## Today's Outline - February 18, 2020

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- Liquid scattering


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- Small angle x-ray scattering


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- Liquid scattering
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- Calculating $R_{g}$


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Homework Assignment \#03:
Chapter 3:1,3,4,6,8
due Thursday, February 27, 2020

## Scattering from molecules

From the atomic form factor, we would like to abstract to the next level of complexity, a molecule (we will leave crystals for Chapter 5).

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\begin{aligned}
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$-z=\left(\overline{O O^{\prime}}+\overline{O^{\prime} A}\right) \cdot\left(\overline{O O^{\prime}}+\overline{O^{\prime} B}\right)$


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0 & =3 z^{2}+2 z-1
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u & =\cos ^{-1}(-z)=109.5^{\circ}
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$F_{ \pm}^{\text {mol }}=f^{C}(Q)+f^{F}(Q)\left[3 e^{\mp i Q R / 3}+e^{ \pm i Q R}\right]$
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The plot shows the structure factor of $\mathrm{CF}_{4}$,


## The $\mathrm{CF}_{4}$ scattering factor

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The plot shows the structure factor of $\mathrm{CF}_{4}$, its orientationally averaged structure factor,


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The plot shows the structure factor of $\mathrm{CF}_{4}$, its orientationally averaged structure factor, and the form factor factor of Mo which has the same number of electrons as $\mathrm{CF}_{4}$


## The $\mathrm{CF}_{4}$ scattering factor

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The plot shows the structure factor of $\mathrm{CF}_{4}$, its orientationally averaged structure factor, and the form factor factor of Mo which has the same number of electrons as $\mathrm{CF}_{4}$

The logarithmic plot shows the spherically averaged structure factor compared to the inelastic scattering for $\mathrm{CF}_{4}$


## The radial distribution function

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Ordered 2D crystal
Amorphous solid or liquid

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Ordered 2D crystal


Amorphous solid or liquid


## The radial distribution function

## Ordered 2D crystal



Amorphous solid or liquid


Take a circle (sphere) of radius $r$ and thickness $d r$ and count the number of atom centers lying within the ring.

## The radial distribution function

## Ordered 2D crystal



Amorphous solid or liquid


Take a circle (sphere) of radius $r$ and thickness $d r$ and count the number of atom centers lying within the ring. Then expand the ring radius by $d r$ to map out the radial distribution function $g(r)$

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## Total scattered intensity

Consider a mono-atomic (-molecular) system where the total scattered intensity is given by

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$$

## Total scattered intensity

Consider a mono-atomic (-molecular) system where the total scattered intensity is given by

$$
\begin{aligned}
I(\vec{Q}) & =f(\vec{Q})^{2} \sum_{n} e^{i \vec{Q} \cdot \overrightarrow{r_{n}}} \sum_{m} e^{-i \vec{Q} \cdot \overrightarrow{r_{m}}}=f(\vec{Q})^{2} \sum_{n} \sum_{m} e^{i \vec{Q} \cdot\left(\overrightarrow{r_{n}}-\overrightarrow{r_{m}}\right)} \\
& =N f(\vec{Q})^{2}+f(\vec{Q})^{2} \sum_{n} \sum_{m \neq n} e^{i \vec{Q} \cdot\left(\overrightarrow{r_{n}}-\overrightarrow{r_{m}}\right)}
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\end{aligned}
$$

The sum over $m \neq m$ is now replaced with an integral of the continuous atomic pair density function, $\rho_{n}\left(\vec{r}_{n m}\right)$
$I(\vec{Q})=N f(\vec{Q})^{2}+f(\vec{Q})^{2} \sum_{n} \int_{V} \rho_{n}\left(\vec{r}_{n m}\right) e^{i \vec{Q} \cdot\left(\overrightarrow{r_{n}}-\overrightarrow{r_{m}}\right)} d V_{m}$

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The sum over $m \neq m$ is now replaced with an integral of the continuous atomic pair density function, $\rho_{n}\left(\vec{r}_{n m}\right)$ and adding and subtracting the average atomic density $\rho_{\text {at }}$

$$
\begin{aligned}
& I(\vec{Q})= N f(\vec{Q})^{2}+f(\vec{Q})^{2} \sum_{n} \int_{V} \rho_{n}\left(\overrightarrow{r_{n m}}\right) e^{i \vec{Q} \cdot\left(\overrightarrow{r_{n}}-\overrightarrow{r_{m}}\right)} d V_{m} \\
&=N f(\vec{Q})^{2}+f(\vec{Q})^{2} \sum_{n} \int_{V}\left[\rho_{n}\left(\overrightarrow{r_{n m}}\right)-\rho_{a t}\right] e^{i \vec{Q} \cdot\left(\overrightarrow{r_{n}}-\overrightarrow{r_{m}}\right)} d V_{m} \\
&+f(\vec{Q})^{2} \rho_{a t} \sum_{n} \int_{V} e^{i \vec{Q} \cdot\left(\overrightarrow{\left.r_{n}-r_{m}\right)} d V_{m}\right.}
\end{aligned}
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& I(\vec{Q})= N f(\vec{Q})^{2}+f(\vec{Q})^{2} \sum_{n} \int_{V} \rho_{n}\left(\overrightarrow{r_{n m}}\right) e^{i \vec{Q} \cdot\left(\overrightarrow{r_{n}}-\overrightarrow{r_{m}}\right)} d V_{m} \\
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&+f(\vec{Q})^{2} \rho_{a t} \sum_{n} \int_{V} e^{i \vec{Q} \cdot\left(\overrightarrow{r_{n}}-\overrightarrow{r_{m}}\right)} d V_{m}=I^{S R O}(\vec{Q})
\end{aligned}
$$

## Total scattered intensity

Consider a mono-atomic (-molecular) system where the total scattered intensity is given by

$$
\begin{aligned}
I(\vec{Q}) & =f(\vec{Q})^{2} \sum_{n} e^{i \vec{Q} \cdot \overrightarrow{r_{n}}} \sum_{m} e^{-i \vec{Q} \cdot \overrightarrow{r_{m}}}=f(\vec{Q})^{2} \sum_{n} \sum_{m} e^{i \vec{Q} \cdot\left(\overrightarrow{r_{n}}-\overrightarrow{r_{m}}\right)} \\
& =N f(\vec{Q})^{2}+f(\vec{Q})^{2} \sum_{n} \sum_{m \neq n} e^{i \vec{Q} \cdot\left(\overrightarrow{r_{n}}-\overrightarrow{r_{m}}\right)}
\end{aligned}
$$

The sum over $m \neq m$ is now replaced with an integral of the continuous atomic pair density function, $\rho_{n}\left(\vec{r}_{n m}\right)$ and adding and subtracting the average atomic density $\rho_{\text {at }}$

$$
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& I(\vec{Q})= N f(\vec{Q})^{2}+f(\vec{Q})^{2} \sum_{n} \int_{V} \rho_{n}\left(\vec{r}_{n m}\right) e^{i \vec{Q} \cdot\left(\overrightarrow{r_{n}}-\overrightarrow{r_{m}}\right)} d V_{m} \\
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Which is the sine Fourier Transform of the deviation of the atomic density from its average, $\mathcal{H}(r)=4 \pi r[g(r)-1]$

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The relation between radial distribution function and structure factor can be extended to multi-component systems where $g(r) \rightarrow g_{i j}(r)$ and $S(Q) \rightarrow S_{i j}(Q)$.

## Structure in supercooled liquid metals

Measurement of the liquid structure factor of molten metals have shown that there is short range order which leads to the phenomenon of supercooling.

"Difference in Icosahedral Short-Range Order in Early and Late Transition Metal Liquids",
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## Water dynamics

## Liquid scattering can also be used to study dynamics

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In this article, the authors measured the liquid scattering as a function of both momentum, $Q$, and energy, $E$, transfer by using analyzers set for a specific energy ( 21.747 keV ) but varying $Q$ and then scanning the incident energy at fixed incident angle

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The Van Hoff function can be obtained by a double Fourier transform

$$
g(r, t)-1=\frac{1}{2 \rho \pi^{2} r} \iint e^{i \omega t} \sin (Q r) Q S(Q, E) d Q d E
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[^1]
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The first and second peaks are highly coupled in space and time and merge within 0.8 ps . This behavior is different from liquid metals and leads to the viscosity of water.
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Where we have assumed sufficient averaging and introduced $\rho_{s l}=f \rho_{a t}$. This final expression looks just like an atomic form factor but the charge density that we consider here is on a much longer length scale than an atom.

## The SAXS experiment



## Scattering from a dilute solution

The simplest case is for a dilute solution of non-interacting molecules.

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I^{S A X S}(\vec{Q})=\left|\int_{V_{p}} \rho_{s l} e^{i \vec{Q} \cdot \vec{r}} d V_{p}\right|^{2}
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Where $\Delta \rho=\left(\rho_{s l, p}-\rho_{s l, 0}\right)$, and the form factor depends on the morphology of the particle (size and shape).

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$\mathrm{p}=10 \%$
$p=20 \%$



[^0]:    "Seeing real-space dynamics of liquid water through inelastic x-ray scattering," T. Iwashita et al. Sci. Adv. 3, e1603079 (2017).

[^1]:    "Seeing real-space dynamics of liquid water through inelastic x-ray scattering," T. Iwashita et al. Sci. Adv. 3, e1603079 (2017).

