

# Today's outline - September 09, 2021



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Problems on Blackboard  
due Tuesday, September 21, 2021

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Reading Assignment: Chapter 3.4

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1. Knowing that the photoelectric absorption of an element scales as the inverse of the energy cubed, calculate:
  - (a) the absorption coefficient at 10keV for copper when the value at 5keV is  $1698.3 \text{ cm}^{-1}$ ;
  - (b) The actual absorption coefficient of copper at 10keV is  $1942.1 \text{ cm}^{-1}$ , why is this so different than your calculated value?
2. A 30 cm long, ionization chamber, filled with 80% helium and 20% nitrogen gases at 1 atmosphere, is being used to measure the photon rate (photons/sec) in a synchrotron beamline at 12 keV. If a current of 10 nA is measured, what is the photon flux entering the ionization chamber?
3. A 5 cm deep ionization chamber is used to measure the fluorescence from a sample containing arsenic (As). Using any noble gases or nitrogen, determine a gas fill (at 1 atmosphere) for this chamber which absorbs at least 60% of the incident photons. How does this change if you are measuring the fluorescence from ruthenium (Ru)?



4. Calculate the critical angle of reflection of 10 keV and 30 keV x-rays for:

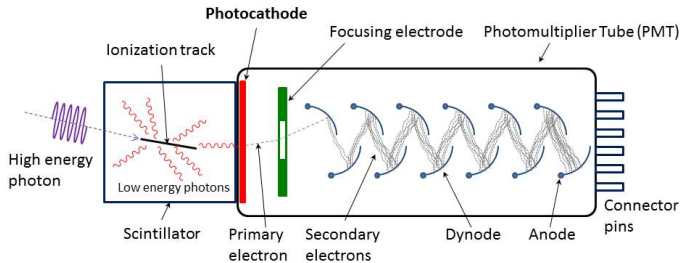
- (a) A slab of glass ( $\text{SiO}_2$ );
- (b) A thick chromium mirror;
- (c) A thick platinum mirror.
- (d) If the incident x-ray beam is 2 mm high, what length of mirror is required to reflect the entire beam for each material?

5. Calculate the fraction of silver (Ag) fluorescence x-rays which are absorbed in a 1 mm thick silicon (Si) detector and the charge pulse expected for each absorbed photon. Repeat the calculation for a 1 mm thick germanium (Ge) detector.

# Scintillation detector



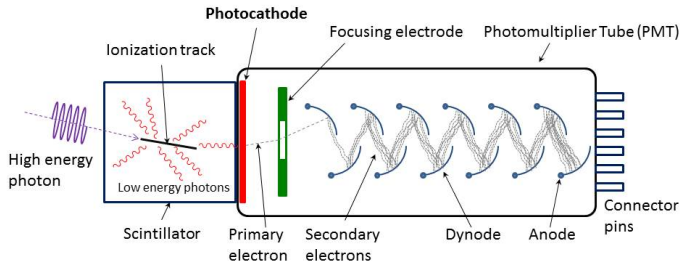
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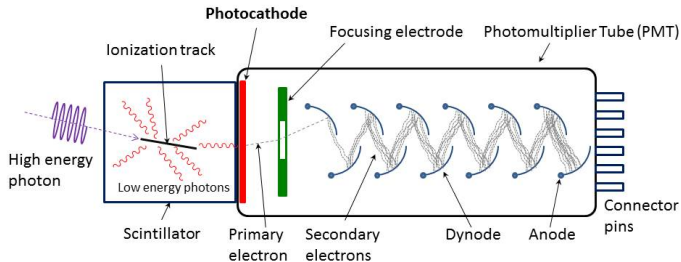


- NaI(Tl), Yttrium Aluminum Perovskite (YAP) or plastic which, absorb x-rays and fluoresce in the visible spectrum.

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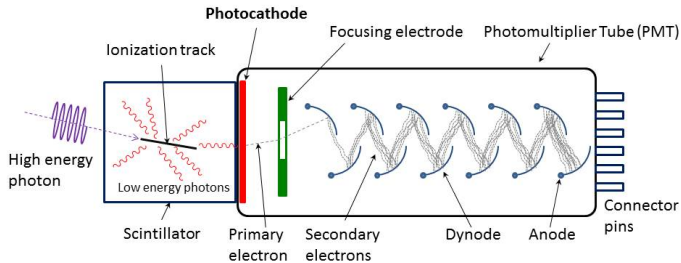


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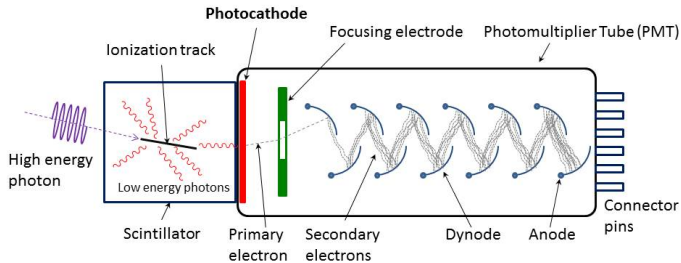


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- Output voltage pulse is proportional to initial x-ray energy.

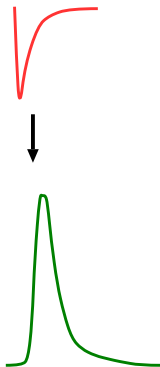


# Counting a scintillator pulse



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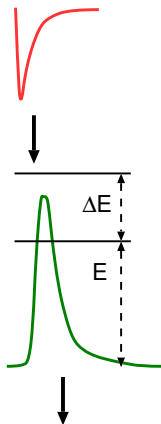
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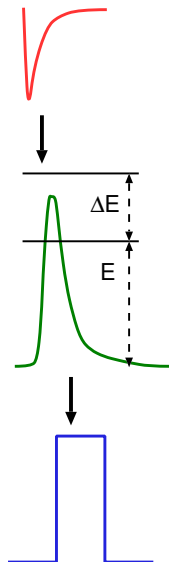


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if the **voltage pulse** falls within the discriminator window, a short **digital pulse** is output from the discriminator and into a scaler for counting

# Solid state detectors



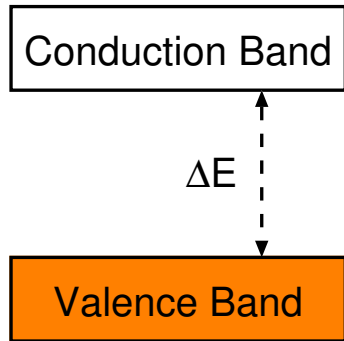
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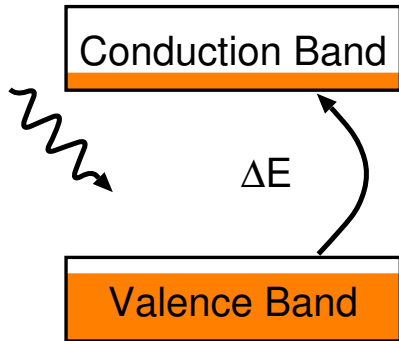
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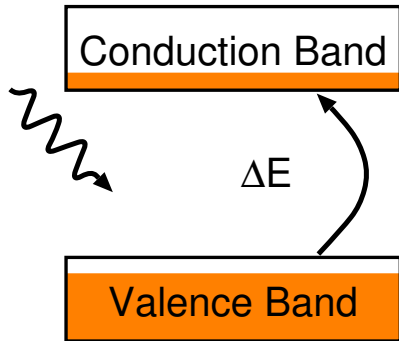
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because of the small energy required to produce an electron-hole pair, one x-ray photon will create many and its energy can be detected with very high resolution



# Semiconductor junctions



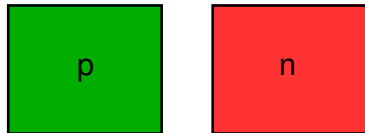
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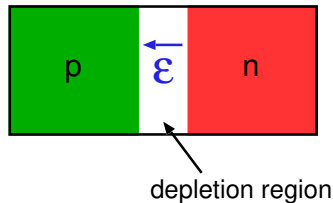
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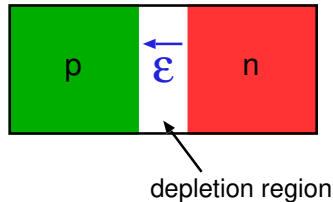


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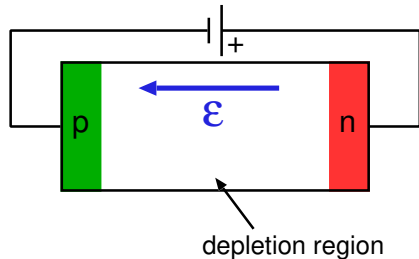
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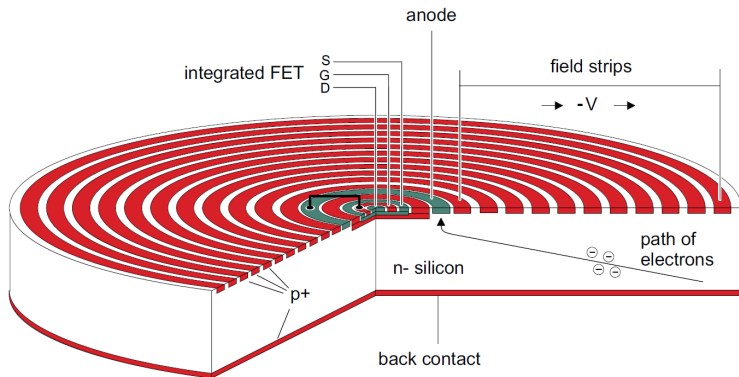
by applying a reverse bias voltage, it is possible to extend the depleted region, make the effective volume of the detector larger and increase the **electric field** to get faster charge collection times



# Silicon Drift Detector

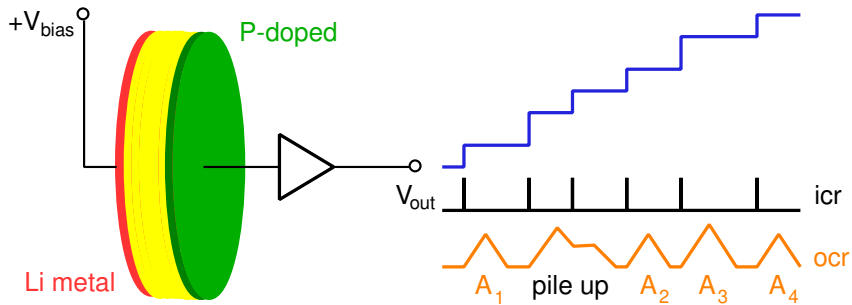


Same principle as intrinsic or p-i-n detector but much more compact and operates at higher temperatures

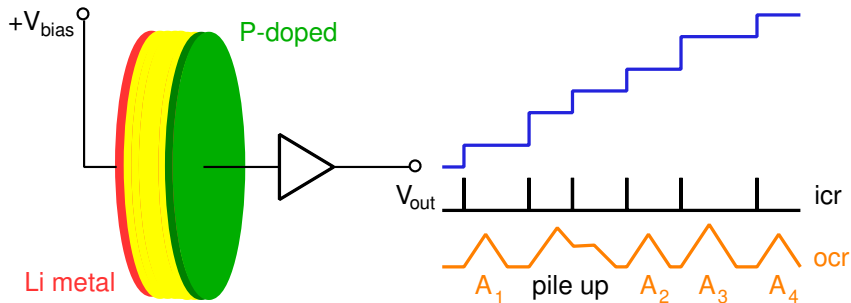


Relatively low stopping power is a drawback

# Detector operation



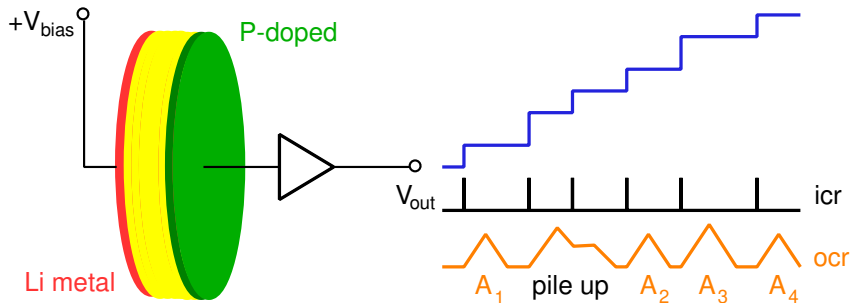
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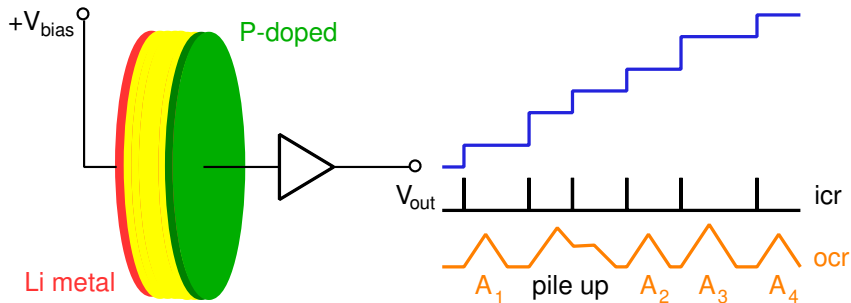
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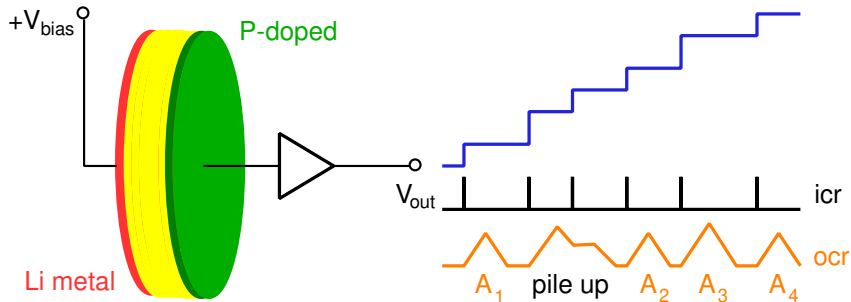
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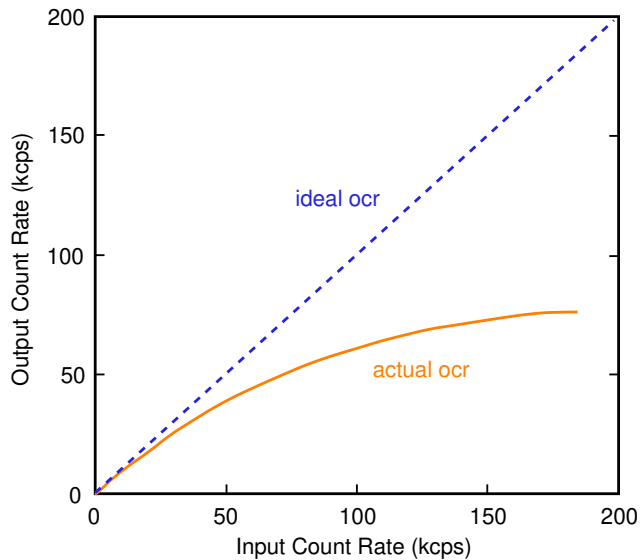


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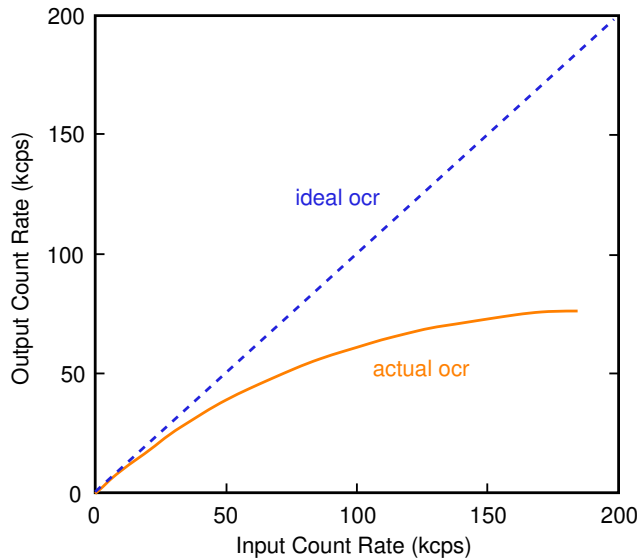
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electronics outputs input count rate (icr), output count rate (ocr), and areas of integrated pulses ( $A_n$ )

# Dead time correction

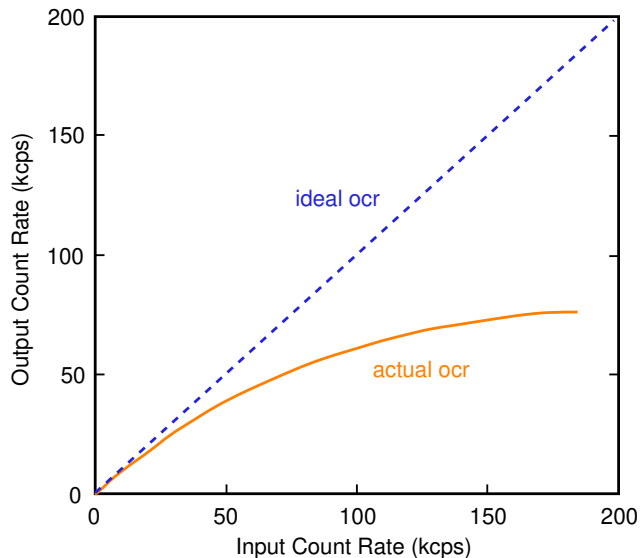


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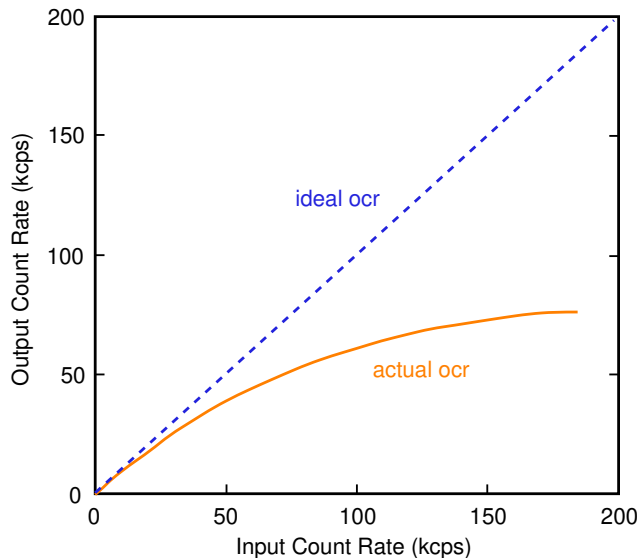
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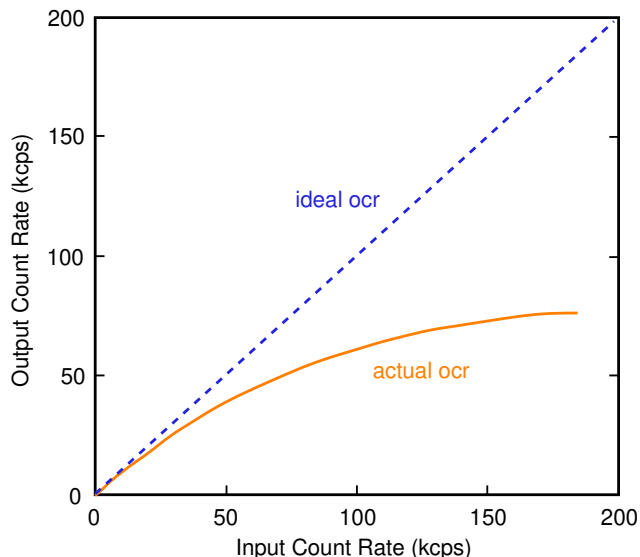


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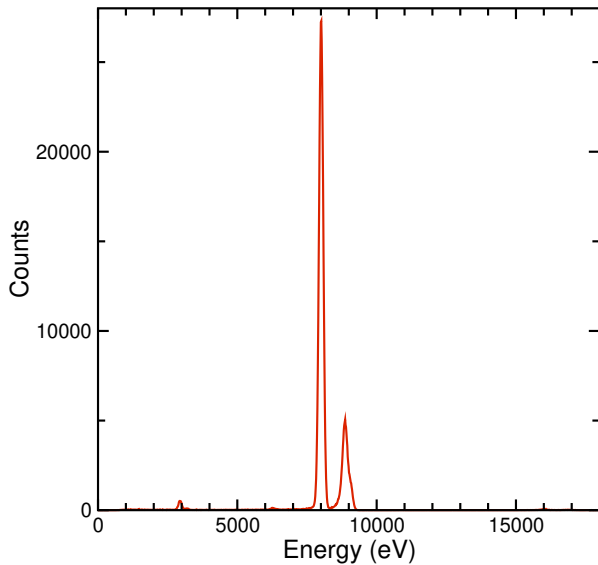
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When dead time is too large, correction will not be accurate!

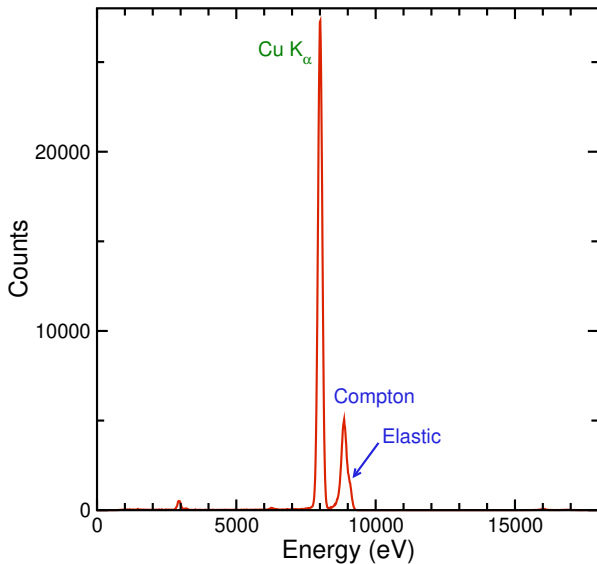


# SDD spectrum



Fluorescence spectrum of Cu foil in air  
using 9200 eV x-rays

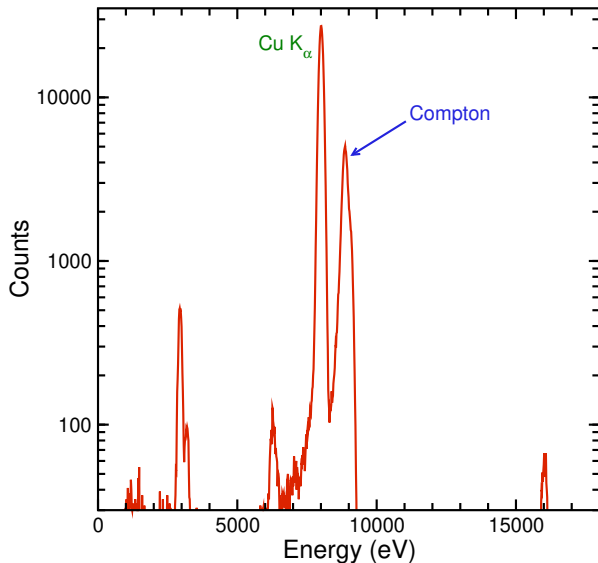
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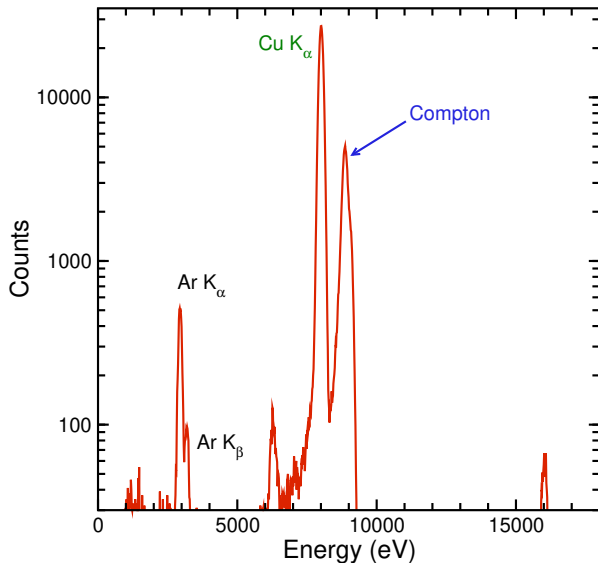
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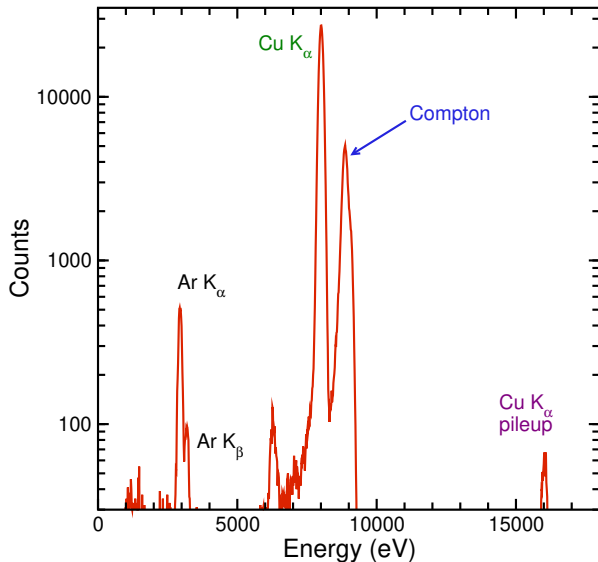


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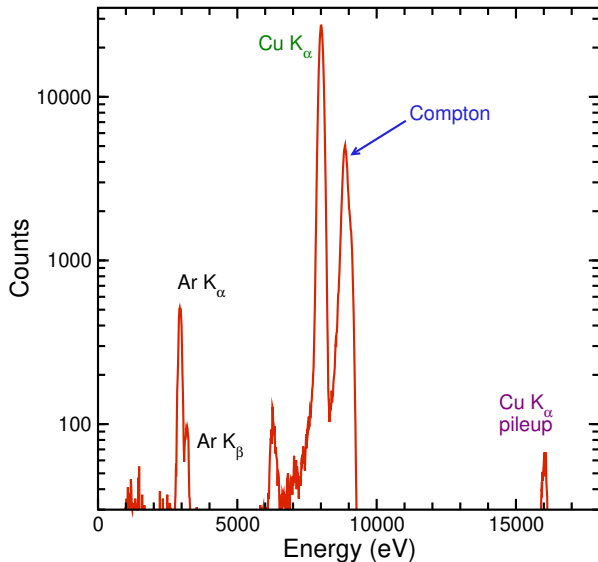
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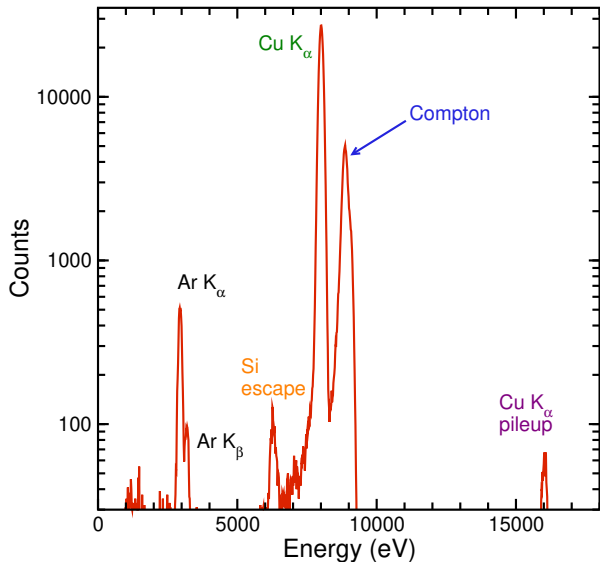
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Si escape peak

$$E_{esc} = 8046 - 1839 = 6207 \text{ eV}$$

# Area detectors



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The most advanced detectors can easily cost over a million dollars!

# CCD detectors - direct



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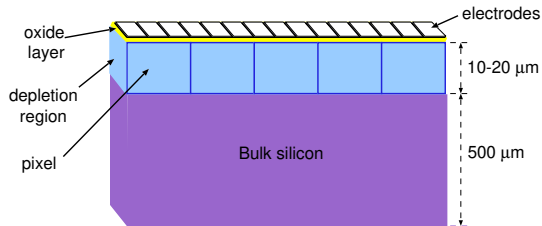


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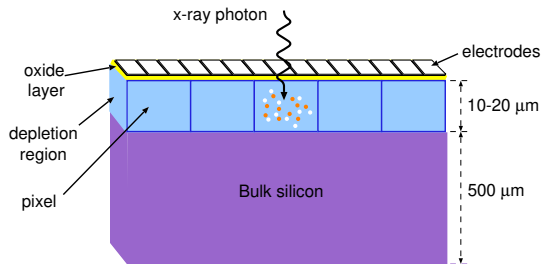
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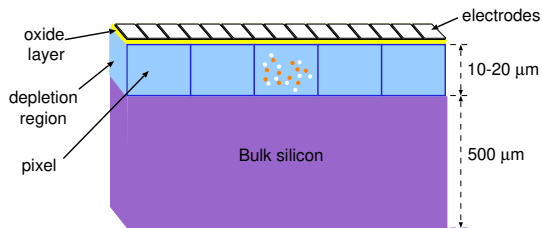
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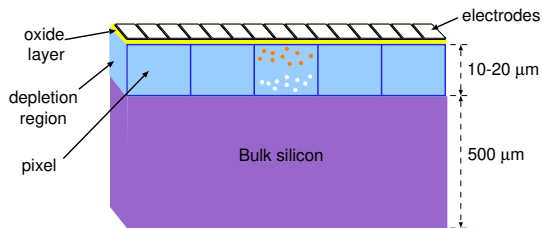
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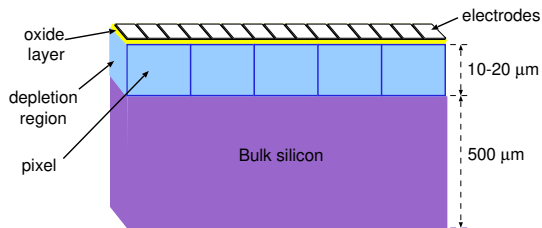
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expensive to make very large, limited sensitivity to high energies

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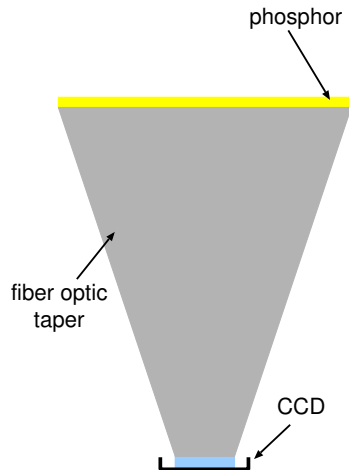
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The CCD is coupled optically to a fiber optic taper which ends at a large phosphor



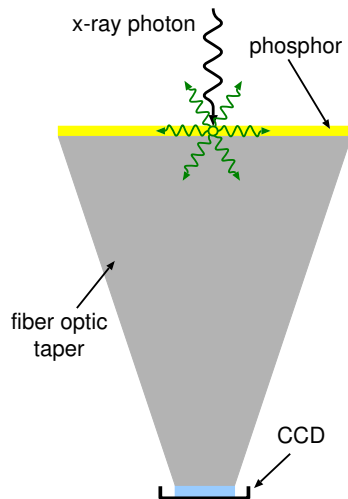
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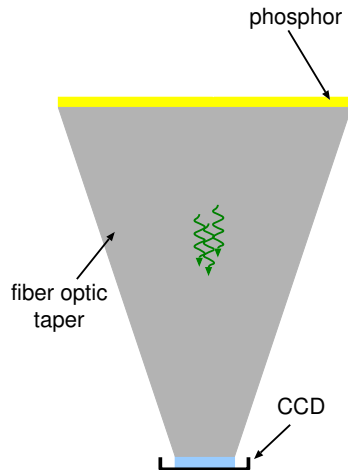


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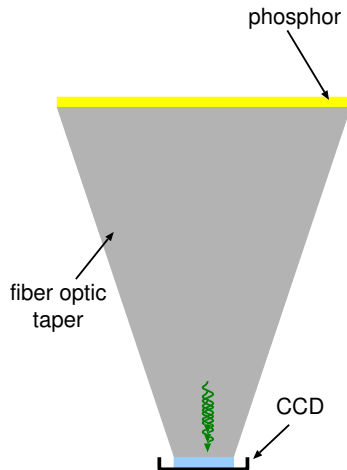
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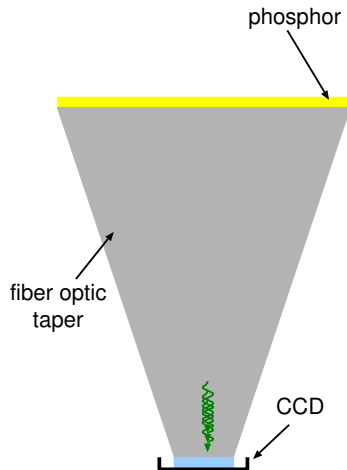
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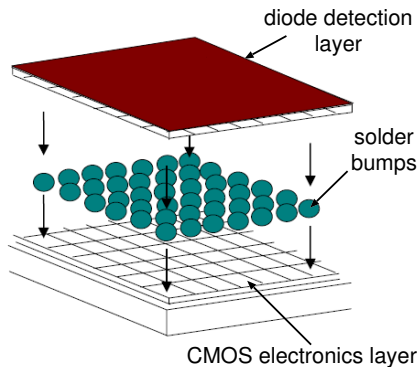


# Pixel array detectors - schematic



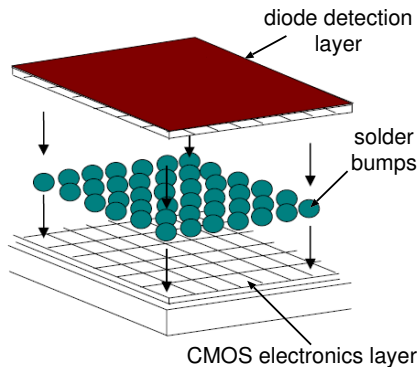
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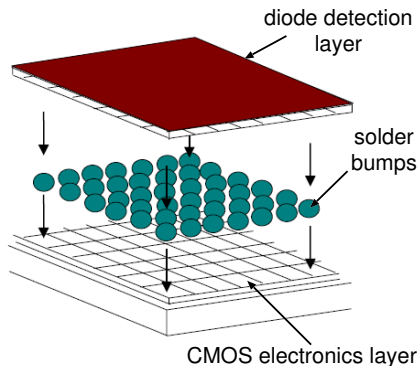
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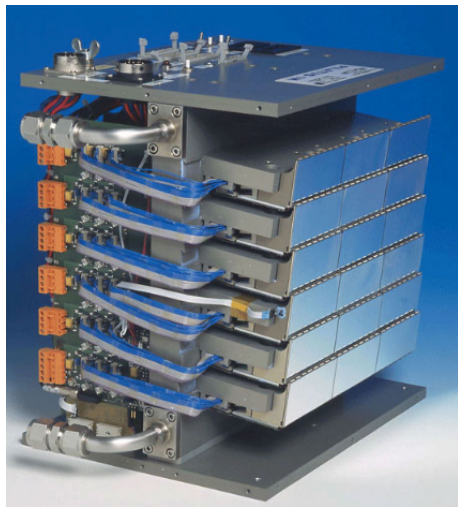


The Pixel Array Detector combines area detection with on-board electronics for fast signal processing

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This permits fast processing and possibly energy discrimination on a per-pixel level

# Pixel array detectors - Pilatus



Pixel array detector with 1,000,000 pixels.

Each pixel has energy resolving capabilities & high speed readout.

Silicon sensor limits energy range of operation.

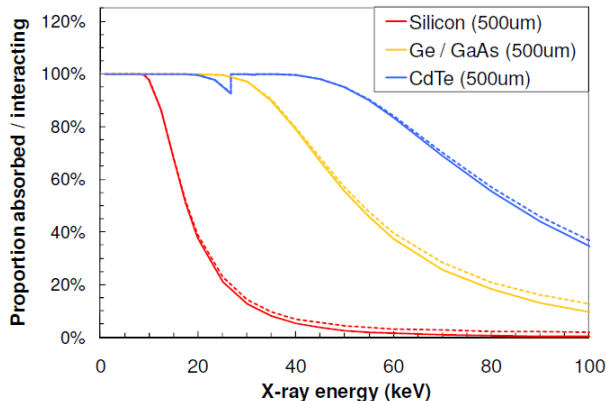
from Swiss Light Source





One of the major problems with pixel array detectors and SDDs is the low absorption cross section at high energies

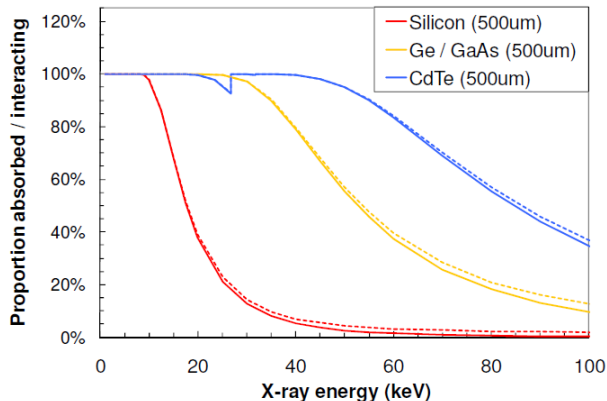
# High energy solutions



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The absorption can be significantly enhanced with these higher Z elements while maintaining good energy discrimination capabilities.

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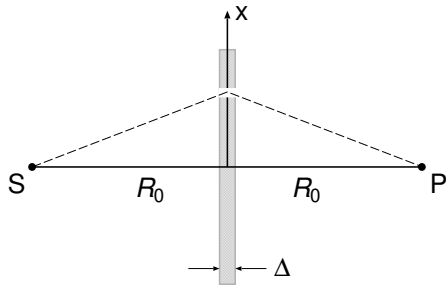
Initially assume that all interfaces are perfectly flat and ignore all absorption processes.



## Thin plate response - scattering approach



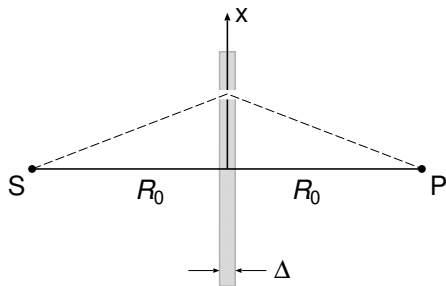
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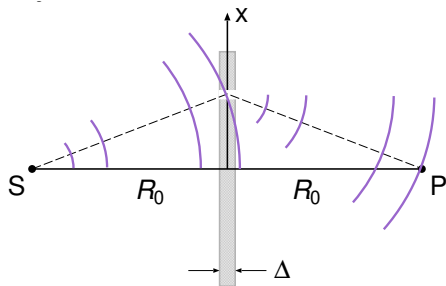
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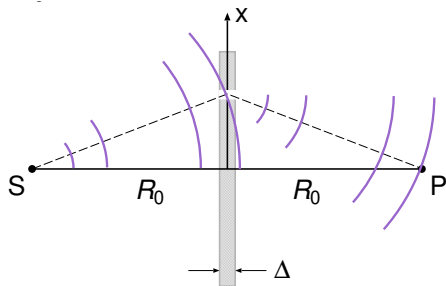
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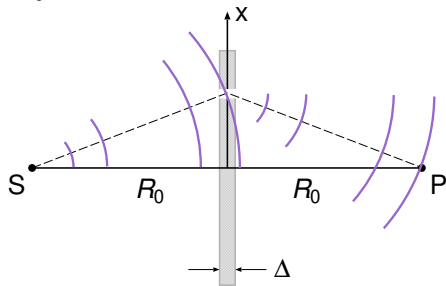


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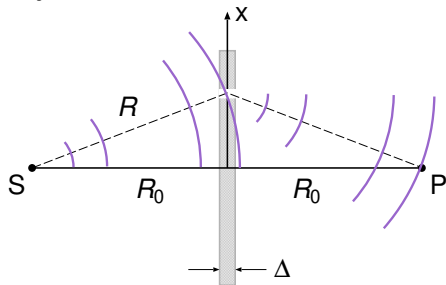


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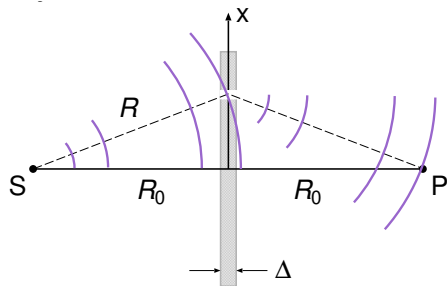
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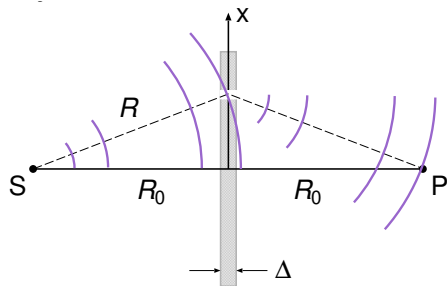
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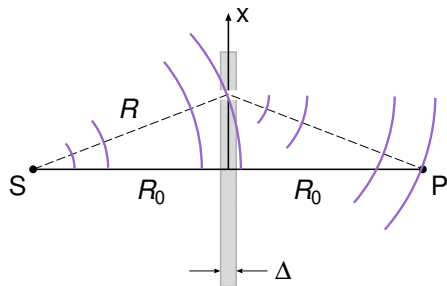
$$R = R_0 \sqrt{1 + \frac{x^2 + y^2}{R_0^2}} \approx R_0 \left[ 1 + \frac{x^2 + y^2}{2R_0^2} \right]$$



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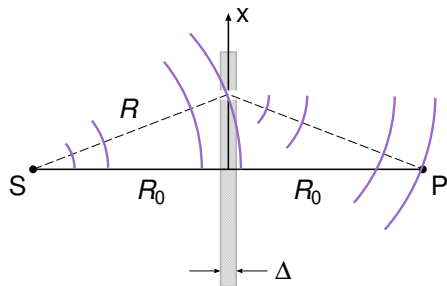
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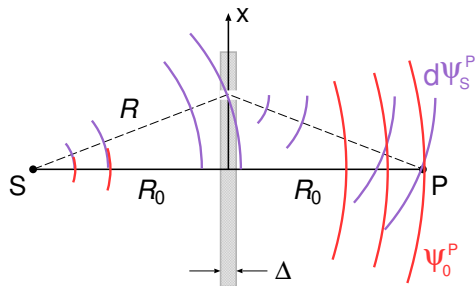


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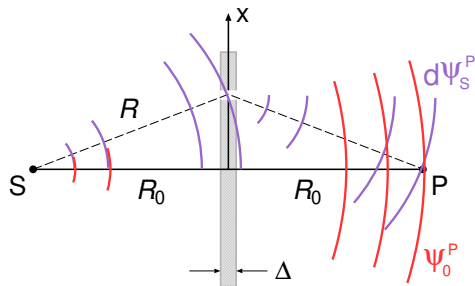
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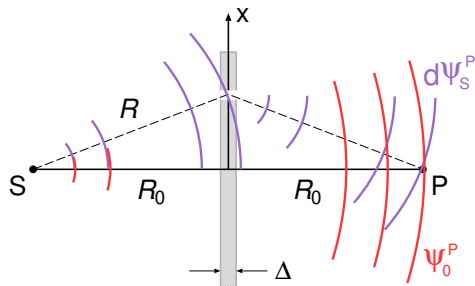
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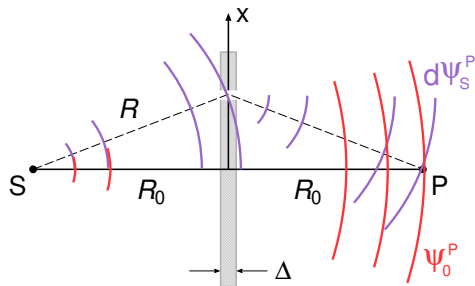
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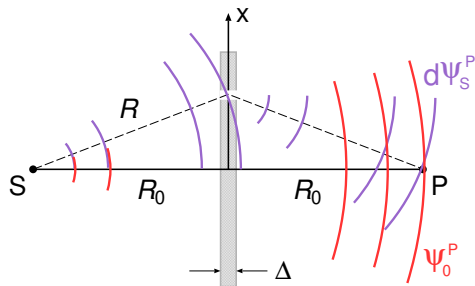
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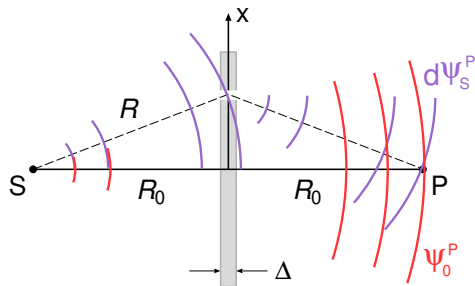
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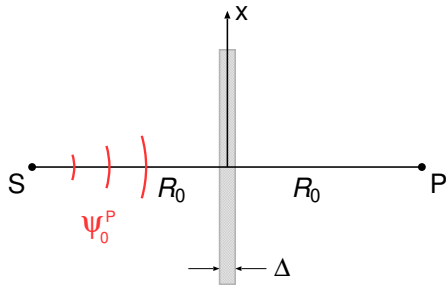
Thus the total wave (electric field) at  $P$  can be written

$$\psi^P = \psi_0^P + \psi_S^P = \frac{e^{i2kR_0}}{2R_0} - i\rho b \Delta \frac{\pi R_0}{k} \frac{e^{i2kR_0}}{R_0^2} = \psi_0^P \left[ 1 - i \frac{2\pi \rho b \Delta}{k} \right]$$

# Thin plate response - refraction approach



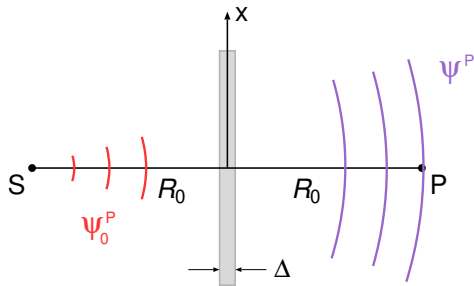
Now let's look at this phenomenon from a different point of view, that of refraction.



# Thin plate response - refraction approach



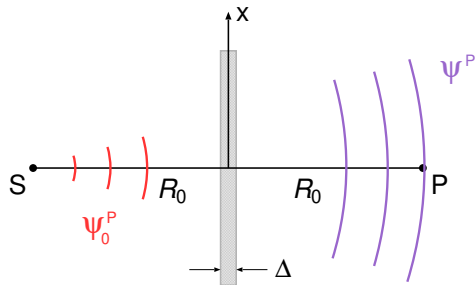
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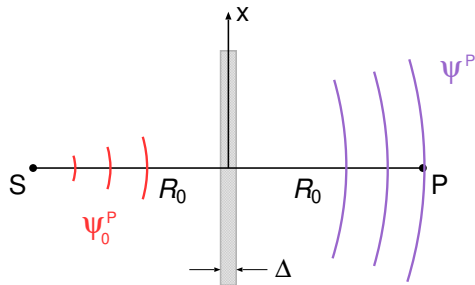


The phase shift depends on the thickness and the difference between the index of refraction of the medium and that of vacuum

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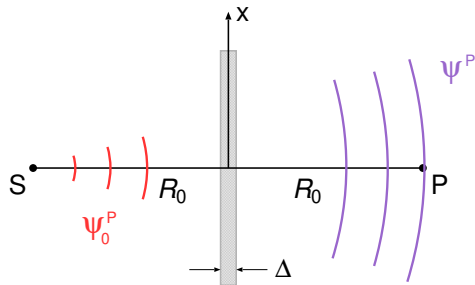
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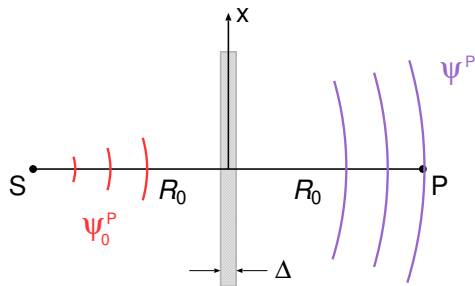
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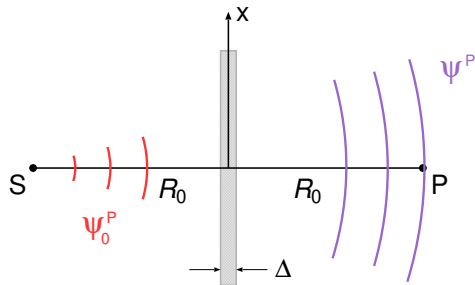
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$$\begin{aligned}\phi &= 2\pi \left( \frac{n\Delta}{\lambda} - \frac{\Delta}{\lambda} \right) \\ &= \frac{2\pi}{\lambda} \Delta (n - 1)\end{aligned}$$

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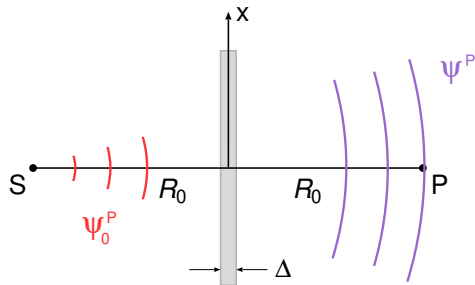
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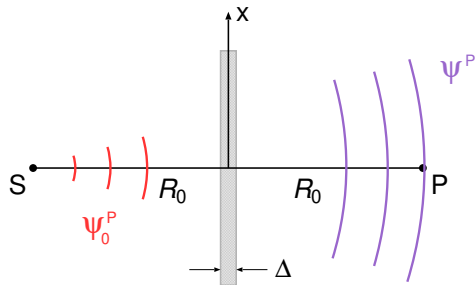
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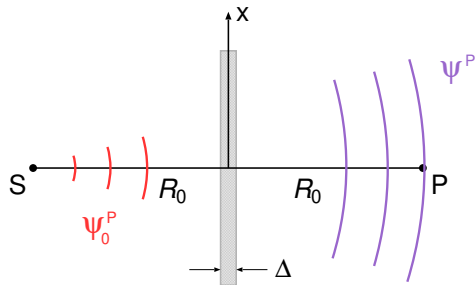
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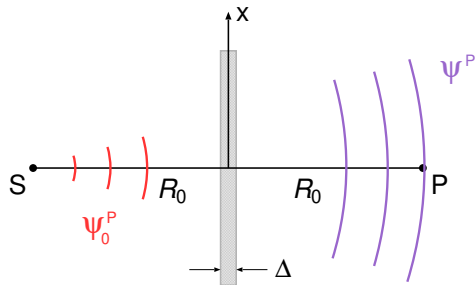
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The wave function at  $P$  is then:

$$\psi^P = \psi_0^P e^{i(n-1)k\Delta} = \psi_0^P [1 + i(n-1)k\Delta + \dots] \approx \psi_0^P [1 + i(n-1)k\Delta]$$

## Calculating $n$



We can now compare the expressions obtained by the scattering and refraction approaches.

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