

Today's outline - August 31, 2021



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- Refraction and reflection of x-rays

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- Coherence of x-ray sources

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- The bending magnet source
 - Segmented arc approximation
 - Off-axis emission

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Reading Assignment: Chapter 2.3–2.4

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Reading Assignment: Chapter 2.3–2.4

Homework Assignment #01:

Chapter 2: 2,3,5,6,8

due Tuesday, September 07, 2021

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Reading Assignment: Chapter 2.3–2.4

Homework Assignment #01:
Chapter 2: 2,3,5,6,8
due Tuesday, September 07, 2021

Homework Assignment #02:
Problems on Blackboard
due Tuesday, September 21, 2021

Refraction & reflection of x-rays



X-rays can be treated like light when interaction with a medium. However, unlike visible light, the index of refraction of x-rays in matter is very close to unity:

Refraction & reflection of x-rays



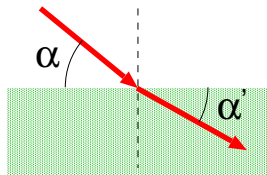
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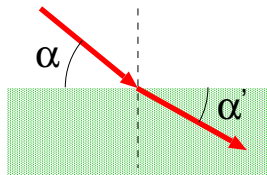


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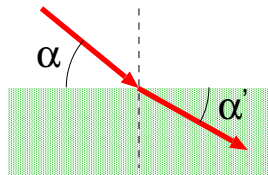
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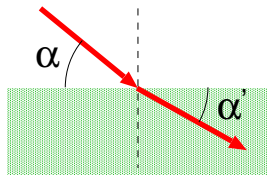
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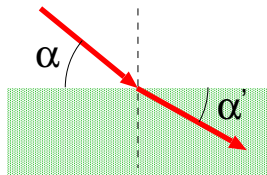
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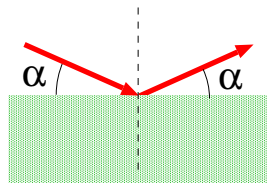


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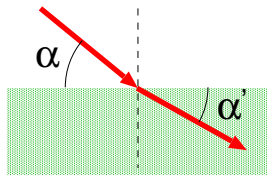
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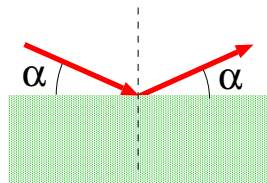


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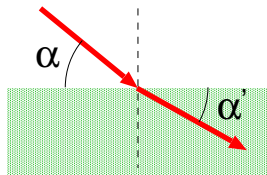


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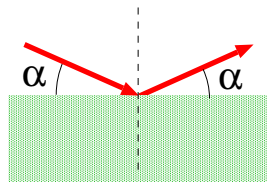


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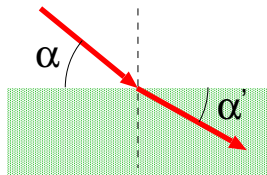
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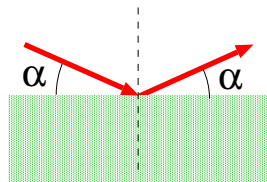


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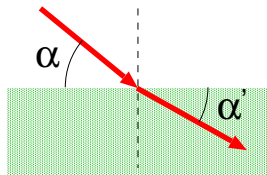
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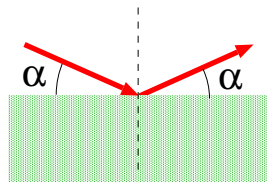


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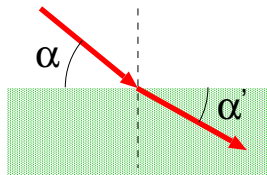
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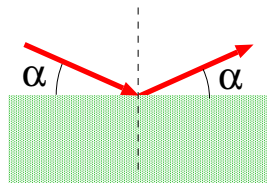


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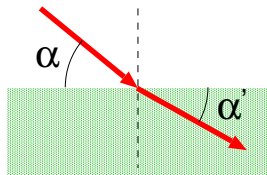
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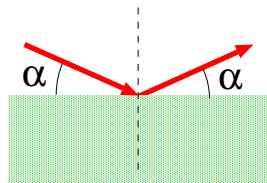


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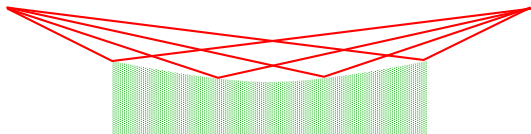


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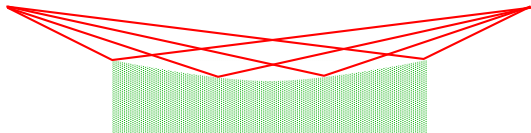
$$\delta = \frac{\alpha_c^2}{2} \quad \longrightarrow \quad \alpha_c = \sqrt{2\delta}$$

Uses of total external reflection



X-ray mirrors

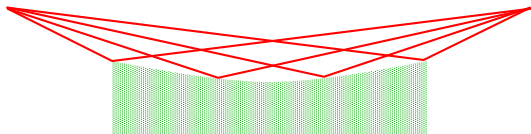
Uses of total external reflection



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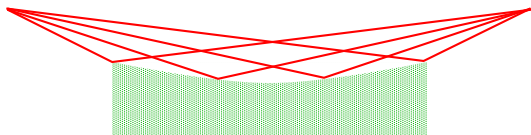
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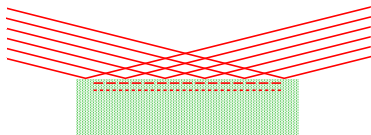
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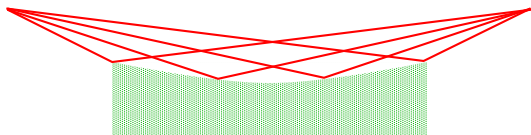
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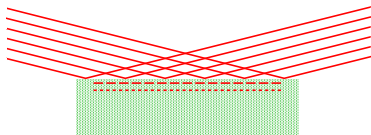
Evanscent wave experiments

Uses of total external reflection



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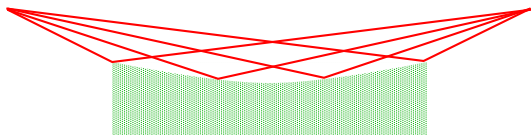
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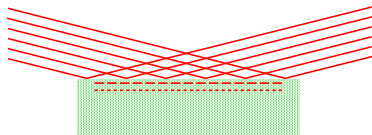
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Uses of total external reflection



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Evanscent wave experiments

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Magnetic interactions



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Coherence: what is it?



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Because of these imperfections the “coherence length” of an x-ray beam is finite and we can calculate it.

Longitudinal coherence

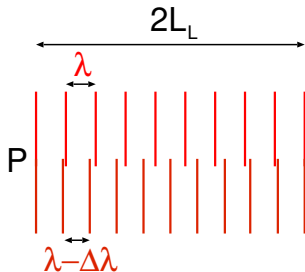


Definition: *Distance over which two waves from the same source point with slightly different wavelengths will completely dephase.*

Longitudinal coherence



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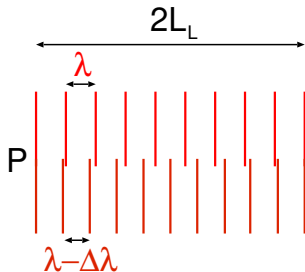


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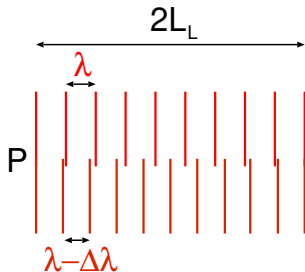
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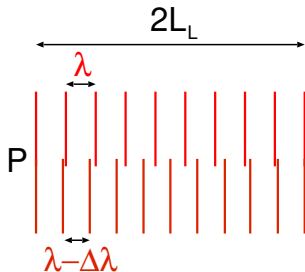
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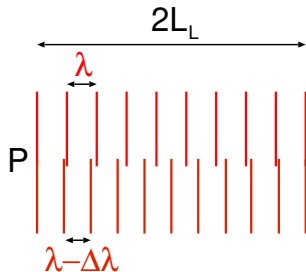
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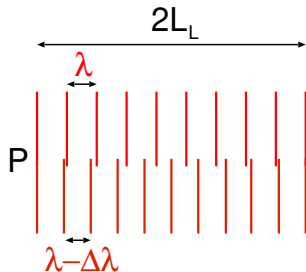
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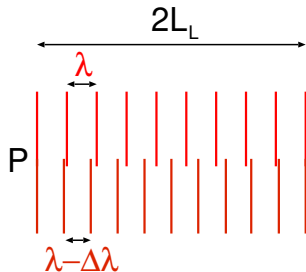
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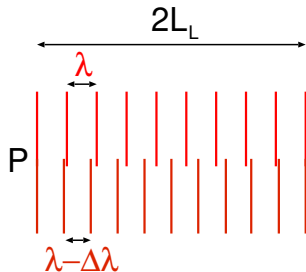
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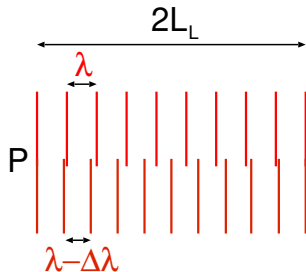
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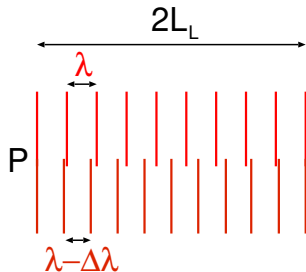
$$\cancel{N\lambda} = \cancel{N\lambda} + \lambda - N\Delta\lambda - \Delta\lambda$$

$$0 = \lambda - N\Delta\lambda - \Delta\lambda \rightarrow \lambda = (N+1)\Delta\lambda \rightarrow N \approx \frac{\lambda}{\Delta\lambda}$$

Longitudinal coherence



Definition: Distance over which two waves from the same source point with slightly different wavelengths will completely dephase.



Two waves of slightly different wavelengths λ and $\lambda - \Delta\lambda$ are emitted from the same point in space simultaneously.

After a distance L_L , the two waves will be exactly out of phase and after $2L_L$ they will once again be in phase.

$$2L_L = N\lambda$$

$$2L_L = (N+1)(\lambda - \Delta\lambda)$$

$$\cancel{N\lambda} = \cancel{N\lambda} + \lambda - N\Delta\lambda - \Delta\lambda$$

$$0 = \lambda - N\Delta\lambda - \Delta\lambda \rightarrow \lambda = (N+1)\Delta\lambda \rightarrow N \approx \frac{\lambda}{\Delta\lambda} \rightarrow L_L = \frac{\lambda^2}{2\Delta\lambda}$$

Transverse coherence

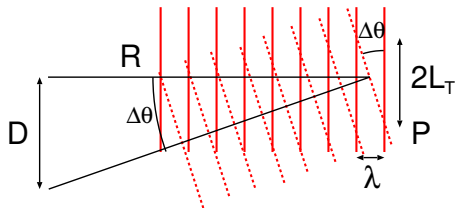


Definition: *The lateral distance along a wavefront over which there is a complete dephasing between two waves, of the same wavelength, which originate from two separate points in space.*

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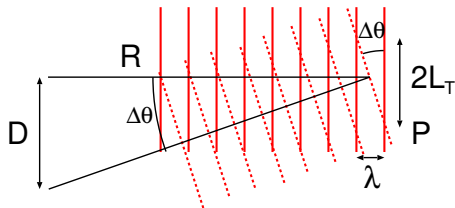


If we assume that the two waves originate from points with a small angular separation $\Delta\theta$, The transverse coherence length is given by:

Transverse coherence



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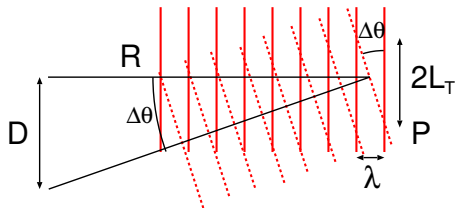
$$\frac{\lambda}{2L_T} = \tan \Delta\theta$$

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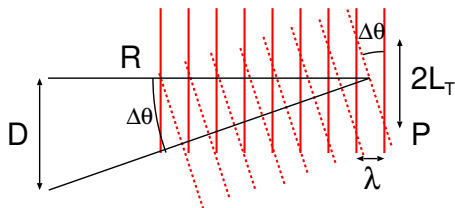
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$$L_T = \frac{\lambda R}{2D}$$

Coherence lengths at the APS



For a typical 3rd generation undulator source, such as at the Advanced Photon Source the vertical source size is $D_v = 10\mu\text{m}$ and we are typically $R = 50\text{m}$ away with our experiment. If we assume a typical wavelength of $\lambda = 1\text{\AA}$, and a monochromator resolution of $\Delta\lambda/\lambda = 10^{-5}$ we have for the vertical direction:

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$$L_L = \frac{\lambda^2}{2\Delta\lambda} = \frac{\lambda}{2} \cdot \frac{\lambda}{\Delta\lambda} = \frac{1 \times 10^{-10}}{2 \cdot 10^{-5}} = 5\mu\text{m}$$

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$$L_T = \frac{\lambda R}{2D} = \frac{(1 \times 10^{-10}) \cdot 50}{2 \cdot (10 \times 10^{-6})}$$

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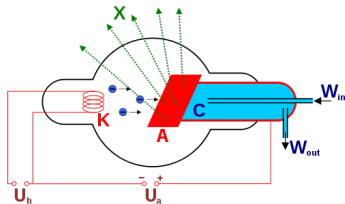
$$L_L = \frac{\lambda^2}{2\Delta\lambda} = \frac{\lambda}{2} \cdot \frac{\lambda}{\Delta\lambda} = \frac{1 \times 10^{-10}}{2 \cdot 10^{-5}} = 5\mu\text{m}$$

$$L_T = \frac{\lambda R}{2D} = \frac{(1 \times 10^{-10}) \cdot 50}{2 \cdot (10 \times 10^{-6})} = 250\mu\text{m}$$

Lab x-ray source schematics



Fixed anode tube

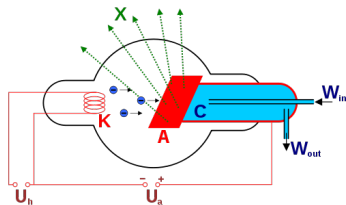


- low power
- low maintenance

Lab x-ray source schematics

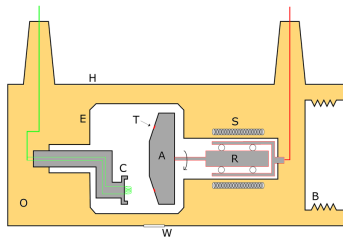


Fixed anode tube



- low power
- low maintenance

Rotating anode tube

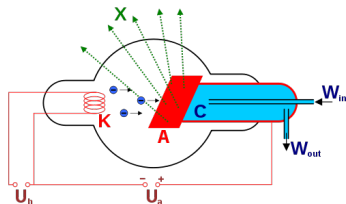


- high power
- high maintenance

Lab x-ray source schematics

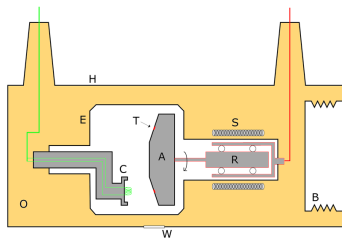


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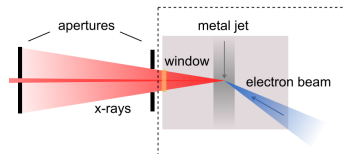
- low power
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Rotating anode tube



- high power
- high maintenance

Liquid metal jet

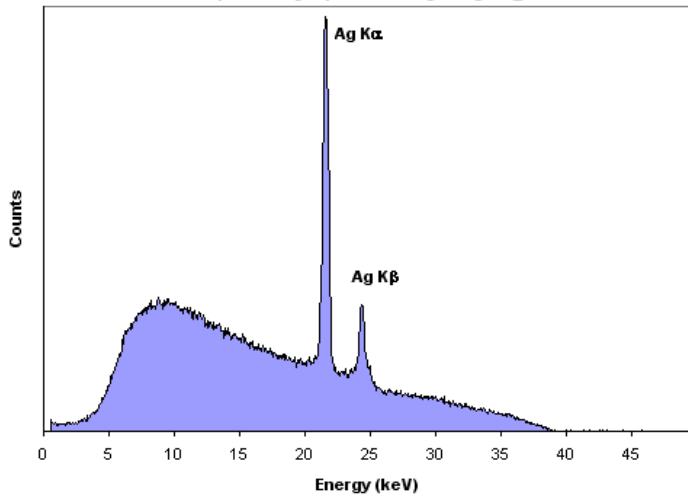


- high brightness
- small spot size

X-ray tube spectrum



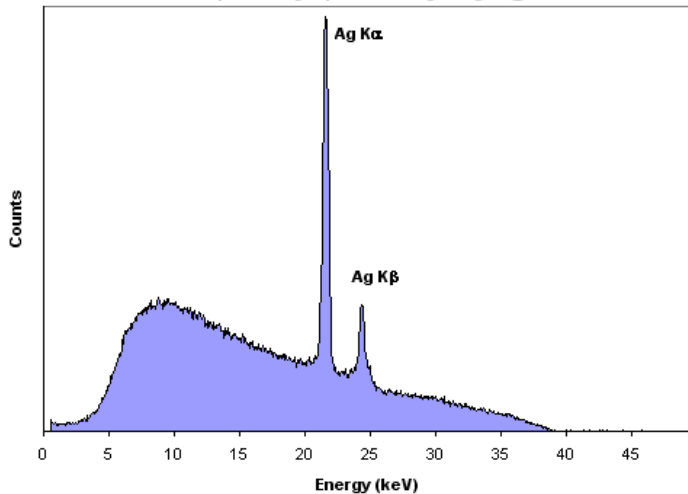
Mini-X Output X-Ray Spectrum: Ag Target @ 40 kV



X-ray tube spectrum



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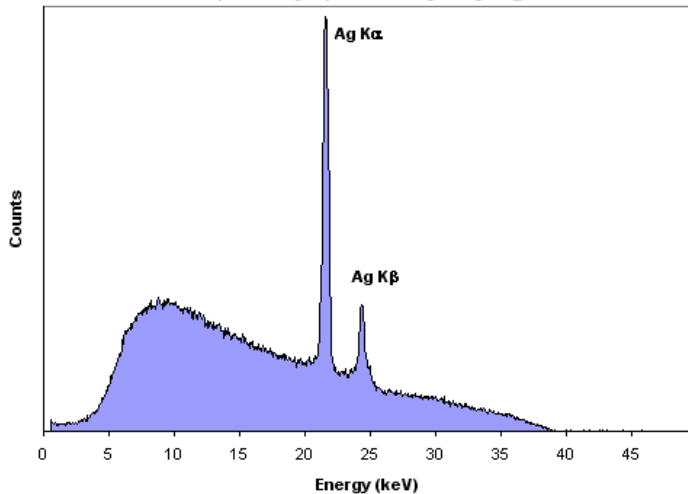


- Minimum wavelength (maximum energy) set by accelerating potential

X-ray tube spectrum



Mini-X Output X-Ray Spectrum: Ag Target @ 40 kV

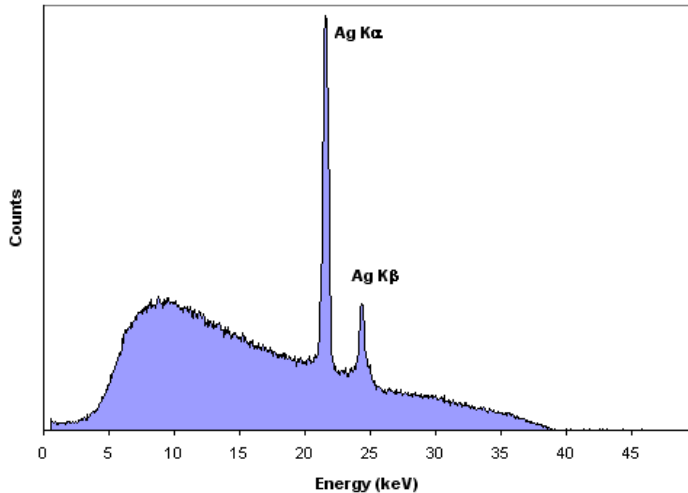


- Minimum wavelength (maximum energy) set by accelerating potential
- Bremsstrahlung radiation provides smooth background (charged particle deceleration)

X-ray tube spectrum



Mini-X Output X-Ray Spectrum: Ag Target @ 40 kV

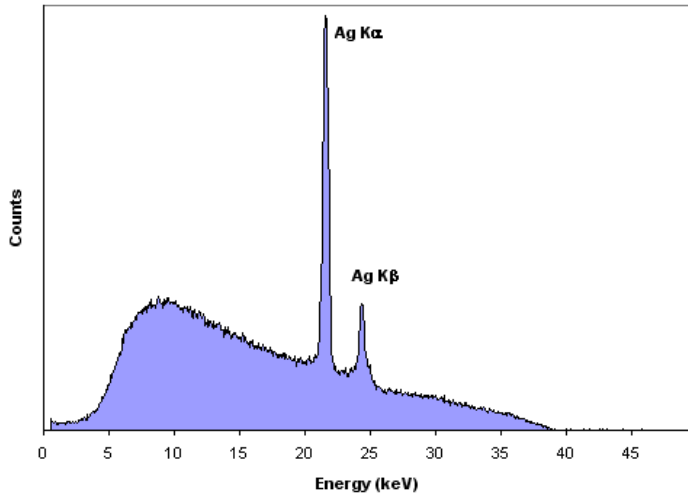


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X-ray tube spectrum



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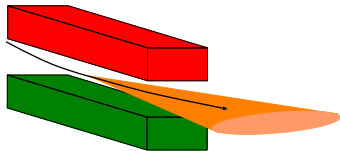


- Minimum wavelength (maximum energy) set by accelerating potential
- Bremsstrahlung radiation provides smooth background (charged particle deceleration)
- Highest intensity at the characteristic fluorescence emission energy of the anode material
- Unpolarized, incoherent x-rays emitted in all directions from anode surface, must be collimated with slits

Synchrotron sources



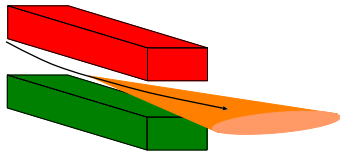
Bending magnet



Synchrotron sources



Bending magnet

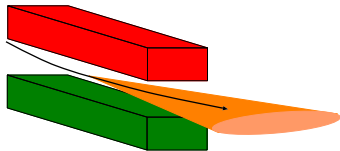


- Wide horizontal beam

Synchrotron sources



Bending magnet

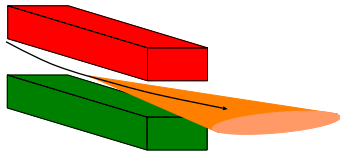


- Wide horizontal beam
- Broad spectrum

Synchrotron sources



Bending magnet

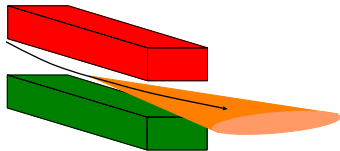


- Wide horizontal beam
- Broad spectrum
- Low brilliance

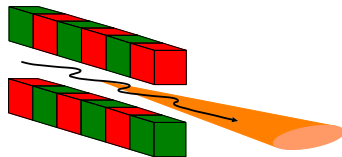
Synchrotron sources



Bending magnet



Wiggler

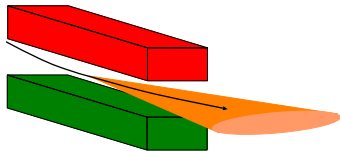


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Synchrotron sources

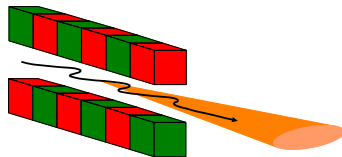


Bending magnet



- Wide horizontal beam
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Wiggler

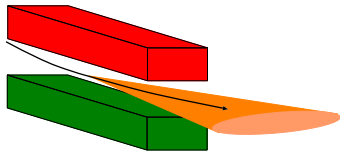


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Synchrotron sources

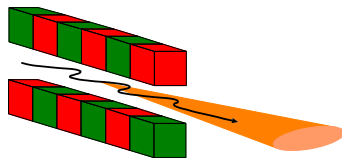


Bending magnet



- Wide horizontal beam
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Wiggler

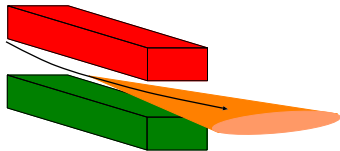


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Synchrotron sources

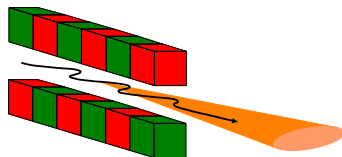


Bending magnet



- Wide horizontal beam
- Broad spectrum
- Low brilliance

Wiggler

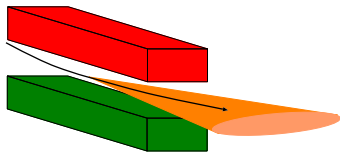


- Wide horizontal beam
- Broad spectrum
- Higher critical energy

Synchrotron sources

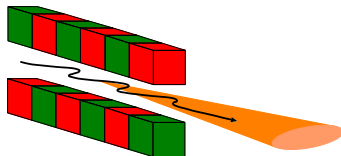


Bending magnet



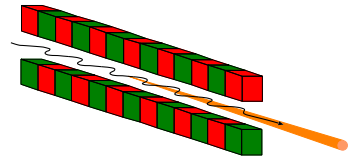
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Wiggler



- Wide horizontal beam
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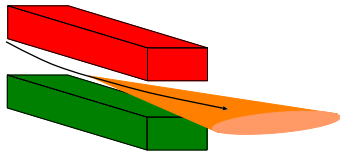
Undulator



Synchrotron sources

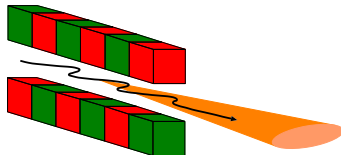


Bending magnet



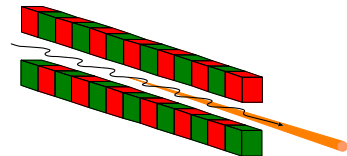
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Wiggler



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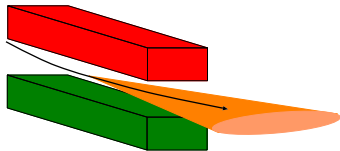


- Highly collimated beam

Synchrotron sources

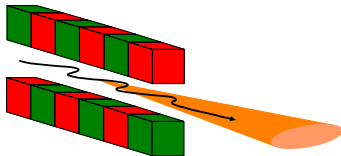


Bending magnet



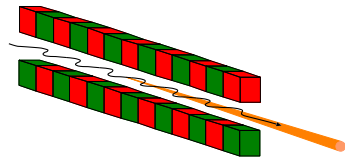
- Wide horizontal beam
- Broad spectrum
- Low brilliance

Wiggler



- Wide horizontal beam
- Broad spectrum
- Higher critical energy

Undulator

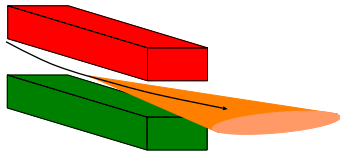


- Highly collimated beam
- Highly peaked spectrum

Synchrotron sources

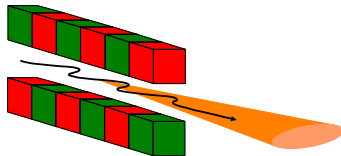


Bending magnet



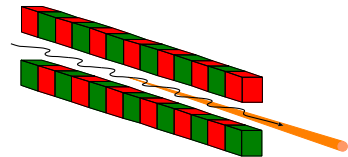
- Wide horizontal beam
- Broad spectrum
- Low brilliance

Wiggler



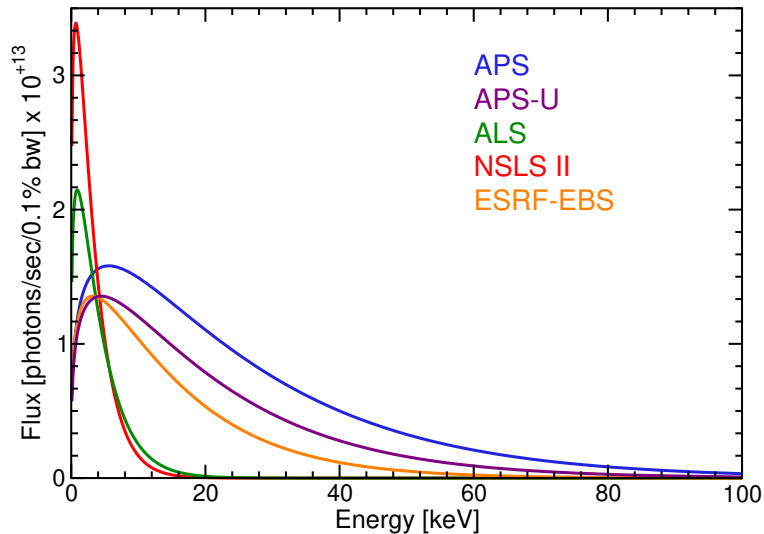
- Wide horizontal beam
- Broad spectrum
- Higher critical energy

Undulator

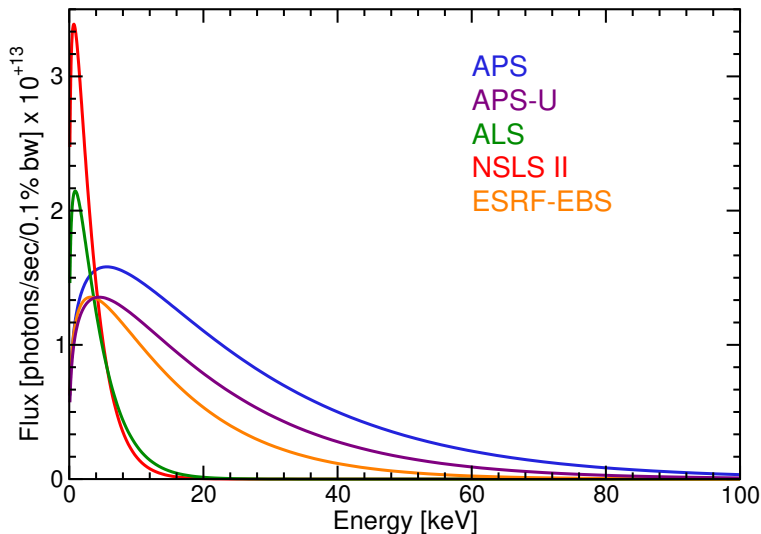


- Highly collimated beam
- Highly peaked spectrum
- High brightness

Bending magnet spectra



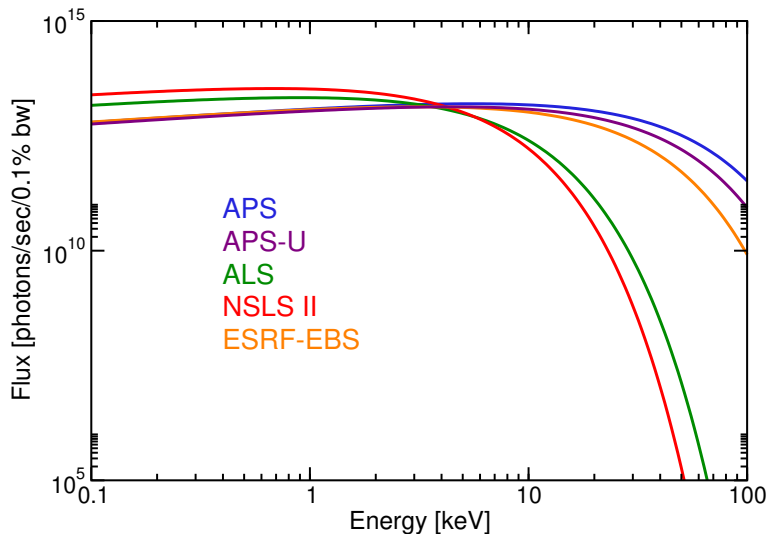
Bending magnet spectra



Lower energy sources, such as NSLS have lower peak energy and higher intensity at the peak.

Logarithmic scale shows clearly how much more energetic and intense the bending magnet sources at the 6 GeV and 7 GeV sources are.

Bending magnet spectra

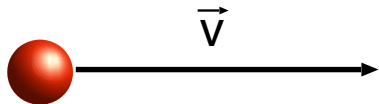


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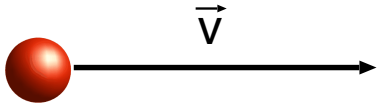
Higher energy sources, such as APS have higher energy spectrum and are only off by a factor of 2 intensity at low energy.

Logarithmic scale shows clearly how much more energetic and intense the bending magnet sources at the 6 GeV and 7 GeV sources are.

Review of special relativity

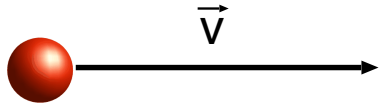


Review of special relativity



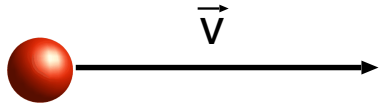
$$\beta = \frac{v}{c}$$

Review of special relativity



$$\beta = \frac{v}{c} \quad \gamma = \sqrt{\frac{1}{1 - \beta^2}}$$

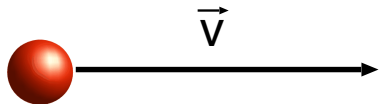
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Review of special relativity

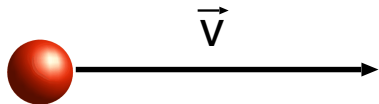


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$$E = \gamma mc^2$$

$$\beta = \sqrt{1 - \frac{1}{\gamma^2}}$$

Review of special relativity



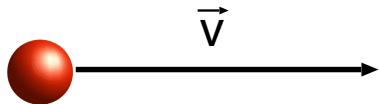
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$$\beta = \sqrt{1 - \frac{1}{\gamma^2}} \rightarrow \beta \approx 1 - \frac{1}{2} \frac{1}{\gamma^2}$$

use binomial expansion since $1/\gamma^2 \ll 1$

Review of special relativity



Let's calculate these quantities
for an electron at NSLS and
APS

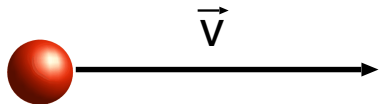
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$$m_e = 0.511 \text{ MeV}/c^2$$

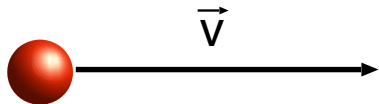
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Review of special relativity



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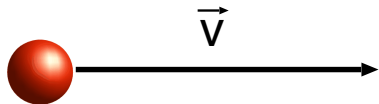
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NSLS II: $E = 3.0 \text{ GeV}$

$$\gamma = \frac{3.0 \times 10^9}{0.511 \times 10^6} = \mathbf{5871}$$

Review of special relativity



$$\beta = \frac{v}{c} \quad \gamma = \sqrt{\frac{1}{1 - \beta^2}}$$

$$E = \gamma mc^2$$

$$\beta = \sqrt{1 - \frac{1}{\gamma^2}} \rightarrow \beta \approx 1 - \frac{1}{2\gamma^2}$$

use binomial expansion since $1/\gamma^2 \ll 1$

Let's calculate these quantities
for an electron at NSLS and
APS

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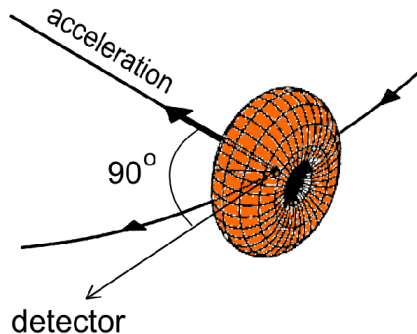
APS: $E = 7.0 \text{ GeV}$

$$\gamma = \frac{7.0 \times 10^9}{0.511 \times 10^6} = 13700$$

“Headlight” effect



In electron rest frame:

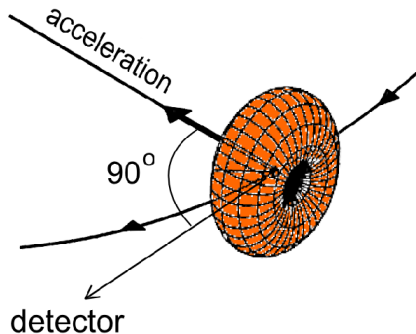


emission is symmetric about the axis of the acceleration vector

"Headlight" effect

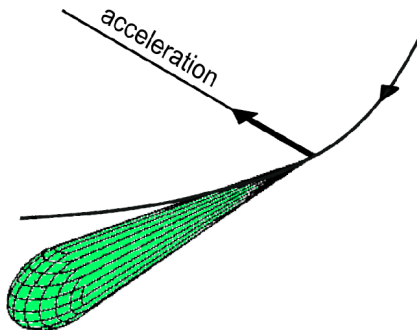


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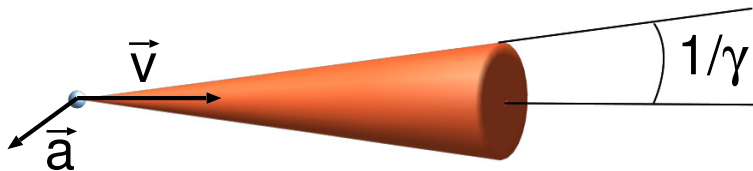
emission is symmetric about the axis of the acceleration vector

In lab frame:



emission is pushed into the direction of motion of the electron

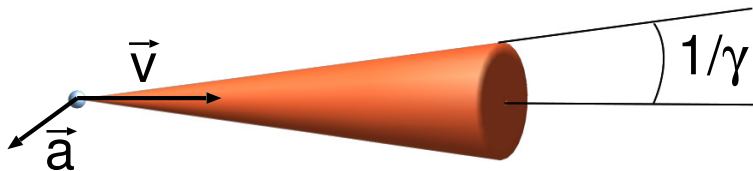
Relativistic emission



the electron is in constant transverse acceleration due to the Lorentz force from the magnetic field of the bending magnet

$$\vec{F} = e\vec{v} \times \vec{B} = m_e \vec{a}$$

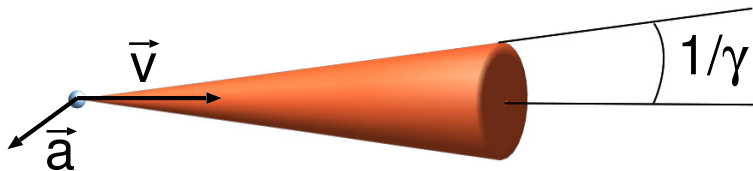
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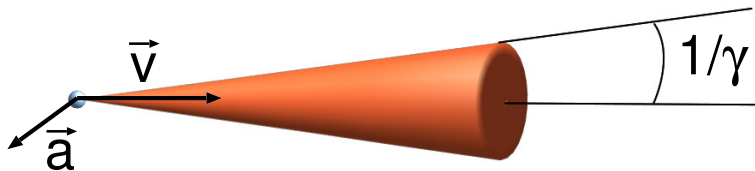


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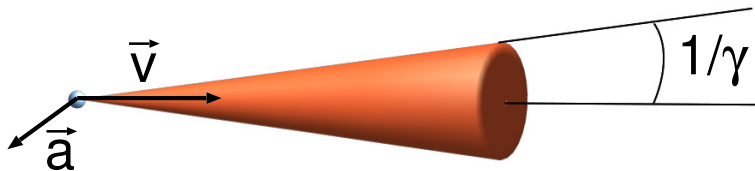
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for the APS, with $\gamma \approx 10^4$ we have

$$E_{max} \approx (10^4)^3 \cdot 10^6 = 10^{18}$$

Flux and brightness



There are a number of important quantities which are relevant to the quality of an x-ray source:

Flux and brightness



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photon flux

source type

optics

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Computing brightness

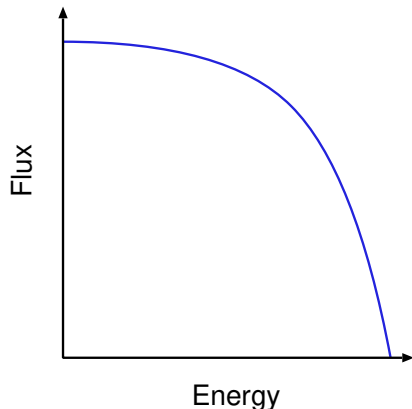


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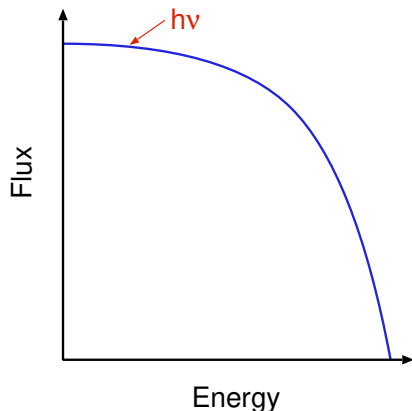


For a specific photon flux distribution, we would normally integrate to get the total flux.

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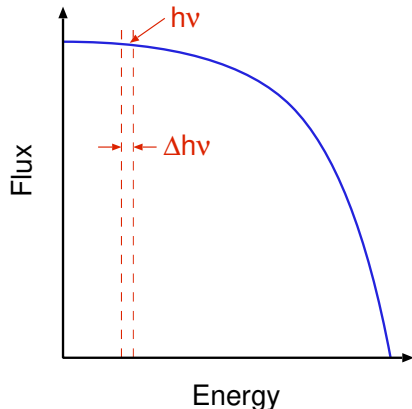


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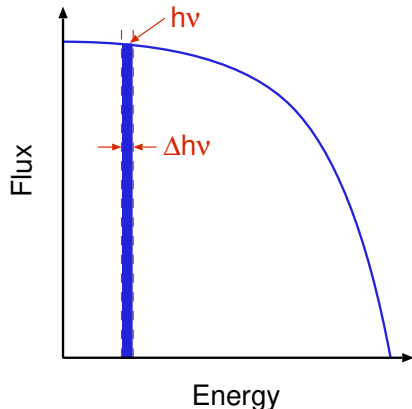
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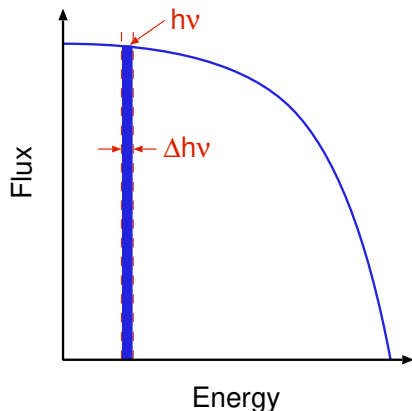
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Compute the integrated photon flux in that bandwidth.

Computing brightness



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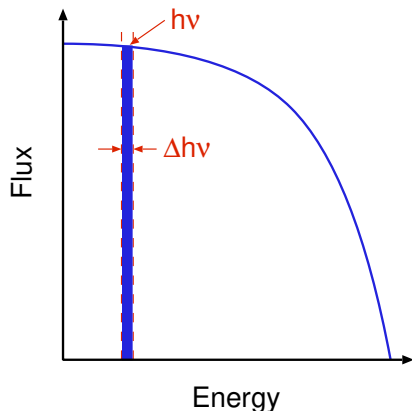


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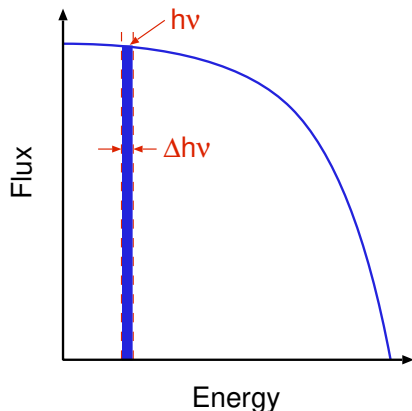


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$$\alpha \approx x/z$$

$$\beta \approx y/z,$$

where z is the distance from the source over which there is a lateral spread x and y in each direction