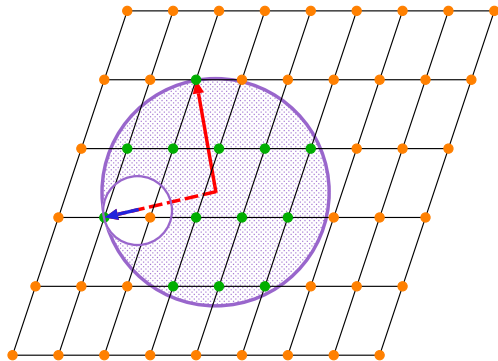
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This technique is useful for taking data on crystals which are changing or may degrade in the beam with a single shot of x-rays on a 2D detector.





XRayView

<http://www.phillipslab.org/downloads>

Diffraction resources



XRayView

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Bilbao Crystallography Server

<http://www.cryst.ehu.es/>



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GSAS-II

<https://subversion.xray.aps.anl.gov/trac/pyGSAS>



Exercise 1 - Ewald sphere

Exercise 4 - Wavelength

Exercise 8 - Laue diffraction

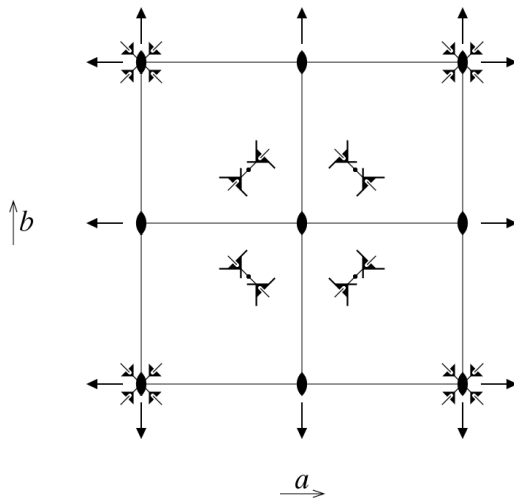
Exercise 9 - Serial crystallography

$P23$

$P 2 3$

23

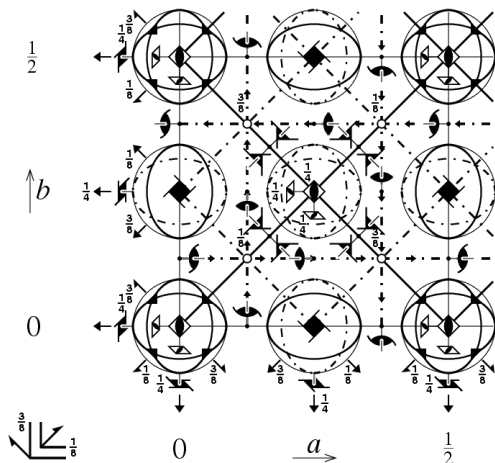
No. 195



- 1 x, y, z
- 2 x, \bar{y}, \bar{z}
- 3 \bar{x}, y, \bar{z}
- 4 \bar{x}, \bar{y}, z
- 5 z, x, y
- 6 \bar{z}, \bar{x}, y
- 7 z, \bar{x}, \bar{y}
- 8 \bar{z}, x, \bar{y}
- 9 y, z, x
- 10 \bar{y}, z, \bar{x}
- 11 \bar{y}, \bar{z}, x
- 12 y, \bar{z}, \bar{x}

$Fd\bar{3}m$ $F4_1/d\bar{3}2/m$ $m\bar{3}m$

No. 227



- | | |
|--------------------------------------------------------|--------------------------------------------------------|
| 1 x, y, z | 25 $\frac{1}{4} - x, \frac{1}{4} - y, \frac{1}{4} - z$ |
| 2 x, \bar{y}, \bar{z} | 26 $\frac{1}{4} - x, \frac{1}{4} + y, \frac{1}{4} + z$ |
| 3 \bar{x}, y, \bar{z} | 27 $\frac{1}{4} + x, \frac{1}{4} - y, \frac{1}{4} + z$ |
| 4 \bar{x}, \bar{y}, z | 28 $\frac{1}{4} + x, \frac{1}{4} + y, \frac{1}{4} - z$ |
| 5 z, x, y | 29 $\frac{1}{4} - z, \frac{1}{4} - x, \frac{1}{4} - y$ |
| 6 \bar{z}, \bar{x}, y | 30 $\frac{1}{4} + z, \frac{1}{4} + x, \frac{1}{4} - y$ |
| 7 z, \bar{x}, \bar{y} | 31 $\frac{1}{4} - z, \frac{1}{4} + x, \frac{1}{4} + y$ |
| 8 \bar{z}, x, \bar{y} | 32 $\frac{1}{4} + z, \frac{1}{4} - x, \frac{1}{4} + y$ |
| 9 y, z, x | 33 $\frac{1}{4} - y, \frac{1}{4} - z, \frac{1}{4} - x$ |
| 10 \bar{y}, z, \bar{x} | 34 $\frac{1}{4} + y, \frac{1}{4} - z, \frac{1}{4} + x$ |
| 11 \bar{y}, \bar{z}, x | 35 $\frac{1}{4} + y, \frac{1}{4} + z, \frac{1}{4} - x$ |
| 12 y, \bar{z}, \bar{x} | 36 $\frac{1}{4} - y, \frac{1}{4} + z, \frac{1}{4} + x$ |
| 13 $\frac{1}{4} + x, \frac{1}{4} - z, \frac{1}{4} + y$ | 37 \bar{x}, z, \bar{y} |
| 14 $\frac{1}{4} + x, \frac{1}{4} + z, \frac{1}{4} - y$ | 38 \bar{x}, \bar{z}, y |
| 15 $\frac{1}{4} - x, \frac{1}{4} - z, \frac{1}{4} - y$ | 39 x, z, y |
| 16 $\frac{1}{4} - x, \frac{1}{4} + z, \frac{1}{4} + y$ | 40 x, \bar{z}, \bar{y} |
| 17 $\frac{1}{4} + z, \frac{1}{4} + y, \frac{1}{4} - x$ | 41 \bar{z}, \bar{y}, x |
| 18 $\frac{1}{4} - z, \frac{1}{4} + y, \frac{1}{4} + x$ | 42 $\bar{z}, \bar{y}, \bar{x}$ |
| 19 $\frac{1}{4} - z, \frac{1}{4} - y, \frac{1}{4} - x$ | 43 z, y, x |
| 20 $\frac{1}{4} + z, \frac{1}{4} - y, \frac{1}{4} + x$ | 44 \bar{z}, y, \bar{x} |
| 21 $\frac{1}{4} - y, \frac{1}{4} + x, \frac{1}{4} + z$ | 45 y, \bar{x}, \bar{z} |
| 22 $\frac{1}{4} + y, \frac{1}{4} - x, \frac{1}{4} + z$ | 46 \bar{y}, x, \bar{z} |
| 23 $\frac{1}{4} - y, \frac{1}{4} - x, \frac{1}{4} - z$ | 47 y, x, z |
| 24 $\frac{1}{4} + y, \frac{1}{4} + x, \frac{1}{4} - z$ | 48 \bar{y}, \bar{x}, z |

 $+ (0, \frac{1}{2}, \frac{1}{2}), (\frac{1}{2}, 0, \frac{1}{2}), (\frac{1}{2}, \frac{1}{2}, 0)$


Wyckoff Positions of Group 195 ($P2_3$)

Multiplicity	Wyckoff letter	Site symmetry	Coordinates
12	j	1	(x,y,z) $(-x,-y,z)$ $(-x,y,-z)$ $(x,-y,-z)$ (z,x,y) $(z,-x,-y)$ $(-z,-x,y)$ $(-z,x,-y)$ (y,z,x) $(-y,z,-x)$ $(y,-z,-x)$ $(-y,-z,x)$
6	i	2..	$(x,1/2,1/2)$ $(-x,1/2,1/2)$ $(1/2,x,1/2)$ $(1/2,-x,1/2)$ $(1/2,1/2,x)$ $(1/2,1/2,-x)$
6	h	2..	$(x,1/2,0)$ $(-x,1/2,0)$ $(0,x,1/2)$ $(0,-x,1/2)$ $(1/2,0,x)$ $(1/2,0,-x)$
6	g	2..	$(x,0,1/2)$ $(-x,0,1/2)$ $(1/2,x,0)$ $(1/2,-x,0)$ $(0,1/2,x)$ $(0,1/2,-x)$
6	f	2..	$(x,0,0)$ $(-x,0,0)$ $(0,x,0)$ $(0,-x,0)$ $(0,0,x)$ $(0,0,-x)$
4	e	.3.	(x,x,x) $(-x,-x,x)$ $(-x,x,-x)$ $(x,-x,-x)$
3	d	222 . .	$(1/2,0,0)$ $(0,1/2,0)$ $(0,0,1/2)$
3	c	222 . .	$(0,1/2,1/2)$ $(1/2,0,1/2)$ $(1/2,1/2,0)$
1	b	23.	$(1/2,1/2,1/2)$
1	a	23.	$(0,0,0)$

Wyckoff Positions of Group 227 (*Fd-3m*) [origin choice 1]

Multiplicity	Wyckoff letter	Site symmetry	Coordinates			
			$(0,0,0) + (0,1/2,1/2) + (1/2,0,1/2) + (1/2,1/2,0) +$			
192	i	1	(x,y,z)	(-x,-y+1/2,z+1/2)	(-x+1/2,y+1/2,-z)	(x+1/2,-y,-z+1/2)
			(z,x,y)	(z+1/2,-x,-y+1/2)	(-z,-x+1/2,y+1/2)	(-z+1/2,x+1/2,-y)
			(y,z,x)	(-y+1/2,z+1/2,-x)	(y+1/2,-z,-x+1/2)	(-y,-z+1/2,x+1/2)
			(y+3/4,x+1/4,-z+3/4)	(-y+1/4,-x+1/4,-z+1/4)	(y+1/4,-x+3/4,z+3/4)	(-y+3/4,x+3/4,z+1/4)
			(x+3/4,z+1/4,-y+3/4)	(-x+3/4,z+3/4,y+1/4)	(-x+1/4,-z+1/4,-y+1/4)	(x+1/4,-z+3/4,y+3/4)
			(z+3/4,y+1/4,-x+3/4)	(z+1/4,-y+3/4,x+3/4)	(-z+3/4,y+3/4,x+1/4)	(-z+1/4,-y+1/4,-x+1/4)
			(-x+1/4,-y+1/4,-z+1/4)	(x+1/4,y+3/4,-z+3/4)	(x+3/4,-y+3/4,z+1/4)	(-x+3/4,y+1/4,z+3/4)
			(-z+1/4,-x+1/4,-y+1/4)	(-z+3/4,x+1/4,y+3/4)	(z+1/4,x+3/4,-y+3/4)	(z+3/4,-x+3/4,y+1/4)
			(-y+1/4,-z+1/4,-x+1/4)	(y+3/4,-z+3/4,x+1/4)	(-y+3/4,z+1/4,x+3/4)	(y+1/4,z+3/4,-x+3/4)
			(-y+1/2,-x,z+1/2)	(y,x,z)	(-y,x+1/2,-z+1/2)	(y+1/2,-x+1/2,-z)
			(-x+1/2,-z,y+1/2)	(x+1/2,-z+1/2,-y)	(x,z,y)	(-x,z+1/2,-y+1/2)
			(-z+1/2,-y,x+1/2)	(-z,y+1/2,-x+1/2)	(z+1/2,-y+1/2,-x)	(z,y,x)
			(1/8,y,-y+1/4)	(7/8,-y+1/2,-y+3/4)	(3/8,y+1/2,y+3/4)	(5/8,-y,y+1/4)
			(-y+1/4,1/8,y)	(-y+3/4,7/8,-y+1/2)	(y+3/4,3/8,y+1/2)	(y+1/4,5/8,-y)
96	h	.2	(y,-y+1/4,1/8)	(-y+1/2,-y+3/4,7/8)	(y+1/2,y+3/4,3/8)	(-y,y+1/4,5/8)
			(1/8,-y+1/4,y)	(3/8,y+3/4,y+1/2)	(7/8,-y+3/4,-y+1/2)	(5/8,y+1/4,-y)
			(y,1/8,-y+1/4)	(y+1/2,3/8,y+3/4)	(-y+1/2,7/8,-y+3/4)	(-y,5/8,y+1/4)
			(-y+1/4,y,1/8)	(y+3/4,y+1/2,3/8)	(-y+3/4,-y+1/2,7/8)	(y+1/4,-y,5/8)
96	g	.m	(x,x,z)	(-x,-x+1/2,z+1/2)	(-x+1/2,x+1/2,-z)	(x+1/2,-x,-z+1/2)
			(z,x,x)	(z+1/2,-x,-x+1/2)	(-z,-x+1/2,x+1/2)	(-z+1/2,x+1/2,-x)
			(x,z,x)	(-x+1/2,z+1/2,-x)	(x+1/2,-z,-x+1/2)	(-x,-z+1/2,x+1/2)
			(x+3/4,x+1/4,-z+3/4)	(-x+1/4,-x+1/4,-z+1/4)	(x+1/4,-x+3/4,z+3/4)	(-x+3/4,x+3/4,z+1/4)
			(x+3/4,z+1/4,-x+3/4)	(-x+3/4,z+3/4,x+1/4)	(-x+1/4,-z+1/4,-x+1/4)	(x+1/4,-z+3/4,x+3/4)
			(z+3/4,x+1/4,-x+3/4)	(z+1/4,-x+3/4,x+3/4)	(-z+3/4,x+3/4,x+1/4)	(-z+1/4,-x+1/4,-x+1/4)
48	f	2 m m	(x,0,0)	(-x,1/2,1/2)	(0,x,0)	(1/2,-x,1/2)
			(0,0,x)	(1/2,1/2,-x)	(3/4,x+1/4,3/4)	(1/4,-x+1/4,1/4)
			(x+3/4,1/4,3/4)	(-x+3/4,3/4,1/4)	(3/4,1/4,-x+3/4)	(1/4,3/4,x+3/4)
32	e	.3m	(x,x,x)	(-x,-x+1/2,x+1/2)	(-x+1/2,x+1/2,-x)	(x+1/2,-x,-x+1/2)
			(x+3/4,x+1/4,-x+3/4)	(-x+1/4,-x+1/4,-x+1/4)	(x+1/4,-x+3/4,x+3/4)	(-x+3/4,x+3/4,x+1/4)
16	d	-.3m	(5/8,5/8,5/8)	(3/8,7/8,1/8)	(7/8,1/8,3/8)	(1/8,3/8,7/8)
16	c	-.3m	(1/8,1/8,1/8)	(7/8,3/8,5/8)	(3/8,5/8,7/8)	(5/8,7/8,3/8)
8	b	-.43m	(1/2,1/2,1/2)	(1/4,3/4,1/4)		
8	a	-.43m	(0,0,0)	(3/4,1/4,3/4)		