

Today's outline - September 30, 2021



Today's outline - September 30, 2021



- Small angle x-ray scattering

Today's outline - September 30, 2021



- Small angle x-ray scattering
- Calculating R_g

Today's outline - September 30, 2021



- Small angle x-ray scattering
- Calculating R_g
- Porod analysis

Today's outline - September 30, 2021



- Small angle x-ray scattering
- Calculating R_g
- Porod analysis
- Polydispersivity

Today's outline - September 30, 2021



- Small angle x-ray scattering
- Calculating R_g
- Porod analysis
- Polydispersivity
- Unified fit model

Today's outline - September 30, 2021



- Small angle x-ray scattering
- Calculating R_g
- Porod analysis
- Polydispersivity
- Unified fit model
- Interparticle interactions

Today's outline - September 30, 2021



- Small angle x-ray scattering
- Calculating R_g
- Porod analysis
- Polydispersivity
- Unified fit model
- Interparticle interactions

Reading Assignment: Chapter 4.5; Chapter 5.1

Today's outline - September 30, 2021



- Small angle x-ray scattering
- Calculating R_g
- Porod analysis
- Polydispersivity
- Unified fit model
- Interparticle interactions

Reading Assignment: Chapter 4.5; Chapter 5.1

Homework Assignment #03:

Chapter 3: 1,3,4,6,8

due Tuesday, October 05, 2021

Today's outline - September 30, 2021



- Small angle x-ray scattering
- Calculating R_g
- Porod analysis
- Polydispersivity
- Unified fit model
- Interparticle interactions

Reading Assignment: Chapter 4.5; Chapter 5.1

Homework Assignment #03:

Chapter 3: 1,3,4,6,8

due Tuesday, October 05, 2021

Homework Assignment #04:

Chapter 4: 2,4,6,7.10

due Tuesday, October 19, 2021



The liquid scattering experiment measured the liquid structure factor, $S(Q)$, a short range order quantity.



The liquid scattering experiment measured the liquid structure factor, $S(Q)$, a short range order quantity.

by applying a Fourier transform, it is possible to extract the radial distribution function for the liquid, $g(r)$



The liquid scattering experiment measured the liquid structure factor, $S(Q)$, a short range order quantity.

by applying a Fourier transform, it is possible to extract the radial distribution function for the liquid, $g(r)$

$$g(r) = 1 + \frac{1}{2\pi^2 r \rho_{at}} \int_0^\infty Q [S(Q) - 1] \sin(Qr) dQ$$



The liquid scattering experiment measured the liquid structure factor, $S(Q)$, a short range order quantity.

by applying a Fourier transform, it is possible to extract the radial distribution function for the liquid, $g(r)$

$$g(r) = 1 + \frac{1}{2\pi^2 r \rho_{at}} \int_0^\infty Q [S(Q) - 1] \sin(Qr) dQ$$



The liquid scattering experiment measured the liquid structure factor, $S(Q)$, a short range order quantity.

by applying a Fourier transform, it is possible to extract the radial distribution function for the liquid, $g(r)$

$$g(r) = 1 + \frac{1}{2\pi^2 r \rho_{at}} \int_0^\infty Q [S(Q) - 1] \sin(Qr) dQ$$

This formalism holds for both non-crystalline solids and liquids, even though inelastic scattering dominates in the latter.



The liquid scattering experiment measured the liquid structure factor, $S(Q)$, a short range order quantity.

by applying a Fourier transform, it is possible to extract the radial distribution function for the liquid, $g(r)$

$$g(r) = 1 + \frac{1}{2\pi^2 r \rho_{at}} \int_0^\infty Q [S(Q) - 1] \sin(Qr) dQ$$

This formalism holds for both non-crystalline solids and liquids, even though inelastic scattering dominates in the latter.

The relation between radial distribution function and structure factor can be extended to multi-component systems where $g(r) \rightarrow g_{ij}(r)$ and $S(Q) \rightarrow S_{ij}(Q)$.



Liquid scattering can be used to study dynamics

“Seeing real-space dynamics of liquid water through inelastic x-ray scattering,” T. Iwashita et al. *Sci. Adv.* **3**, e1603079 (2017).



Liquid scattering can be used to study dynamics

In this article, the authors measured the liquid scattering as a function of both momentum, Q , and energy, E , transfer by using analyzers set for a specific energy (21.747 keV) but varying Q and then scanning the incident energy at fixed incident angle

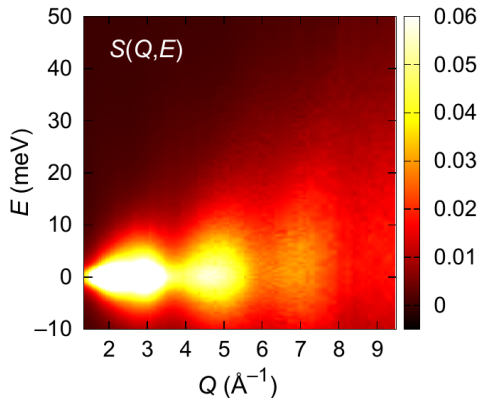
“Seeing real-space dynamics of liquid water through inelastic x-ray scattering,” T. Iwashita et al. *Sci. Adv.* **3**, e1603079 (2017).

Water dynamics



Liquid scattering can be used to study dynamics

In this article, the authors measured the liquid scattering as a function of both momentum, Q , and energy, E , transfer by using analyzers set for a specific energy (21.747 keV) but varying Q and then scanning the incident energy at fixed incident angle



“Seeing real-space dynamics of liquid water through inelastic x-ray scattering,” T. Iwashita et al. *Sci. Adv.* **3**, e1603079 (2017).

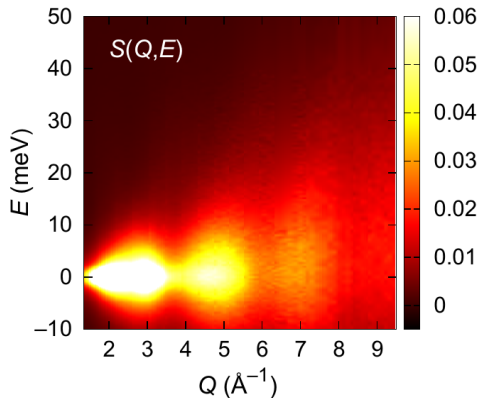
Water dynamics



Liquid scattering can be used to study dynamics

In this article, the authors measured the liquid scattering as a function of both momentum, Q , and energy, E , transfer by using analyzers set for a specific energy (21.747 keV) but varying Q and then scanning the incident energy at fixed incident angle

The Van Hoff function can be obtained by a double Fourier transform



“Seeing real-space dynamics of liquid water through inelastic x-ray scattering,” T. Iwashita et al. *Sci. Adv.* **3**, e1603079 (2017).

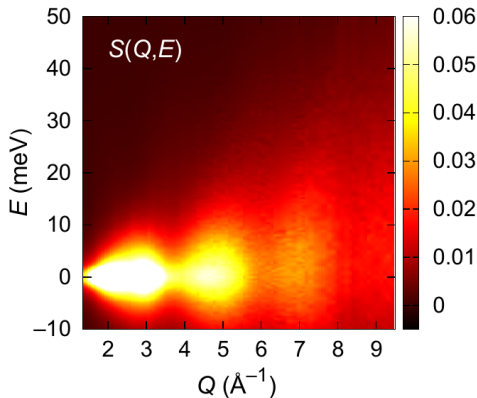


Liquid scattering can be used to study dynamics

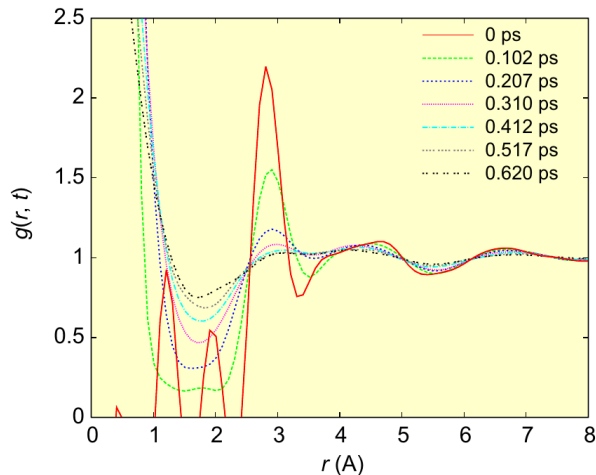
In this article, the authors measured the liquid scattering as a function of both momentum, Q , and energy, E , transfer by using analyzers set for a specific energy (21.747 keV) but varying Q and then scanning the incident energy at fixed incident angle

The Van Hoff function can be obtained by a double Fourier transform

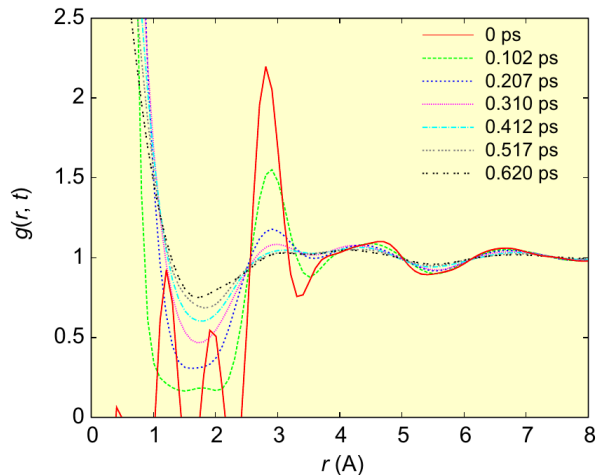
$$g(r, t) - 1 = \frac{1}{2\rho\pi^2 r} \int \int e^{i\omega t} \sin(Qr) Q S(Q, E) dQ dE$$



“Seeing real-space dynamics of liquid water through inelastic x-ray scattering,” T. Iwashita et al. *Sci. Adv.* **3**, e1603079 (2017).

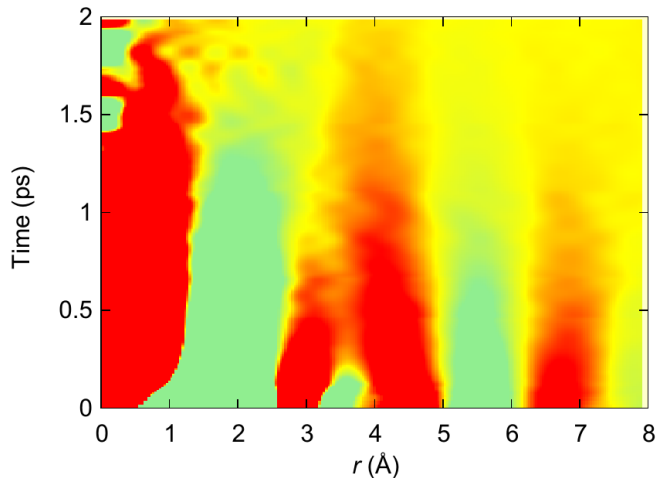


"Seeing real-space dynamics of liquid water through inelastic x-ray scattering," T. Iwashita et al. *Sci. Adv.* **3**, e1603079 (2017).



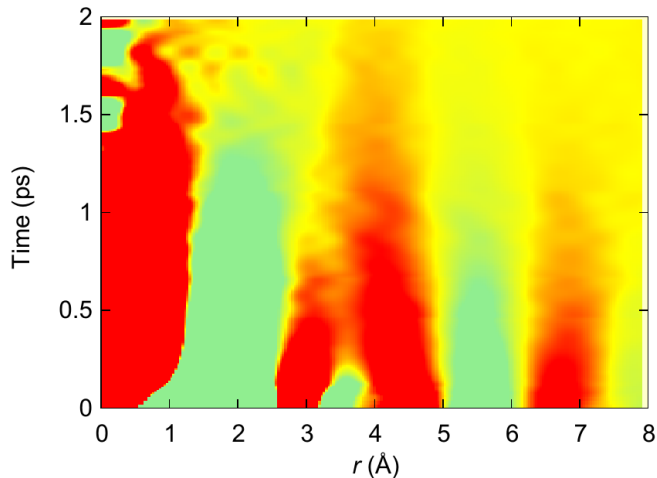
The peak below $1.5\text{\AA}$ represents the self motion of the central atom while the data at longer distances represents the collective motions of two different atoms, in this case the oxygens

“Seeing real-space dynamics of liquid water through inelastic x-ray scattering,” T. Iwashita et al. *Sci. Adv.* **3**, e1603079 (2017).



The peak below $1.5\text{\AA}$ represents the self motion of the central atom while the data at longer distances represents the collective motions of two different atoms, in this case the oxygens

“Seeing real-space dynamics of liquid water through inelastic x-ray scattering,” T. Iwashita et al. *Sci. Adv.* **3**, e1603079 (2017).



The peak below $1.5\text{\AA}$ represents the self motion of the central atom while the data at longer distances represents the collective motions of two different atoms, in this case the oxygens

The first and second peaks are highly coupled in space and time and merge within 0.8 ps. This behavior is different from liquid metals and leads to the viscosity of water.

“Seeing real-space dynamics of liquid water through inelastic x-ray scattering,” T. Iwashita et al. *Sci. Adv.* **3**, e1603079 (2017).

Small angle x-ray scattering



$$I(\vec{Q}) = Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \int_V [\rho_n(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m + f(\vec{Q})^2 \rho_{at} \sum_n \int_V e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m$$

Recall that there was an additional term in the scattering intensity which becomes important at small Q .

Small angle x-ray scattering



$$I(\vec{Q}) = Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \int_V [\rho_n(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m + f(\vec{Q})^2 \rho_{at} \sum_n \int_V e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m$$

Recall that there was an additional term in the scattering intensity which becomes important at small Q .

Small angle x-ray scattering



$$I(\vec{Q}) = Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \int_V [\rho_n(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m + f(\vec{Q})^2 \rho_{at} \sum_n \int_V e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m$$

Recall that there was an additional term in the scattering intensity which becomes important at small Q .

$$I^{SAXS}(\vec{Q}) = f^2 \sum_n \int_V \rho_{at} e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m$$

Small angle x-ray scattering



$$I(\vec{Q}) = Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \int_V [\rho_n(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m + f(\vec{Q})^2 \rho_{at} \sum_n \int_V e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m$$

Recall that there was an additional term in the scattering intensity which becomes important at small Q .

$$I^{SAXS}(\vec{Q}) = f^2 \sum_n \int_V \rho_{at} e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m = f^2 \sum_n e^{i\vec{Q} \cdot \vec{r}_n} \int_V \rho_{at} e^{-i\vec{Q} \cdot \vec{r}_m} dV_m$$



$$I(\vec{Q}) = Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \int_V [\rho_n(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m + f(\vec{Q})^2 \rho_{at} \sum_n \int_V e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m$$

Recall that there was an additional term in the scattering intensity which becomes important at small Q .

$$\begin{aligned} I^{SAXS}(\vec{Q}) &= f^2 \sum_n \int_V \rho_{at} e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m = f^2 \sum_n e^{i\vec{Q} \cdot \vec{r}_n} \int_V \rho_{at} e^{-i\vec{Q} \cdot \vec{r}_m} dV_m \\ &= f^2 \int_V \rho_{at} e^{i\vec{Q} \cdot \vec{r}_n} dV_n \int_V \rho_{at} e^{-i\vec{Q} \cdot \vec{r}_m} dV_m \end{aligned}$$

Small angle x-ray scattering



$$I(\vec{Q}) = Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \int_V [\rho_n(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m + f(\vec{Q})^2 \rho_{at} \sum_n \int_V e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m$$

Recall that there was an additional term in the scattering intensity which becomes important at small Q .

$$\begin{aligned} I^{SAXS}(\vec{Q}) &= f^2 \sum_n \int_V \rho_{at} e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m = f^2 \sum_n e^{i\vec{Q} \cdot \vec{r}_n} \int_V \rho_{at} e^{-i\vec{Q} \cdot \vec{r}_m} dV_m \\ &= f^2 \int_V \rho_{at} e^{i\vec{Q} \cdot \vec{r}_n} dV_n \int_V \rho_{at} e^{-i\vec{Q} \cdot \vec{r}_m} dV_m = \left| \int_V \rho_{sl} e^{i\vec{Q} \cdot \vec{r}} dV \right|^2 \end{aligned}$$

Small angle x-ray scattering



$$I(\vec{Q}) = Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \int_V [\rho_n(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m + f(\vec{Q})^2 \rho_{at} \sum_n \int_V e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m$$

Recall that there was an additional term in the scattering intensity which becomes important at small Q .

$$\begin{aligned} I^{SAXS}(\vec{Q}) &= f^2 \sum_n \int_V \rho_{at} e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m = f^2 \sum_n e^{i\vec{Q} \cdot \vec{r}_n} \int_V \rho_{at} e^{-i\vec{Q} \cdot \vec{r}_m} dV_m \\ &= f^2 \int_V \rho_{at} e^{i\vec{Q} \cdot \vec{r}_n} dV_n \int_V \rho_{at} e^{-i\vec{Q} \cdot \vec{r}_m} dV_m = \left| \int_V \rho_{sl} e^{i\vec{Q} \cdot \vec{r}} dV \right|^2 \end{aligned}$$

Where we have assumed sufficient averaging and introduced $\rho_{sl} = f\rho_{at}$.

Small angle x-ray scattering



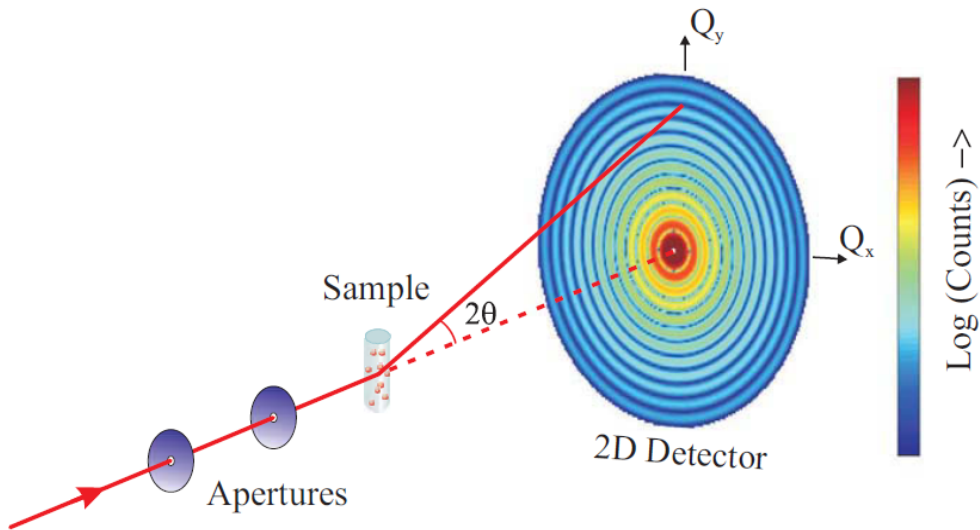
$$I(\vec{Q}) = Nf(\vec{Q})^2 + f(\vec{Q})^2 \sum_n \int_V [\rho_n(\vec{r}_{nm}) - \rho_{at}] e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m + f(\vec{Q})^2 \rho_{at} \sum_n \int_V e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m$$

Recall that there was an additional term in the scattering intensity which becomes important at small Q .

$$\begin{aligned} I^{SAXS}(\vec{Q}) &= f^2 \sum_n \int_V \rho_{at} e^{i\vec{Q} \cdot (\vec{r}_n - \vec{r}_m)} dV_m = f^2 \sum_n e^{i\vec{Q} \cdot \vec{r}_n} \int_V \rho_{at} e^{-i\vec{Q} \cdot \vec{r}_m} dV_m \\ &= f^2 \int_V \rho_{at} e^{i\vec{Q} \cdot \vec{r}_n} dV_n \int_V \rho_{at} e^{-i\vec{Q} \cdot \vec{r}_m} dV_m = \left| \int_V \rho_{sl} e^{i\vec{Q} \cdot \vec{r}} dV \right|^2 \end{aligned}$$

Where we have assumed sufficient averaging and introduced $\rho_{sl} = f \rho_{at}$. This final expression looks just like an atomic form factor but the charge density that we consider here is on a much longer length scale than an atom.

The SAXS experiment



Scattering from a dilute solution



The simplest case is for a dilute solution of non-interacting molecules.

$$I^{SAXS}(\vec{Q}) = \left| \int_{V_p} \rho_s e^{i\vec{Q} \cdot \vec{r}} dV_p \right|^2$$

Scattering from a dilute solution



The simplest case is for a dilute solution of non-interacting molecules.

Assume that the scattering length density of each identical particle (molecule) is given by $\rho_{sl,p}$ and the scattering length density of the solvent is $\rho_{sl,0}$.

$$I^{SAXS}(\vec{Q}) = \left| \int_{V_p} \rho_{sl} e^{i\vec{Q} \cdot \vec{r}} dV_p \right|^2$$

Scattering from a dilute solution



The simplest case is for a dilute solution of non-interacting molecules.

Assume that the scattering length density of each identical particle (molecule) is given by $\rho_{sl,p}$ and the scattering length density of the solvent is $\rho_{sl,0}$.

$$I^{SAXS}(\vec{Q}) = \left| \int_{V_p} \rho_{sl} e^{i\vec{Q} \cdot \vec{r}} dV_p \right|^2 = (\rho_{sl,p} - \rho_{sl,0})^2 \left| \int_{V_p} e^{i\vec{Q} \cdot \vec{r}} dV_p \right|^2$$

Scattering from a dilute solution



The simplest case is for a dilute solution of non-interacting molecules.

Assume that the scattering length density of each identical particle (molecule) is given by $\rho_{sl,p}$ and the scattering length density of the solvent is $\rho_{sl,0}$.

$$I^{SAXS}(\vec{Q}) = \left| \int_{V_p} \rho_{sl} e^{i\vec{Q} \cdot \vec{r}} dV_p \right|^2 = (\rho_{sl,p} - \rho_{sl,0})^2 \left| \int_{V_p} e^{i\vec{Q} \cdot \vec{r}} dV_p \right|^2$$

If we introduce the single-particle form factor $\mathcal{F}(\vec{Q})$:

Scattering from a dilute solution



The simplest case is for a dilute solution of non-interacting molecules.

Assume that the scattering length density of each identical particle (molecule) is given by $\rho_{sl,p}$ and the scattering length density of the solvent is $\rho_{sl,0}$.

$$I^{SAXS}(\vec{Q}) = \left| \int_{V_p} \rho_{sl} e^{i\vec{Q} \cdot \vec{r}} dV_p \right|^2 = (\rho_{sl,p} - \rho_{sl,0})^2 \left| \int_{V_p} e^{i\vec{Q} \cdot \vec{r}} dV_p \right|^2$$

If we introduce the single-particle form factor $\mathcal{F}(\vec{Q})$:

$$\mathcal{F}(\vec{Q}) = \frac{1}{V_p} \int_{V_p} e^{i\vec{Q} \cdot \vec{r}} dV_p$$

Scattering from a dilute solution



The simplest case is for a dilute solution of non-interacting molecules.

Assume that the scattering length density of each identical particle (molecule) is given by $\rho_{sl,p}$ and the scattering length density of the solvent is $\rho_{sl,0}$.

$$I^{SAXS}(\vec{Q}) = \left| \int_{V_p} \rho_{sl} e^{i\vec{Q} \cdot \vec{r}} dV_p \right|^2 = (\rho_{sl,p} - \rho_{sl,0})^2 \left| \int_{V_p} e^{i\vec{Q} \cdot \vec{r}} dV_p \right|^2$$

If we introduce the single-particle form factor $\mathcal{F}(\vec{Q})$:

$$\mathcal{F}(\vec{Q}) = \frac{1}{V_p} \int_{V_p} e^{i\vec{Q} \cdot \vec{r}} dV_p$$

$$I^{SAXS}(\vec{Q}) = \Delta\rho^2 V_p^2 |\mathcal{F}(\vec{Q})|^2$$

Scattering from a dilute solution



The simplest case is for a dilute solution of non-interacting molecules.

Assume that the scattering length density of each identical particle (molecule) is given by $\rho_{sl,p}$ and the scattering length density of the solvent is $\rho_{sl,0}$.

$$I^{SAXS}(\vec{Q}) = \left| \int_{V_p} \rho_{sl} e^{i\vec{Q} \cdot \vec{r}} dV_p \right|^2 = (\rho_{sl,p} - \rho_{sl,0})^2 \left| \int_{V_p} e^{i\vec{Q} \cdot \vec{r}} dV_p \right|^2$$

If we introduce the single-particle form factor $\mathcal{F}(\vec{Q})$:

$$\mathcal{F}(\vec{Q}) = \frac{1}{V_p} \int_{V_p} e^{i\vec{Q} \cdot \vec{r}} dV_p \qquad I^{SAXS}(\vec{Q}) = \Delta\rho^2 V_p^2 |\mathcal{F}(\vec{Q})|^2$$

Where $\Delta\rho = (\rho_{sl,p} - \rho_{sl,0})$, and the form factor depends on the morphology of the particle (size and shape).

Scattering from a sphere



There are only a few morphologies which can be computed exactly and the simplest is a constant density sphere of radius R .

Scattering from a sphere



There are only a few morphologies which can be computed exactly and the simplest is a constant density sphere of radius R .

$$\mathcal{F}(Q) = \frac{1}{V_p} \int_0^R \int_0^{2\pi} \int_0^\pi e^{iQr \cos \theta} r^2 \sin \theta \, d\theta \, d\phi \, dr$$

Scattering from a sphere



There are only a few morphologies which can be computed exactly and the simplest is a constant density sphere of radius R .

$$\mathcal{F}(Q) = \frac{1}{V_p} \int_0^R \int_0^{2\pi} \int_0^\pi e^{iQr \cos \theta} r^2 \sin \theta d\theta d\phi dr$$

$$\mathcal{F}(Q) = \frac{1}{V_p} \int_0^R 4\pi \frac{\sin(Qr)}{Qr} r^2 dr$$

Scattering from a sphere



There are only a few morphologies which can be computed exactly and the simplest is a constant density sphere of radius R .

$$\mathcal{F}(Q) = \frac{1}{V_p} \int_0^R \int_0^{2\pi} \int_0^\pi e^{iQr \cos \theta} r^2 \sin \theta \, d\theta \, d\phi \, dr$$

$$\begin{aligned} \mathcal{F}(Q) &= \frac{1}{V_p} \int_0^R 4\pi \frac{\sin(Qr)}{Qr} r^2 \, dr \\ &= 3 \left[\frac{\sin(QR) - QR \cos(QR)}{Q^3 R^3} \right] \end{aligned}$$

Scattering from a sphere



There are only a few morphologies which can be computed exactly and the simplest is a constant density sphere of radius R .

$$\mathcal{F}(Q) = \frac{1}{V_p} \int_0^R \int_0^{2\pi} \int_0^\pi e^{iQr \cos \theta} r^2 \sin \theta \, d\theta \, d\phi \, dr$$

$$\begin{aligned} \mathcal{F}(Q) &= \frac{1}{V_p} \int_0^R 4\pi \frac{\sin(Qr)}{Qr} r^2 \, dr \\ &= 3 \left[\frac{\sin(QR) - QR \cos(QR)}{Q^3 R^3} \right] \equiv \frac{3J_1(QR)}{QR} \end{aligned}$$

Scattering from a sphere



There are only a few morphologies which can be computed exactly and the simplest is a constant density sphere of radius R .

$$\mathcal{F}(Q) = \frac{1}{V_p} \int_0^R \int_0^{2\pi} \int_0^\pi e^{iQr \cos \theta} r^2 \sin \theta d\theta d\phi dr$$

$$\begin{aligned} \mathcal{F}(Q) &= \frac{1}{V_p} \int_0^R 4\pi \frac{\sin(Qr)}{Qr} r^2 dr \\ &= 3 \left[\frac{\sin(QR) - QR \cos(QR)}{Q^3 R^3} \right] \equiv \frac{3J_1(QR)}{QR} \end{aligned}$$

Where $J_1(x)$ is the Bessel function of the first kind

Scattering from a sphere

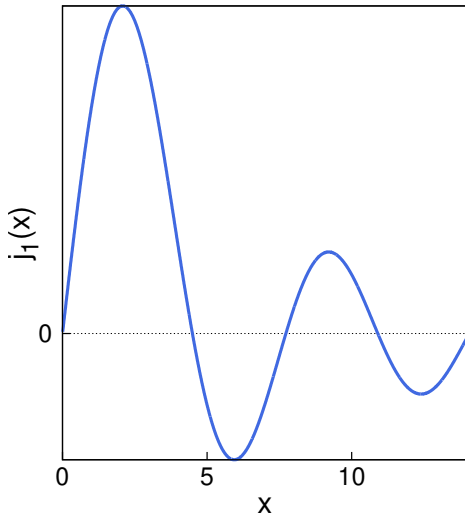


There are only a few morphologies which can be computed exactly and the simplest is a constant density sphere of radius R .

$$\mathcal{F}(Q) = \frac{1}{V_p} \int_0^R \int_0^{2\pi} \int_0^\pi e^{iQr \cos \theta} r^2 \sin \theta d\theta d\phi dr$$

$$\begin{aligned} \mathcal{F}(Q) &= \frac{1}{V_p} \int_0^R 4\pi \frac{\sin(Qr)}{Qr} r^2 dr \\ &= 3 \left[\frac{\sin(QR) - QR \cos(QR)}{Q^3 R^3} \right] \equiv \frac{3J_1(QR)}{QR} \end{aligned}$$

Where $J_1(x)$ is the Bessel function of the first kind



Scattering from a sphere



$$I(Q) = \Delta\rho^2 V_p^2 \left| \frac{3J_1(QR)}{QR} \right|^2$$

Scattering from a sphere

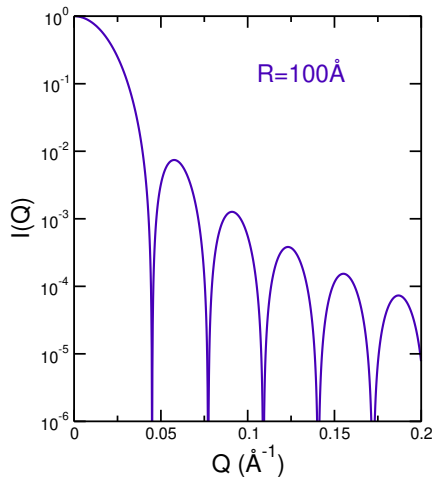


$$I(Q) = \Delta\rho^2 V_p^2 \left| \frac{3J_1(QR)}{QR} \right|^2 = \Delta\rho^2 V_p^2 \left| 3 \frac{\sin(QR) - QR \cos(QR)}{Q^3 R^3} \right|^2$$

Scattering from a sphere



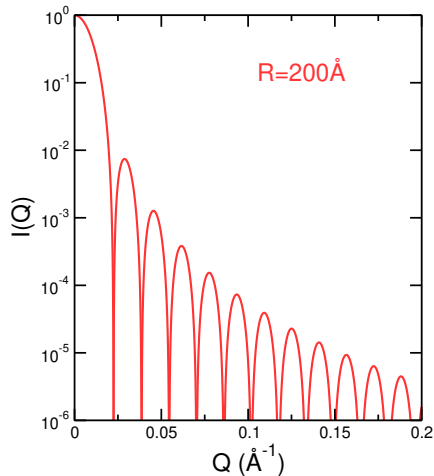
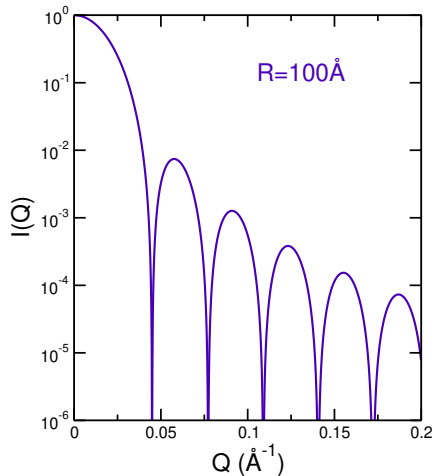
$$I(Q) = \Delta\rho^2 V_p^2 \left| \frac{3J_1(QR)}{QR} \right|^2 = \Delta\rho^2 V_p^2 \left| 3 \frac{\sin(QR) - QR \cos(QR)}{Q^3 R^3} \right|^2$$



Scattering from a sphere



$$I(Q) = \Delta\rho^2 V_p^2 \left| \frac{3J_1(QR)}{QR} \right|^2 = \Delta\rho^2 V_p^2 \left| 3 \frac{\sin(QR) - QR \cos(QR)}{Q^3 R^3} \right|^2$$



Guinier analysis



$$I^{SAXS}(\vec{Q}) = \Delta\rho^2 V_p^2 |\mathcal{F}(\vec{Q})|^2,$$

Guinier analysis



$$I^{SAXS}(\vec{Q}) = \Delta\rho^2 V_p^2 |\mathcal{F}(\vec{Q})|^2,$$

$$\mathcal{F}(\vec{Q}) = 3 \left[\frac{\sin(QR) - QR \cos(QR)}{Q^3 R^3} \right]$$

Guinier analysis



$$I^{SAXS}(\vec{Q}) = \Delta\rho^2 V_p^2 |\mathcal{F}(\vec{Q})|^2,$$

$$\mathcal{F}(\vec{Q}) = 3 \left[\frac{\sin(QR) - QR \cos(QR)}{Q^3 R^3} \right]$$

In the long wavelength limit $QR \rightarrow 0$ we can approximate the scattering factor with the first terms of the sum



$$I^{SAXS}(\vec{Q}) = \Delta\rho^2 V_p^2 |\mathcal{F}(\vec{Q})|^2,$$

$$\mathcal{F}(\vec{Q}) = 3 \left[\frac{\sin(QR) - QR \cos(QR)}{Q^3 R^3} \right]$$

In the long wavelength limit $QR \rightarrow 0$ we can approximate the scattering factor with the first terms of the sum

$$\mathcal{F}(Q) \approx \frac{3}{Q^3 R^3} \left[QR - \frac{Q^3 R^3}{6} + \frac{Q^5 R^5}{120} - \dots - QR \left(1 - \frac{Q^2 R^2}{2} + \frac{Q^4 R^4}{24} - \dots \right) \right]$$



$$I^{SAXS}(\vec{Q}) = \Delta\rho^2 V_p^2 |\mathcal{F}(\vec{Q})|^2,$$

$$\mathcal{F}(\vec{Q}) = 3 \left[\frac{\sin(QR) - QR \cos(QR)}{Q^3 R^3} \right]$$

In the long wavelength limit $QR \rightarrow 0$ we can approximate the scattering factor with the first terms of the sum

$$\mathcal{F}(Q) \approx \frac{3}{Q^3 R^3} \left[QR - \frac{Q^3 R^3}{6} + \frac{Q^5 R^5}{120} - \dots - QR \left(1 - \frac{Q^2 R^2}{2} + \frac{Q^4 R^4}{24} - \dots \right) \right]$$

this simplifies to $\mathcal{F}(Q) \approx 1 - \frac{Q^2 R^2}{10}$ and



$$I^{SAXS}(\vec{Q}) = \Delta\rho^2 V_p^2 |\mathcal{F}(\vec{Q})|^2,$$

$$\mathcal{F}(\vec{Q}) = 3 \left[\frac{\sin(QR) - QR \cos(QR)}{Q^3 R^3} \right]$$

In the long wavelength limit $QR \rightarrow 0$ we can approximate the scattering factor with the first terms of the sum

$$\mathcal{F}(Q) \approx \frac{3}{Q^3 R^3} \left[QR - \frac{Q^3 R^3}{6} + \frac{Q^5 R^5}{120} - \dots - QR \left(1 - \frac{Q^2 R^2}{2} + \frac{Q^4 R^4}{24} - \dots \right) \right]$$

this simplifies to $\mathcal{F}(Q) \approx 1 - \frac{Q^2 R^2}{10}$ and

$$I^{SAXS}(Q) \approx \Delta\rho^2 V_p^2 \left[1 - \frac{Q^2 R^2}{10} \right]^2$$



$$I^{SAXS}(\vec{Q}) = \Delta\rho^2 V_p^2 |\mathcal{F}(\vec{Q})|^2, \quad \mathcal{F}(\vec{Q}) = 3 \left[\frac{\sin(QR) - QR \cos(QR)}{Q^3 R^3} \right]$$

In the long wavelength limit $QR \rightarrow 0$ we can approximate the scattering factor with the first terms of the sum

$$\mathcal{F}(Q) \approx \frac{3}{Q^3 R^3} \left[QR - \frac{Q^3 R^3}{6} + \frac{Q^5 R^5}{120} - \dots - QR \left(1 - \frac{Q^2 R^2}{2} + \frac{Q^4 R^4}{24} - \dots \right) \right]$$

this simplifies to $\mathcal{F}(Q) \approx 1 - \frac{Q^2 R^2}{10}$ and

$$I^{SAXS}(Q) \approx \Delta\rho^2 V_p^2 \left[1 - \frac{Q^2 R^2}{10} \right]^2 \approx \Delta\rho^2 V_p^2 \left[1 - \frac{Q^2 R^2}{5} \right]$$



$$I^{SAXS}(\vec{Q}) = \Delta\rho^2 V_p^2 |\mathcal{F}(\vec{Q})|^2, \quad \mathcal{F}(\vec{Q}) = 3 \left[\frac{\sin(QR) - QR \cos(QR)}{Q^3 R^3} \right]$$

In the long wavelength limit $QR \rightarrow 0$ we can approximate the scattering factor with the first terms of the sum

$$\mathcal{F}(Q) \approx \frac{3}{Q^3 R^3} \left[QR - \frac{Q^3 R^3}{6} + \frac{Q^5 R^5}{120} - \dots - QR \left(1 - \frac{Q^2 R^2}{2} + \frac{Q^4 R^4}{24} - \dots \right) \right]$$

this simplifies to $\mathcal{F}(Q) \approx 1 - \frac{Q^2 R^2}{10}$ and

$$I^{SAXS}(Q) \approx \Delta\rho^2 V_p^2 \left[1 - \frac{Q^2 R^2}{10} \right]^2 \approx \Delta\rho^2 V_p^2 \left[1 - \frac{Q^2 R^2}{5} \right] \approx \Delta\rho^2 V_p^2 e^{-Q^2 R^2/5}, \quad QR \ll 1$$

Guinier analysis



In the long wavelength limit ($QR \rightarrow 0$), the form factor becomes

Guinier analysis



In the long wavelength limit ($QR \rightarrow 0$), the form factor becomes

$$\mathcal{F}(Q) \approx 1 - \frac{Q^2 R^2}{10}$$

Guinier analysis



In the long wavelength limit ($QR \rightarrow 0$), the form factor becomes

$$\mathcal{F}(Q) \approx 1 - \frac{Q^2 R^2}{10}$$
$$I(Q) \approx \Delta\rho^2 V_p^2 e^{-Q^2 R^2/5}$$

Guinier analysis



In the long wavelength limit ($QR \rightarrow 0$), the form factor becomes

$$\mathcal{F}(Q) \approx 1 - \frac{Q^2 R^2}{10}$$
$$I(Q) \approx \Delta\rho^2 V_p^2 e^{-Q^2 R^2/5}$$

and the initial slope of the $\log(I)$ vs Q^2 plot is $-R^2/5$

Guinier analysis

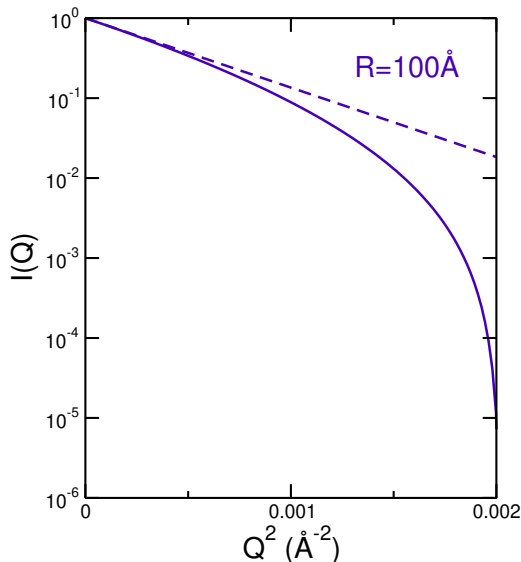


In the long wavelength limit ($QR \rightarrow 0$), the form factor becomes

$$\mathcal{F}(Q) \approx 1 - \frac{Q^2 R^2}{10}$$

$$I(Q) \approx \Delta\rho^2 V_p^2 e^{-Q^2 R^2/5}$$

and the initial slope of the $\log(I)$ vs Q^2 plot is $-R^2/5$



Guinier analysis

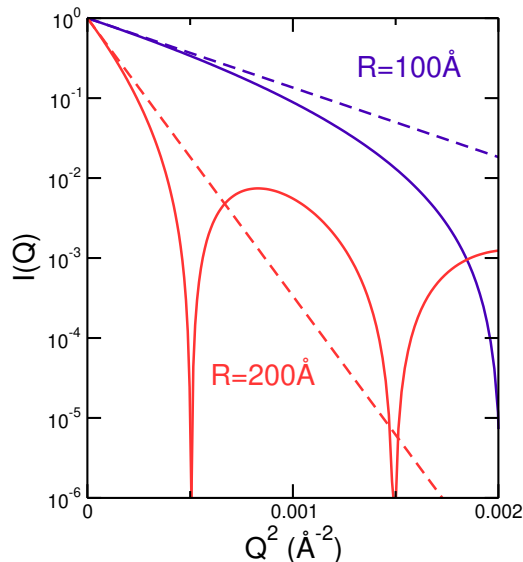


In the long wavelength limit ($QR \rightarrow 0$), the form factor becomes

$$\mathcal{F}(Q) \approx 1 - \frac{Q^2 R^2}{10}$$

$$I(Q) \approx \Delta\rho^2 V_p^2 e^{-Q^2 R^2/5}$$

and the initial slope of the $\log(I)$ vs Q^2 plot is $-R^2/5$



Guinier analysis



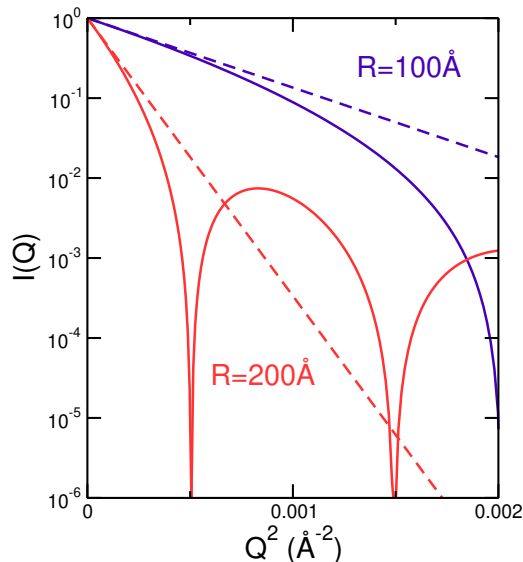
In the long wavelength limit ($QR \rightarrow 0$), the form factor becomes

$$\mathcal{F}(Q) \approx 1 - \frac{Q^2 R^2}{10}$$

$$I(Q) \approx \Delta\rho^2 V_p^2 e^{-Q^2 R^2/5}$$

and the initial slope of the $\log(I)$ vs Q^2 plot is $-R^2/5$

In terms of the radius of gyration, R_g , which for a sphere is given by $\sqrt{\frac{3}{5}}R$



Guinier analysis



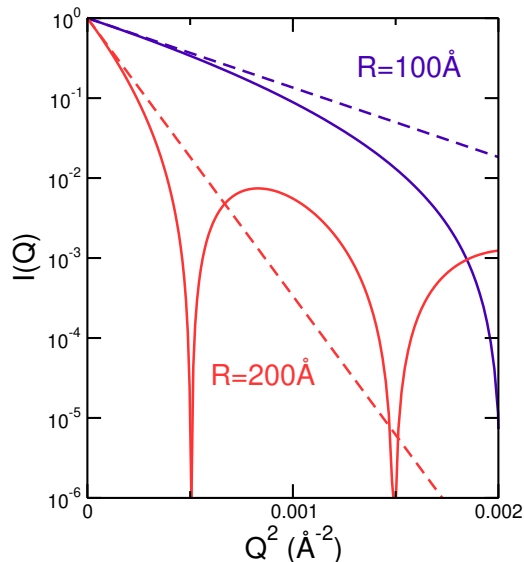
In the long wavelength limit ($QR \rightarrow 0$), the form factor becomes

$$\mathcal{F}(Q) \approx 1 - \frac{Q^2 R^2}{10}$$
$$I(Q) \approx \Delta\rho^2 V_p^2 e^{-Q^2 R^2/5}$$

and the initial slope of the $\log(I)$ vs Q^2 plot is $-R^2/5$

In terms of the radius of gyration, R_g , which for a sphere is given by $\sqrt{\frac{3}{5}}R$

$$I(Q) \approx \Delta\rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$



Calculation of R_g



$$I(Q) \approx \Delta \rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$

Calculation of R_g



$$I(Q) \approx \Delta\rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$

The radius of gyration R_g is defined as the second moment of the volume occupied by the particle



$$I(Q) \approx \Delta\rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$

The radius of gyration R_g is defined as the second moment of the volume occupied by the particle

$$R_g^2 = \frac{1}{V_p} \int_{V_p} r^2 dV_p$$



$$I(Q) \approx \Delta\rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$

The radius of gyration R_g is defined as the second moment of the volume occupied by the particle

in terms of the scattering length density, it can be rewritten as

$$R_g^2 = \frac{1}{V_p} \int_{V_p} r^2 dV_p$$



$$I(Q) \approx \Delta\rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$

The radius of gyration R_g is defined as the second moment of the volume occupied by the particle

in terms of the scattering length density, it can be rewritten as

$$R_g^2 = \frac{1}{V_p} \int_{V_p} r^2 dV_p = \frac{\int_{V_p} \rho_{sl,p}(\vec{r}) r^2 dV_p}{\int_{V_p} \rho_{sl,p}(\vec{r}) dV_p}$$



$$I(Q) \approx \Delta\rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$

The radius of gyration R_g is defined as the second moment of the volume occupied by the particle

in terms of the scattering length density, it can be rewritten as

$$R_g^2 = \frac{1}{V_p} \int_{V_p} r^2 dV_p = \frac{\int_{V_p} \rho_{sl,p}(\vec{r}) r^2 dV_p}{\int_{V_p} \rho_{sl,p}(\vec{r}) dV_p}$$

after orientational averaging this expression can be used to extract R_g from experimental data using



$$I(Q) \approx \Delta\rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$

The radius of gyration R_g is defined as the second moment of the volume occupied by the particle

in terms of the scattering length density, it can be rewritten as

$$R_g^2 = \frac{1}{V_p} \int_{V_p} r^2 dV_p = \frac{\int_{V_p} \rho_{sl,p}(\vec{r}) r^2 dV_p}{\int_{V_p} \rho_{sl,p}(\vec{r}) dV_p}$$

after orientational averaging this expression can be used to extract R_g from experimental data using

$$I_1^{SAXS}(Q) \approx \Delta\rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$



$$I(Q) \approx \Delta\rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$

The radius of gyration R_g is defined as the second moment of the volume occupied by the particle

in terms of the scattering length density, it can be rewritten as

$$R_g^2 = \frac{1}{V_p} \int_{V_p} r^2 dV_p = \frac{\int_{V_p} \rho_{sl,p}(\vec{r}) r^2 dV_p}{\int_{V_p} \rho_{sl,p}(\vec{r}) dV_p}$$

after orientational averaging this expression can be used to extract R_g from experimental data using

$$I_1^{SAXS}(Q) \approx \Delta\rho^2 V_p^2 e^{-Q^2 R_g^2/3}$$

this expression holds for uniform and non-uniform densities

Porod analysis



In the short wavelength limit ($QR \gg 1$), the form factor for a sphere can be approximated

Porod analysis



In the short wavelength limit ($QR \gg 1$), the form factor for a sphere can be approximated

$$\mathcal{F}(Q) = 3 \left[\frac{\sin(QR)}{Q^3 R^3} - \frac{\cos(QR)}{Q^2 R^2} \right]$$

Porod analysis



In the short wavelength limit ($QR \gg 1$), the form factor for a sphere can be approximated

$$\mathcal{F}(Q) = 3 \left[\frac{\sin(QR)}{Q^3 R^3} - \frac{\cos(QR)}{Q^2 R^2} \right] \approx 3 \left[-\frac{\cos(QR)}{Q^2 R^2} \right]$$



In the short wavelength limit ($QR \gg 1$), the form factor for a sphere can be approximated

$$\mathcal{F}(Q) = 3 \left[\frac{\sin(QR)}{Q^3 R^3} - \frac{\cos(QR)}{Q^2 R^2} \right] \approx 3 \left[-\frac{\cos(QR)}{Q^2 R^2} \right]$$

$$I(Q) = 9\Delta\rho^2 V_p^2 \left[-\frac{\cos(QR)}{Q^2 R^2} \right]^2$$



In the short wavelength limit ($QR \gg 1$), the form factor for a sphere can be approximated

$$\mathcal{F}(Q) = 3 \left[\frac{\sin(QR)}{Q^3 R^3} - \frac{\cos(QR)}{Q^2 R^2} \right] \approx 3 \left[-\frac{\cos(QR)}{Q^2 R^2} \right]$$

$$\begin{aligned} I(Q) &= 9\Delta\rho^2 V_p^2 \left[-\frac{\cos(QR)}{Q^2 R^2} \right]^2 \\ &= 9\Delta\rho^2 V_p^2 \frac{\langle \cos^2(QR) \rangle}{Q^4 R^4} \end{aligned}$$



In the short wavelength limit ($QR \gg 1$), the form factor for a sphere can be approximated

$$\mathcal{F}(Q) = 3 \left[\frac{\sin(QR)}{Q^3 R^3} - \frac{\cos(QR)}{Q^2 R^2} \right] \approx 3 \left[-\frac{\cos(QR)}{Q^2 R^2} \right]$$

$$\begin{aligned} I(Q) &= 9\Delta\rho^2 V_p^2 \left[-\frac{\cos(QR)}{Q^2 R^2} \right]^2 \\ &= 9\Delta\rho^2 V_p^2 \frac{\langle \cos^2(QR) \rangle}{Q^4 R^4} = \frac{9\Delta\rho^2 V_p^2}{Q^4 R^4} \left(\frac{1}{2} \right) \end{aligned}$$



In the short wavelength limit ($QR \gg 1$), the form factor for a sphere can be approximated

$$\mathcal{F}(Q) = 3 \left[\frac{\sin(QR)}{Q^3 R^3} - \frac{\cos(QR)}{Q^2 R^2} \right] \approx 3 \left[-\frac{\cos(QR)}{Q^2 R^2} \right]$$

$$\begin{aligned} I(Q) &= 9\Delta\rho^2 V_p^2 \left[-\frac{\cos(QR)}{Q^2 R^2} \right]^2 \\ &= 9\Delta\rho^2 V_p^2 \frac{\langle \cos^2(QR) \rangle}{Q^4 R^4} = \frac{9\Delta\rho^2 V_p^2}{Q^4 R^4} \left(\frac{1}{2} \right) \end{aligned}$$

$$I(Q) = \frac{2\pi\Delta\rho^2}{Q^4} S_p$$

Porod analysis

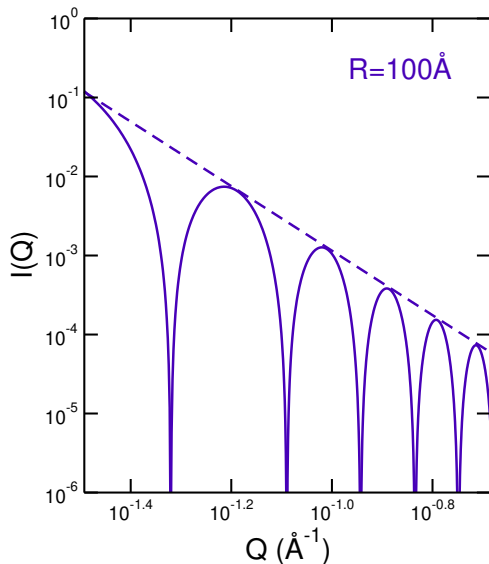


In the short wavelength limit ($QR \gg 1$), the form factor for a sphere can be approximated

$$\mathcal{F}(Q) = 3 \left[\frac{\sin(QR)}{Q^3 R^3} - \frac{\cos(QR)}{Q^2 R^2} \right] \approx 3 \left[-\frac{\cos(QR)}{Q^2 R^2} \right]$$

$$\begin{aligned} I(Q) &= 9\Delta\rho^2 V_p^2 \left[-\frac{\cos(QR)}{Q^2 R^2} \right]^2 \\ &= 9\Delta\rho^2 V_p^2 \frac{\langle \cos^2(QR) \rangle}{Q^4 R^4} = \frac{9\Delta\rho^2 V_p^2}{Q^4 R^4} \left(\frac{1}{2} \right) \end{aligned}$$

$$I(Q) = \frac{2\pi\Delta\rho^2}{Q^4} S_p$$



Porod analysis



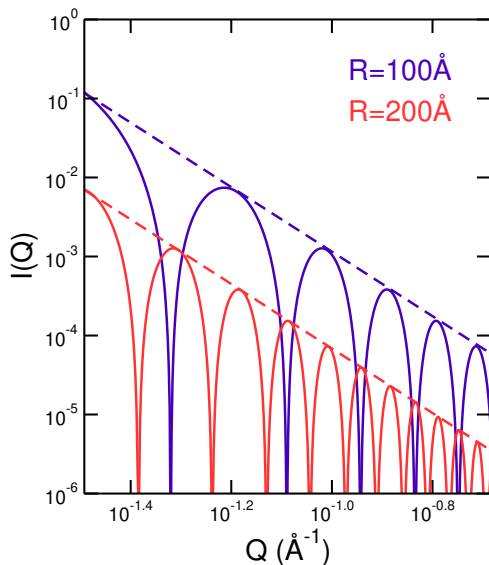
In the short wavelength limit ($QR \gg 1$), the form factor for a sphere can be approximated

$$\mathcal{F}(Q) = 3 \left[\frac{\sin(QR)}{Q^3 R^3} - \frac{\cos(QR)}{Q^2 R^2} \right] \approx 3 \left[-\frac{\cos(QR)}{Q^2 R^2} \right]$$

$$\begin{aligned} I(Q) &= 9\Delta\rho^2 V_p^2 \left[-\frac{\cos(QR)}{Q^2 R^2} \right]^2 \\ &= 9\Delta\rho^2 V_p^2 \frac{\langle \cos^2(QR) \rangle}{Q^4 R^4} = \frac{9\Delta\rho^2 V_p^2}{Q^4 R^4} \left(\frac{1}{2} \right) \end{aligned}$$

$$I(Q) = \frac{2\pi\Delta\rho^2}{Q^4} S_p$$

power law drop as Q^{-4} for spheres





Shape effect on scattering

The shape of the particle will have a significant effect on the SAXS since the form factor is derived from an integral over the particle volume, V_p .

Shape effect on scattering



The shape of the particle will have a significant effect on the SAXS since the form factor is derived from an integral over the particle volume, V_p .

If the particle is not spherical, then its “dimensionality” is not 3 and this will affect the form factor and introduce a different power law in the Porod regime.

Shape effect on scattering



The shape of the particle will have a significant effect on the SAXS since the form factor is derived from an integral over the particle volume, V_p .

If the particle is not spherical, then its “dimensionality” is not 3 and this will affect the form factor and introduce a different power law in the Porod regime.

$$\frac{dV_p}{dr} = 4\pi r^2 \frac{dr}{dr}$$

	shape	order
	sphere	



Shape effect on scattering

The shape of the particle will have a significant effect on the SAXS since the form factor is derived from an integral over the particle volume, V_p .

If the particle is not spherical, then its “dimensionality” is not 3 and this will affect the form factor and introduce a different power law in the Porod regime.

	shape	order
$dV_p = 4\pi r^2 dr$	sphere	
$dA_p = 2\pi r dr$	disk	



Shape effect on scattering

The shape of the particle will have a significant effect on the SAXS since the form factor is derived from an integral over the particle volume, V_p .

If the particle is not spherical, then its “dimensionality” is not 3 and this will affect the form factor and introduce a different power law in the Porod regime.

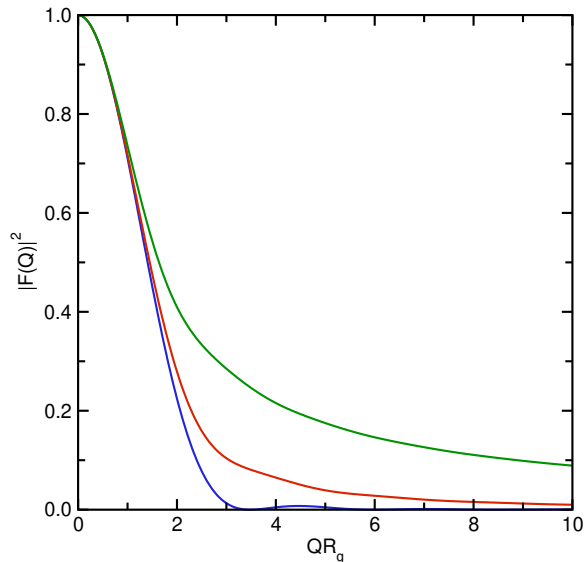
	shape	order
$dV_p = 4\pi r^2 dr$	sphere	
$dA_p = 2\pi r dr$	disk	
$dL_p = dr$	rod	

Shape effect on scattering



The shape of the particle will have a significant effect on the SAXS since the form factor is derived from an integral over the particle volume, V_p .

If the particle is not spherical, then its “dimensionality” is not 3 and this will affect the form factor and introduce a different power law in the Porod regime.



	shape	order
$dV_p = 4\pi r^2 dr$	sphere	
$dA_p = 2\pi r dr$	disk	
$dL_p = dr$	rod	

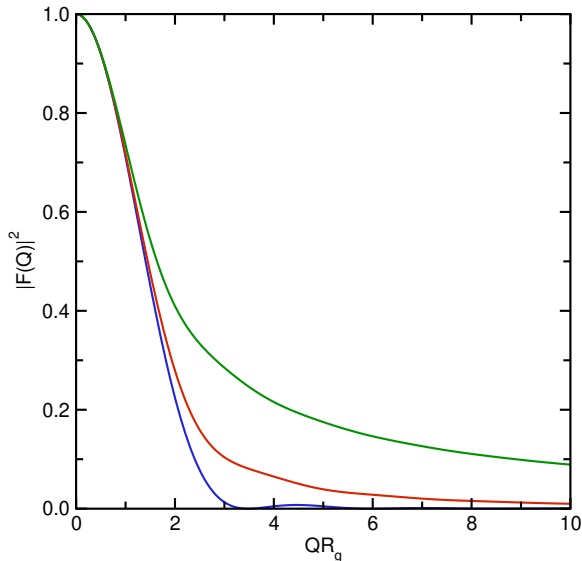
Shape effect on scattering



The shape of the particle will have a significant effect on the SAXS since the form factor is derived from an integral over the particle volume, V_p .

If the particle is not spherical, then its “dimensionality” is not 3 and this will affect the form factor and introduce a different power law in the Porod regime.

	shape	order
$dV_p = 4\pi r^2 dr$	sphere	-4
$dA_p = 2\pi r dr$	disk	
$dL_p = dr$	rod	



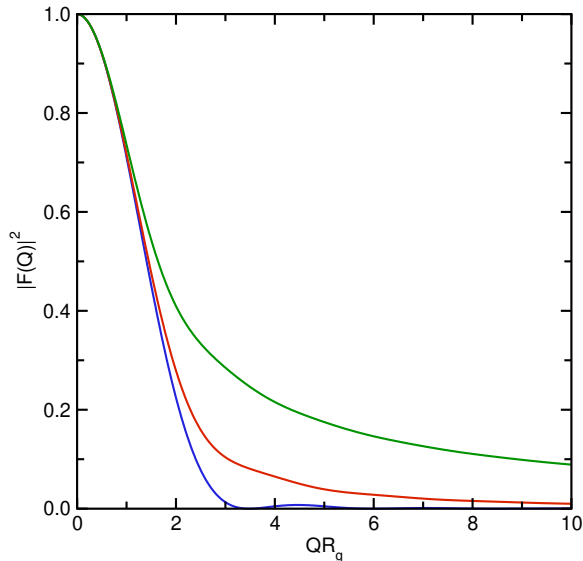
Shape effect on scattering



The shape of the particle will have a significant effect on the SAXS since the form factor is derived from an integral over the particle volume, V_p .

If the particle is not spherical, then its “dimensionality” is not 3 and this will affect the form factor and introduce a different power law in the Porod regime.

	shape	order
$dV_p = 4\pi r^2 dr$	sphere	-4
$dA_p = 2\pi r dr$	disk	-2
$dL_p = dr$	rod	



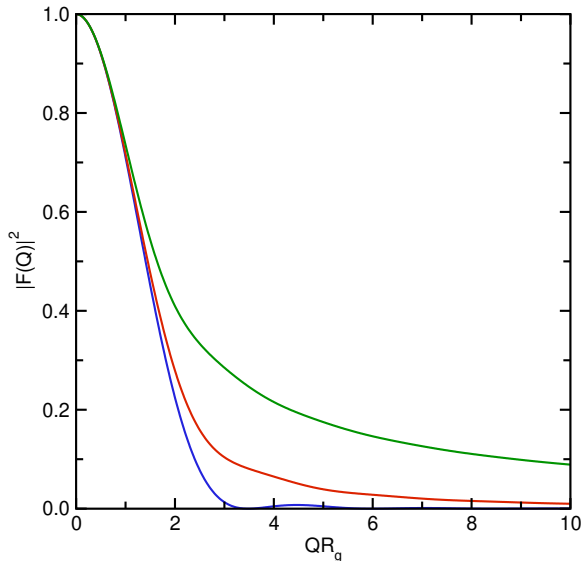
Shape effect on scattering



The shape of the particle will have a significant effect on the SAXS since the form factor is derived from an integral over the particle volume, V_p .

If the particle is not spherical, then its “dimensionality” is not 3 and this will affect the form factor and introduce a different power law in the Porod regime.

	shape	order
$dV_p = 4\pi r^2 dr$	sphere	-4
$dA_p = 2\pi r dr$	disk	-2
$dL_p = dr$	rod	-1



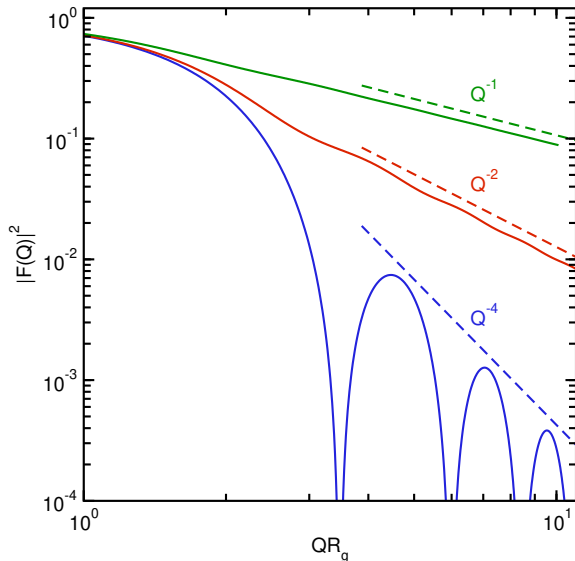
Shape effect on scattering



The shape of the particle will have a significant effect on the SAXS since the form factor is derived from an integral over the particle volume, V_p .

If the particle is not spherical, then its “dimensionality” is not 3 and this will affect the form factor and introduce a different power law in the Porod regime.

	shape	order
$dV_p = 4\pi r^2 dr$	sphere	-4
$dA_p = 2\pi r dr$	disk	-2
$dL_p = dr$	rod	-1



Polydispersivity



Analysis is simple when all particles are the same size, i.e. monodispersed and dilute,

Polydispersivity



Analysis is simple when all particles are the same size, i.e. monodispersed and dilute, more complex models must be used when the particles are polydispersed with a distribution function $D(R)$



Analysis is simple when all particles are the same size, i.e. monodispersed and dilute, more complex models must be used when the particles are polydispersed with a distribution function $D(R)$

$$I^{SAXS}(Q) = \Delta\rho^2 \int_0^\infty D(R) V_p(R)^2 |\mathcal{F}(Q, R)|^2 dR$$



Analysis is simple when all particles are the same size, i.e. monodispersed and dilute, more complex models must be used when the particles are polydispersed with a distribution function $D(R)$

$$I^{SAXS}(Q) = \Delta\rho^2 \int_0^\infty D(R) V_p(R)^2 |\mathcal{F}(Q, R)|^2 dR$$

the Schulz function is commonly used to model $D(R)$ as it goes to a delta function as the percentage polydispersivity, $p \rightarrow 0$

Polydispersity



Analysis is simple when all particles are the same size, i.e. monodispersed and dilute, more complex models must be used when the particles are polydispersed with a distribution function $D(R)$

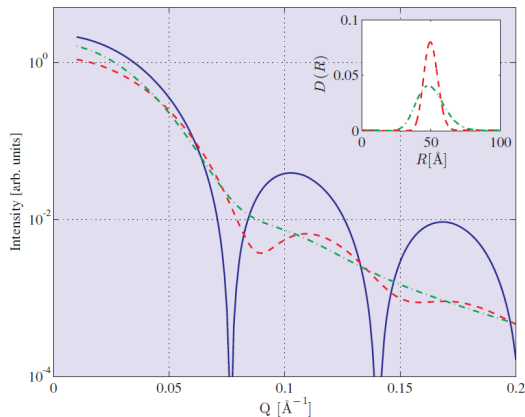
$$I^{SAXS}(Q) = \Delta\rho^2 \int_0^\infty D(R) V_p(R)^2 |\mathcal{F}(Q, R)|^2 dR$$

the Schulz function is commonly used to model $D(R)$ as it goes to a delta function as the percentage polydispersity, $p \rightarrow 0$

$$p = 0$$

$$p = 10\%$$

$$p = 20\%$$



Unified model for SAXS



For inhomogeneous samples, using simple Guinier and Porod models can give an incorrect description of the actual sample

Unified model for SAXS



For inhomogeneous samples, using simple Guinier and Porod models can give an incorrect description of the actual sample

This can be solved, judiciously, using the Unified model proposed by Beaucage in 1995

Unified model for SAXS



For inhomogeneous samples, using simple Guinier and Porod models can give an incorrect description of the actual sample

This can be solved, judiciously, using the Unified model proposed by Beaucage in 1995

The idea is to include multiple populations each with its own Guinier dependence with R_g and a power law dependence that asymptotically approaches the Porod law for spheres

Unified model for SAXS



For inhomogeneous samples, using simple Guinier and Porod models can give an incorrect description of the actual sample

This can be solved, judiciously, using the Unified model proposed by Beaucage in 1995

The idea is to include multiple populations each with its own Guinier dependence with R_g and a power law dependence that asymptotically approaches the Porod law for spheres

$$I(q) = B_{bkg} + \sum_{i=1}^N G_i e^{-\frac{q^2 R_{g,i}^2}{3}} + e^{-\frac{q^2 R_{g,i-1}^2}{3}} B_i \left[\frac{\left(\operatorname{erf} \left\{ \frac{q R_{g,i}}{\sqrt{6}} \right\} \right)^3}{q} \right]^{P_i}$$



For inhomogeneous samples, using simple Guinier and Porod models can give an incorrect description of the actual sample

This can be solved, judiciously, using the Unified model proposed by Beaucage in 1995

The idea is to include multiple populations each with its own Guinier dependence with R_g and a power law dependence that asymptotically approaches the Porod law for spheres

$$I(q) = B_{bkg} + \sum_{i=1}^N G_i e^{-\frac{q^2 R_{g,i}^2}{3}} + e^{-\frac{q^2 R_{g,i-1}^2}{3}} B_i \left[\frac{\left(\operatorname{erf} \left\{ \frac{q R_{g,i}}{\sqrt{6}} \right\} \right)^3}{q} \right]^{P_i}$$

The sum is over structural levels starting with the smallest. For each level there is a Guinier exponential prefactor (G_i), a radius of gyration ($R_{g,i}$), a power law constant prefactor (B_i), and a power law exponent (P_i).



$$I(q) = B_{bkg} + \sum_{i=1}^N G_i e^{-\frac{q^2 R_{g,i}^2}{3}} + e^{-\frac{q^2 R_{g,i-1}^2}{3}} B_i \left[\frac{\left(\operatorname{erf} \left\{ \frac{q R_{g,i}}{\sqrt{6}} \right\} \right)^3}{q} \right]^{P_i}$$



$$I(q) = B_{bkg} + \sum_{i=1}^N G_i e^{-\frac{q^2 R_{g,i}^2}{3}} + e^{-\frac{q^2 R_{g,i-1}^2}{3}} B_i \left[\frac{\left(\operatorname{erf} \left\{ \frac{q R_{g,i}}{\sqrt{6}} \right\} \right)^3}{q} \right]^{P_i}$$

$R_{g,i-1}$ is the high- q power law cutoff for each level and is taken to be the radius of gyration of the previous level to avoid double counting.



$$I(q) = B_{bkg} + \sum_{i=1}^N G_i e^{-\frac{q^2 R_{g,i}^2}{3}} + e^{-\frac{q^2 R_{g,i-1}^2}{3}} B_i \left[\frac{\left(\operatorname{erf} \left\{ \frac{q R_{g,i}}{\sqrt{6}} \right\} \right)^3}{q} \right]^{P_i}$$

$R_{g,i-1}$ is the high- q power law cutoff for each level and is taken to be the radius of gyration of the previous level to avoid double counting.

Additional parameters can be added, such as a structure factor that is different than unity, and interparticle correlation parameters.



$$I(q) = B_{bkg} + \sum_{i=1}^N G_i e^{-\frac{q^2 R_{g,i}^2}{3}} + e^{-\frac{q^2 R_{g,i-1}^2}{3}} B_i \left[\frac{\left(\operatorname{erf} \left\{ \frac{q R_{g,i}}{\sqrt{6}} \right\} \right)^3}{q} \right]^{P_i}$$

$R_{g,i-1}$ is the high- q power law cutoff for each level and is taken to be the radius of gyration of the previous level to avoid double counting.

Additional parameters can be added, such as a structure factor that is different than unity, and interparticle correlation parameters.

It is important not to include more levels than are significant physically