



• Undulator spectrum

V

- Undulator spectrum
- Undulator coherence

V

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- APS-U, ERLs and FELs



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Reading Assignment: Chapter 3.1–3.3



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Reading Assignment: Chapter 3.1–3.3 Homework Assignment #02: Problems on Blackboard

due Tuesday, September 21, 2021

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Reading Assignment: Chapter 3.1–3.3

Homework Assignment #02: Problems on Blackboard due Tuesday, September 21, 2021 Homework Assignment #03: Chapter 3: 1,3,4,6,8 due Tuesday, October 05, 2021



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Now let us look at additional properties

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$$= 0.934 \lambda_u [\text{cm}] B_0[\text{T}]$$

$$S\lambda_u \approx \lambda_u \left(1 + \frac{1}{4} \frac{\kappa^2}{\gamma^2}\right)$$

$$\lambda_1 pprox rac{\lambda_u}{2\gamma^2} \left(1 + rac{\kappa^2}{2}
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This can be rewritten in terms of the coordinates in the figure using the vector of unit length in the observer direction:

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1

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 $\omega_1 \gg \omega_u$ as expected because of the Doppler compression



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The motion of the electron, $\sin \omega_u t'$, is always sinusoidal, but because of the additional terms, the motion as seen by the observer, $\sin \omega_1 t$, is not.

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$$\omega_1 t = \omega_u t' - \frac{\kappa^2/4}{1 + (\gamma \theta)^2 + \kappa^2/2} \sin(2\omega_u t') \\ - \frac{2\kappa\gamma}{1 + (\gamma \theta)^2 + \kappa^2/2} \phi \sin(\omega_u t')$$



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Suppose we have K = 1 and $\theta = 0$ (on axis), then



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Suppose we have $\mathcal{K}=1$ and $\theta=0$ (on axis), then

$$\omega_1 t = \omega_u t' + \frac{1}{6} \sin\left(2\omega_u t'\right)$$



On-axis undulator characteristics

$$\omega_1 t = \omega_u t' - \frac{K^2/4}{1 + (\gamma \theta)^2 + K^2/2} \sin(2\omega_u t') - \frac{2K\gamma}{1 \pm (\gamma \theta)^2 + K^2/2} \phi \sin(\omega_u t')$$

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Similarly, for K = 2



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Similarly, for K = 2 and K = 5, the deviation becomes more pronounced.



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Similarly, for K = 2 and K = 5, the deviation becomes more pronounced. This shows how higher harmonics must be present in the radiation as seen by the observer.



Intensity (arb. units)



$$\omega_1 t = \omega_u t' - \frac{K^2/4}{1 + (\gamma \theta)^2 + K^2/2} \sin(2\omega_u t') - \frac{2K\gamma}{1 + (\gamma \theta)^2 + K^2/2} \phi \sin(\omega_u t')$$

When K=2 and $heta=\phi=1/\gamma$, we have





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Off-axis undulator characteristics

$$\omega_1 t = \omega_u t' - \frac{\kappa^2/4}{1 + (\gamma \theta)^2 + \kappa^2/2} \sin(2\omega_u t') - \frac{2\kappa\gamma}{1 + (\gamma \theta)^2 + \kappa^2/2} \phi \sin(\omega_u t')$$

When K=2 and $\theta=\phi=1/\gamma$, we have

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The last term introduces an antisymmetric term which skews the function and leads to the presence of forbidden harmonics $(2^{nd}, 4^{th}, \text{ etc})$ in the radiation from the undulator compared to the on-axis radiation.







An N period undulator is basically like a diffraction grating, only in the time domain rather than the space domain.



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$$\sum_{m=0}^{N-1} e^{i(\vec{k}\cdot\vec{r}+2\pi m\epsilon)}$$



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the phase shift from each undulator pole depends on the wavelength $\lambda_{\rm u}$

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$$S_N = 1 + kS_{N-1}$$



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The sum is simply a geometric series, S_N with $k = e^{i2\pi\epsilon}$

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$$S_N - kS_N = 1 - k^N \quad \longrightarrow \quad S_N = \frac{1 - k^N}{1 - k}$$

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Intensity from a diffraction grating

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The coherence of an undulator depends on the amount each pole's emission is out of phase with the others, ϵ .

2πε=0



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30

40















































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With the height and width of the peak dependent on the number of poles.









The more poles in the undulator, the more monochromatic the beam since a slight change in $\epsilon = \delta L/\lambda$ implies a slightly different wavelength λ



Intensity (arb units)



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Higher order harmonics have narrower energy bandwidth but lower peak intensity

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Synchrotron time structure





There are two important time scales for a storage ring such as the APS: pulse length and interpulse spacing

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The APS pulse length in 24-bunch mode is 90 ps while the pulses come every 154 ns

Other modes include single-bunch mode for timing experiments and 324-bunch mode (inter pulse timing of 11.7 ns) for a more constant x-ray flux









• Bending magnet

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• Bending magnet

• Broad, nearly white spectrum





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 - Energies extend to 100's of keV





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• Undulator

• Brilliance is 6 orders larger than a bending magnet





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- Both odd and even harmonics appear
Spectral comparison





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Undulator

- Brilliance is 6 orders larger than a bending magnet
- Both odd and even harmonics appear
- Harmonics can be tuned in energy (dashed lines)

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Is there a limit to the brightness of an undulator source at a synchrotron?



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the brightness is inversely proportional to the square of the product of the linear source size and the angular divergence

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this emittance cannot be changed but it can be rotated or deformed by magnetic fields as the electron beam travels around the storage ring as long as the area is kept constant





For photon emission from a single electron in a 2m undulator at $1 \mbox{\AA}$

V

APS emittance

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$$\sigma_{\it radiation} = 9.1 \mu {
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 $\sigma'_{\it radiation} = 7.7 \mu {
m rad}$

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|--|

Paramotor	APS			
rarameter	1995	2001	2005	
Bunches		24 & 324		
σ_{x}	334 μ m	352 μ m	280 μ m	
σ'_{x}	24 μ rad	22 μ rad	11.6 μ rad	
σ_y	89 μ m	18.4 μ m	9.1 μ m	
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Daramatar	APS		AF	APS-U	
Parameter	1995	2001	2005	Timing	Brightness
Bunches		24 & 324		48	324
σ_{x}	334 μ m	352 μ m	280 μ m	18.1 μ m	21.8 μ m
σ'_{x}	24 μ rad	22 μ rad	11.6 μ rad	2.6 μ rad	3.1 μ rad
σ_y	89 μ m	18.4 μ m	9.1 μ m	10.6 μ m	4.1 μ m
σ'_{y}	8.9 μ rad	4.2 μ rad	3.0 μ rad	4.2 μ rad	$1.7~\mu$ rad

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By the end of the first decade of operation, the horizontal source size decreased by about 16% and its horizontal divergence by more than 50% while the vertical source size decreased by over 90% and the vertical divergence by nearly 67%

The APS-U will make the beam smaller and more square in space making for higher performance insertion device beam lines.

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The beam will be nearly square and there will be much more coherence from the undulators

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APS-U magnet layout



The APS upgrade will install a multi-bend achromat instead of the two bending magnets.

APS-U magnet layout

The APS upgrade will install a multi-bend achromat instead of the two bending magnets.

two dipole magnets - double-bend-achromat





Seven dipole magnets – multi-bend-achromat (MBA)

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APSU undulator performance



The multi-bend achromat will produce a diffraction-limited source with a lower energy (6.0 GeV) and doubled current (200 mA).

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Since MRCAT's science is primarily flux driven, the goal will be to replace the 2.4m undulator with one that outperforms the current 33mm period but with only modest increase in power.





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The APS-U is an example of a " 4^{th} " generation synchrotron source





Energy recovery linacs

V

Another way to imcrease the undulator peak brilliance is the energy recovery linac which overcomes the fundamental limitation of a synchrotron by passing the beam through only once.



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Initial electron cloud, each electron emits coherently but independently

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Free electron laser





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Over course of 100 m, electric field of photons, feeds back on the electron bunch

Free electron laser





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Microbunches form with period of FEL (and radiation in electron frame)
Free electron laser





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Microbunches form with period of FEL (and radiation in electron frame)

Each microbunch emits coherently with neighboring ones

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Microbunches form with period of FEL (and radiation in electron frame)

Each microbunch emits coherently with neighboring ones

Again, an alternative way to view this is that the pulse train from a 100m long undulator is long enough in time to produce a monochromatic and coherent frequency distribution when Fourier Transformed

FEL performance















An FEL has a single accelerator whose electron beam is shunted sequentially through different undulators and end stations





An FEL has a single accelerator whose electron beam is shunted sequentially through different undulators and end stations

The single pass of the electron beam permits a very low emittance to be achieved and thus higher coherence





An FEL has a single accelerator whose electron beam is shunted sequentially through different undulators and end stations

The single pass of the electron beam permits a very low emittance to be achieved and thus higher coherence

The high brightness usually results in destruction of the sample during the illumination, thus the need for multiple samples and multiple shot experiments

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Small low energy, high current electron ring







Small low energy, high current electron ring

Straight section intersects a laser cavity

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Undulator is the standing wave of the laser, alternatively can consider this an inverse Compton effect





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Low energy x-rays produced are suitable for protein crystallography, SAXS and imaging

Cost ${\sim}\$5$ M plus ${\sim}\$1$ M per year service contract







Gas detectors

• Ionization chamber



- Ionization chamber
- Proportional counter



- Ionization chamber
- Proportional counter
- Geiger-Muller tube



- Ionization chamber
- Proportional counter
- Geiger-Muller tube
- Scintillation counters



- Ionization chamber
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- Geiger-Muller tube Scintillation counters Solid state detectors



Gas detectors

- Ionization chamber
- Proportional counter
- Geiger-Muller tube

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• Intrinsic semiconductor



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Indirect



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- Indirect
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The most useful regime is the ionization region where the output pulse is independent of the applied voltage over a wide range









Useful for beam monitoring, flux measurement, fluorescence measurement, spectroscopy.



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- 22-41 eV per electron-ion pair (depending on the gas) makes this useful for quantitative measurements.


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the digital pulse train is counted by a scaler for a user-definable length of time