

Today's outline - February 06, 2020

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- Designing a multilayer

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- Reflection from a graded index

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Homework Assignment #02:
Problems on Blackboard
due Tuesday, February 18, 2020

Multilayer design

Materials for multilayer monochromator chosen to reflect 12 keV x-rays at ~ 2 degrees with 0.5% and 1.0% bandwidth

A. Khounsary et al., "A dual-bandwidth multilayer monochromator system," *Proc. SPIE* **10760**, 107600j (2018).

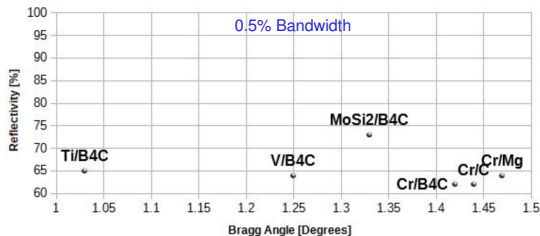
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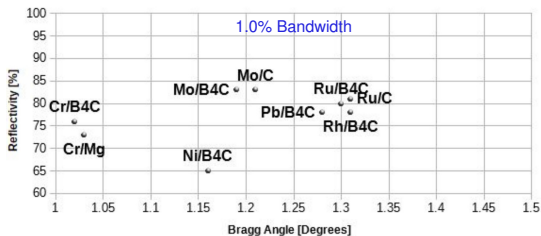
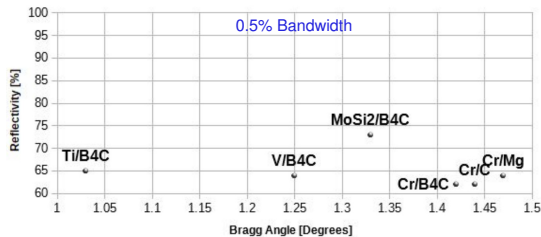


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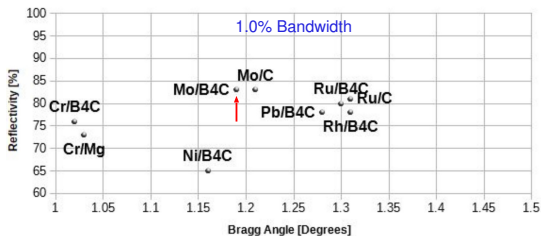
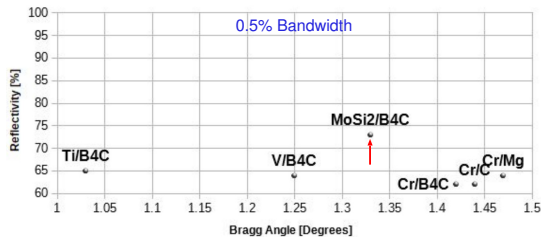


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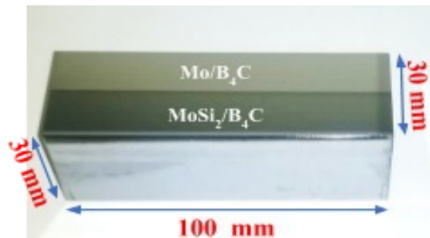
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MoSi₂/B₄C and Mo/B₄C were selected for the 0.5% and 1.0% bandwidth coatings, respectively

Multilayer fabrication & testing

The 0.5% and 1.0% bandwidth layers were deposited side-by-side on a monolithic 20 mm × 30 mm × 100 mm polished silicon block



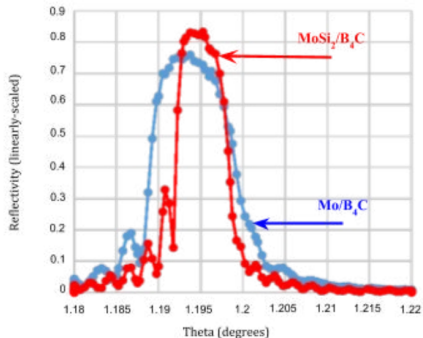
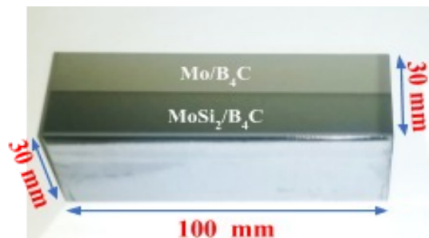
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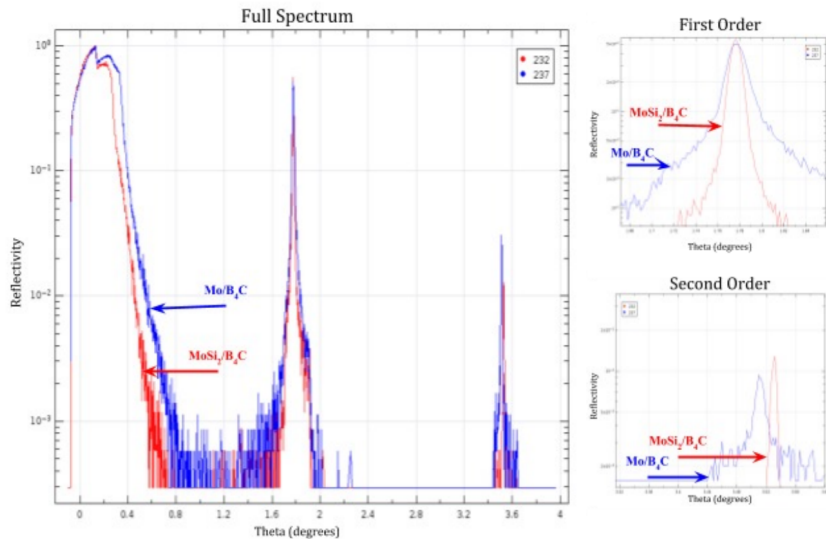
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When illuminated with 12 keV x-rays the two multilayers showed diffraction peaks at nearly the same angle. The reflectivities were both over 75% and the bandwidths were 0.52% and 0.86%, respectively.

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Multilayer spectrum



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Graded interfaces

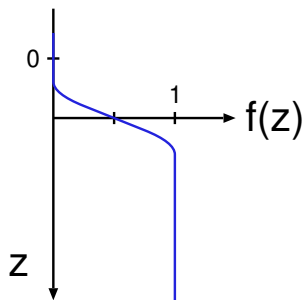
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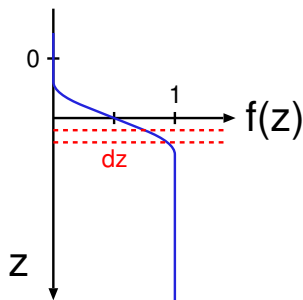


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The density profile of the interface can be described by the function $f(z)$ which approaches 1 as $z \rightarrow \infty$.

The reflectivity can be computed as the superposition of the reflectivity of a series of infinitesimal slabs of thickness dz at a depth z .

Reflectivity of a graded interface

The differential reflectivity from a slab of thickness dz at depth z is:

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Calculating the full reflection coefficient relative to the Fresnel reflection coefficient

$$\frac{R(Q)}{R_F(Q)} = \left| \int_{-\infty}^{\infty} \left(\frac{df}{dz} \right) e^{iQz} dz \right|^2$$

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The error function - a specific case

The error function is often chosen as a model for the density gradient

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whose Fourier transform is also a Gaussian, which when squared to obtain the reflection coefficient, gives.

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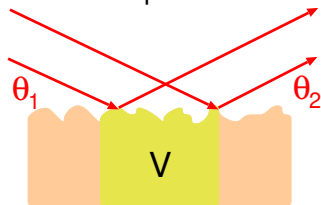
$$Q = k \sin \theta, \quad Q' = k' \sin \theta'$$

Rough surfaces

When a surface or interface is not perfectly smooth but has some roughness the reflectivity is no longer simply specular but has a non-zero diffuse component which we must include in the model.

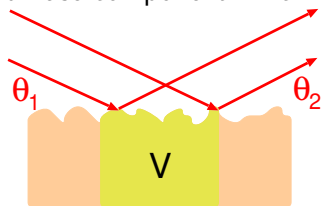
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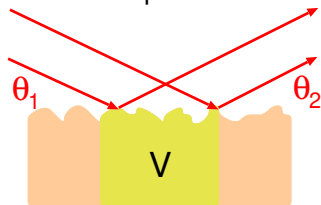
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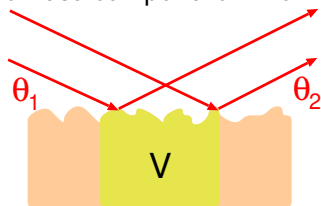
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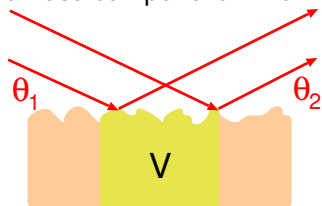


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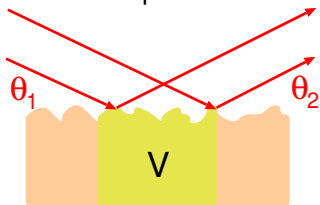


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The actual scattering cross section is the square of this integral

$$\frac{d\sigma}{d\Omega} = \left(\frac{r_0 \rho}{Q_z} \right)^2 \int_S \int_{S'} e^{i Q_z (h(x, y) - h(x', y'))} e^{i Q_x (x - x')} e^{i Q_y (y - y')} dx dy dx' dy'$$

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Finally, it is assumed that the statistics of the height variation are Gaussian and

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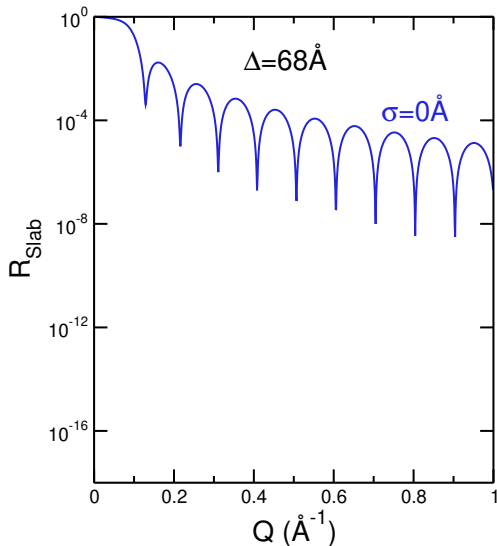
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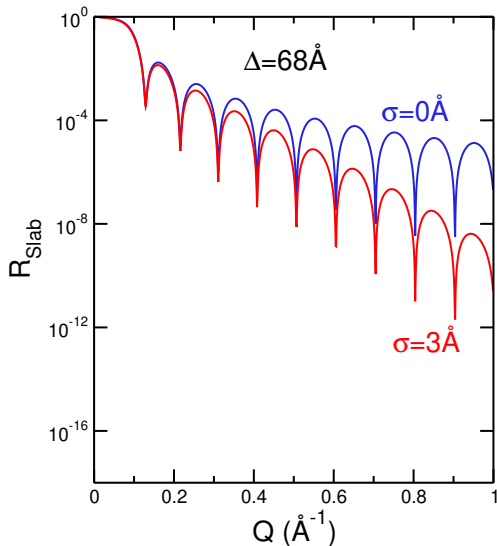


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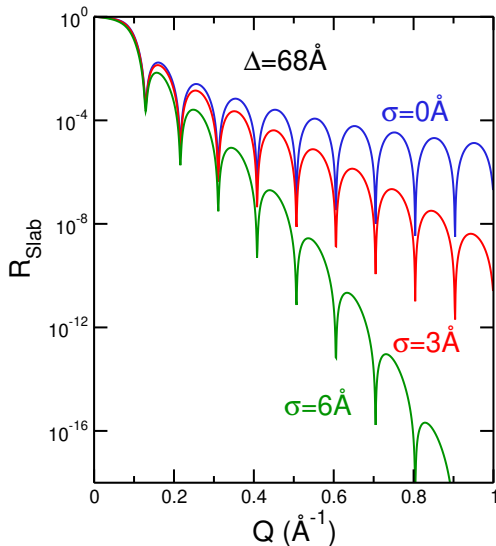
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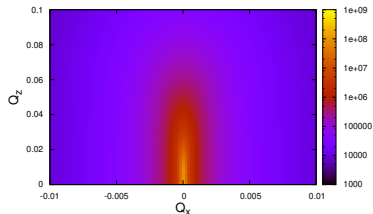
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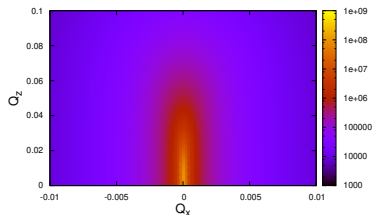
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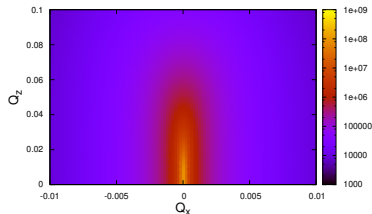
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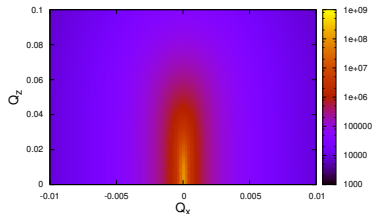
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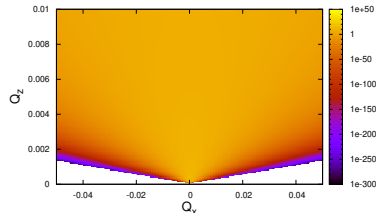
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$$C(x, y) = \sigma^2 e^{-(r/\xi)^{2h}}$$

$$\begin{aligned} \left(\frac{d\sigma}{d\Omega} \right) &= \left(\frac{r_0 \rho}{Q_z} \right)^2 \frac{A_0}{\sin \theta_1} e^{-Q_z^2 \sigma^2} \int e^{Q_z^2 C(x, y)} e^{iQ_x x} e^{iQ_y y} dx dy \\ &= \left(\frac{r_0 \rho}{Q_z} \right)^2 \frac{A_0}{\sin \theta_1} e^{-Q_z^2 \sigma^2} \int \left[e^{Q_z^2 C(x, y)} - 1 + 1 \right] e^{iQ_x x} e^{iQ_y y} dx dy \end{aligned}$$

Bounded correlations

If the correlations remain bounded as $r \rightarrow \infty$

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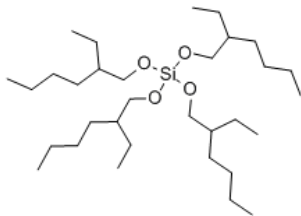
Layering in liquid films

TEHOS, tetrakis-(2-ethylhexoxy)-silane, a non-polar, roughly spherical molecule, was deposited on Si(111) single crystals

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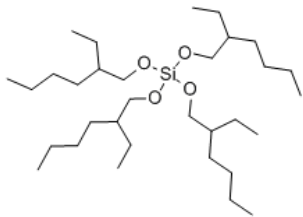
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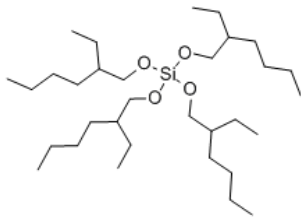


Specular reflection measurements were made at MRCAT (Sector 10 at APS) and at X18A (at NSLS).

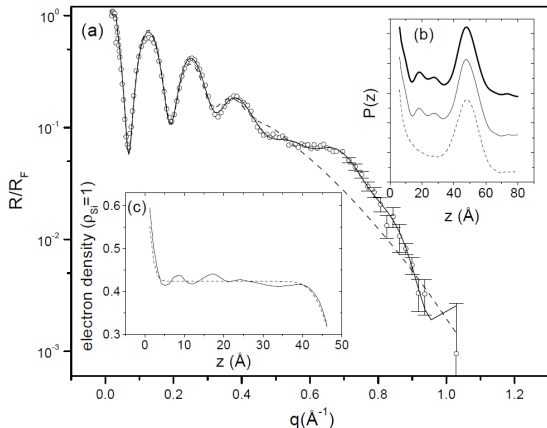
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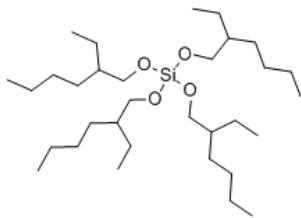
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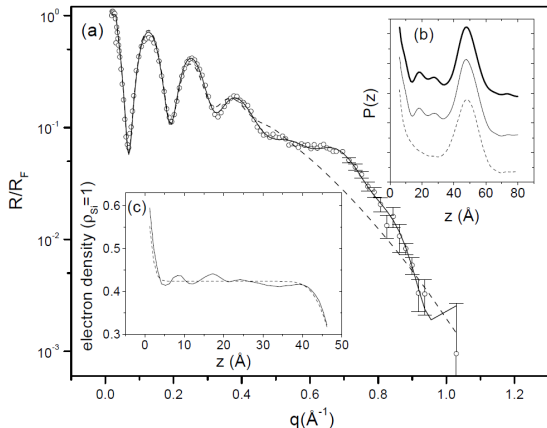
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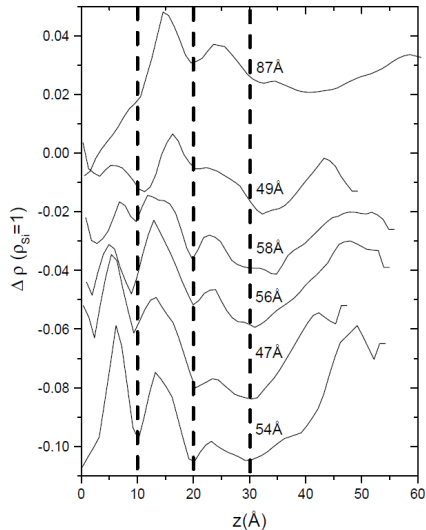
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Deviations from uniform density are used to fit experimental reflectivity

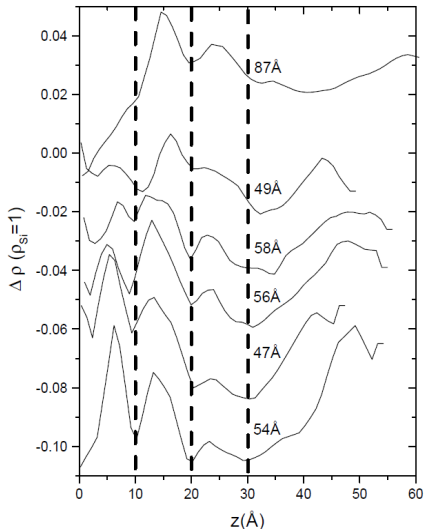
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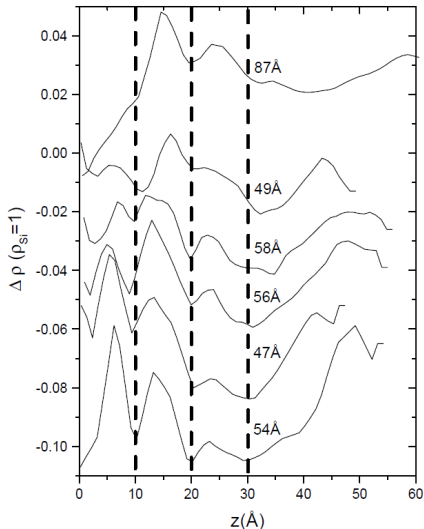
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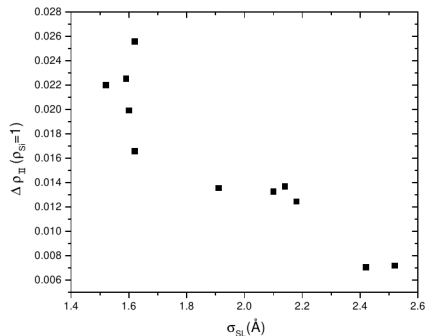


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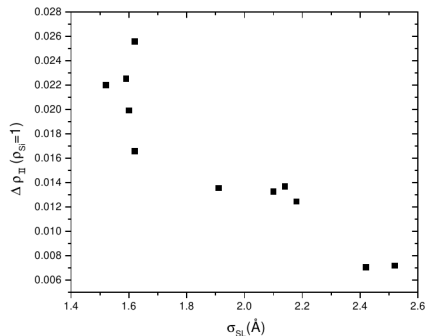
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The authors conclude that the presence of a hard smooth surface is required for ordering and therefore deviations from an ideal, isotropic liquid.

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Film growth kinetics

The goal of this project was to understand the evolution of surface roughness during the growth of a silver thin film.

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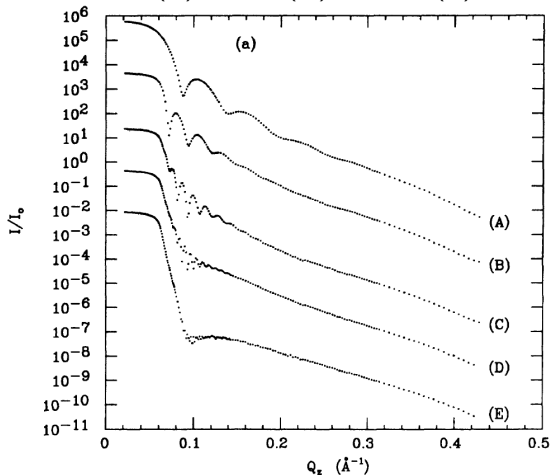
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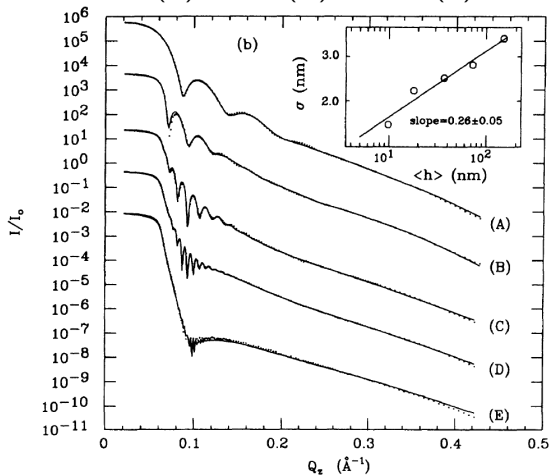
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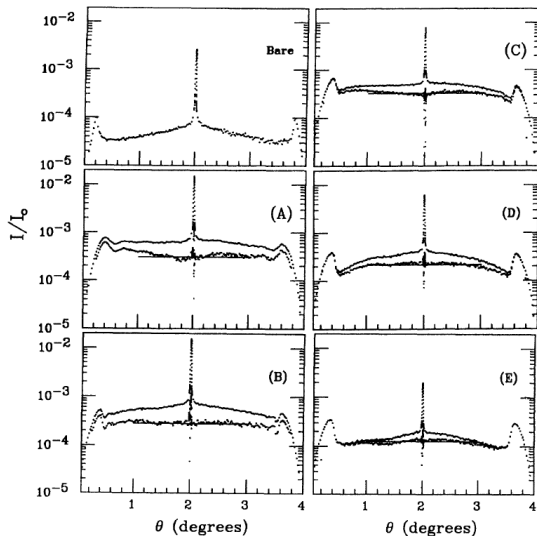
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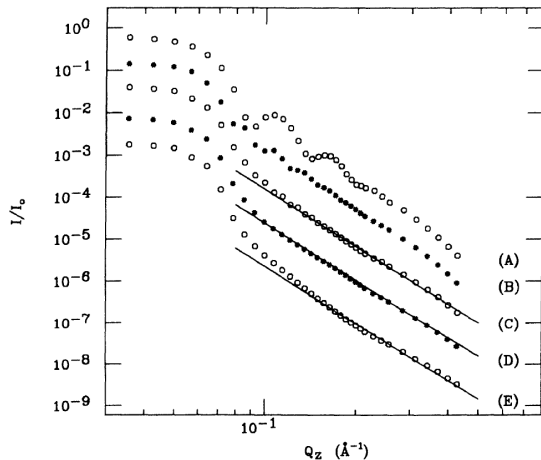
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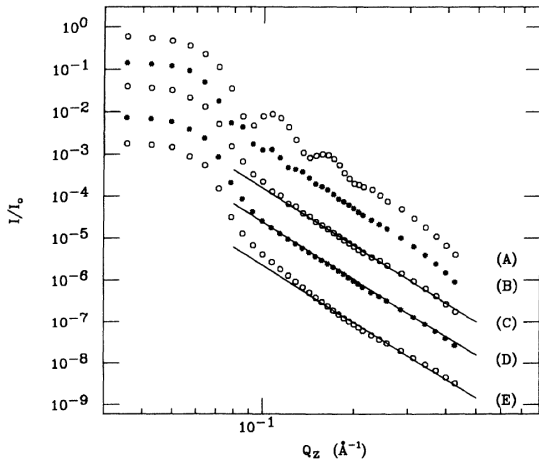


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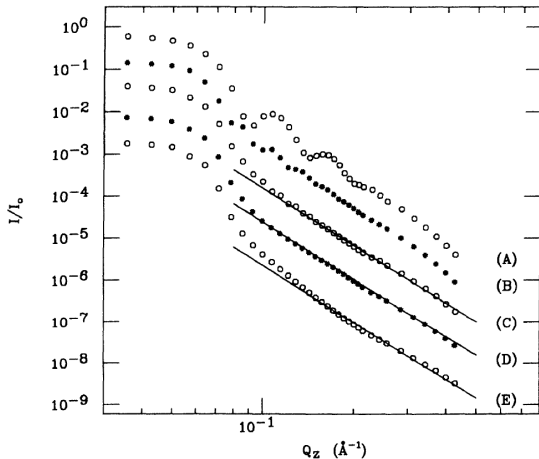
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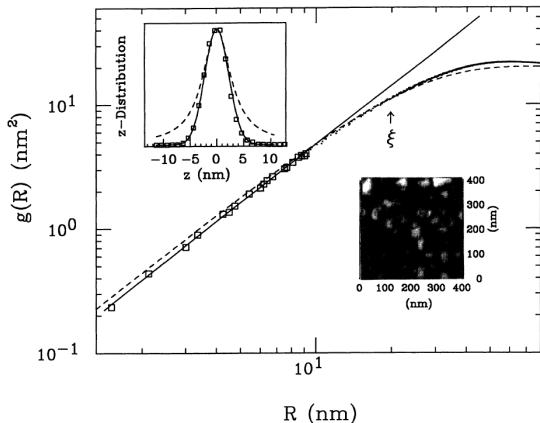
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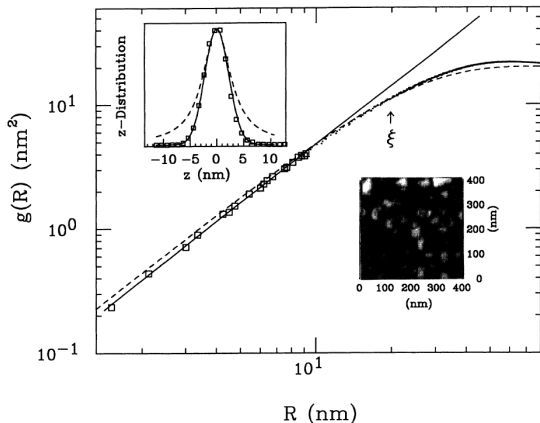
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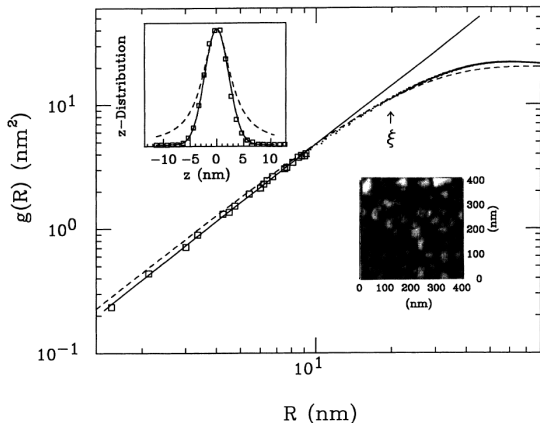
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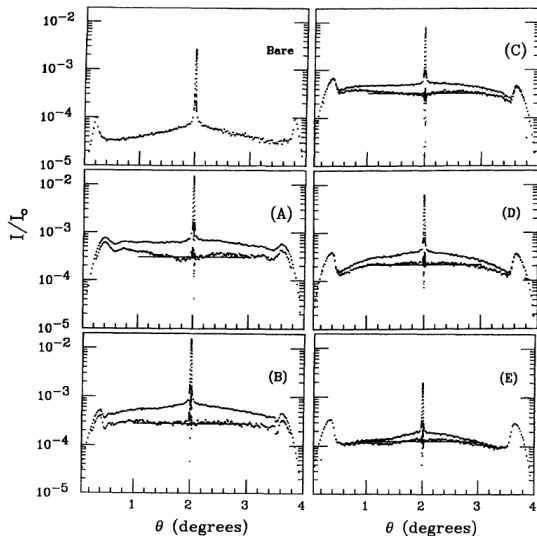
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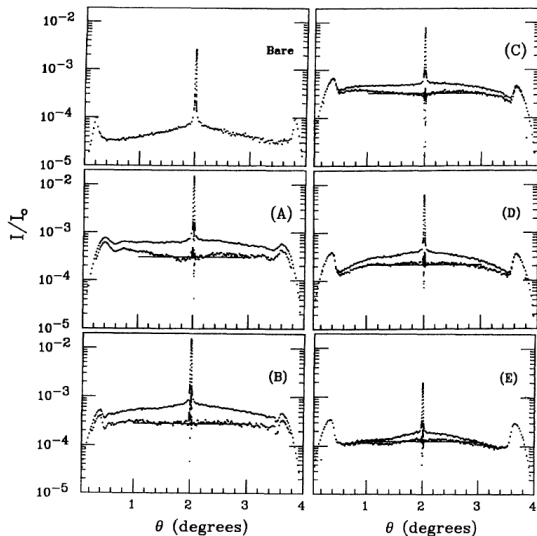
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X-ray reflectivity using synchrotron radiation has made possible the study of the surface of liquid metals

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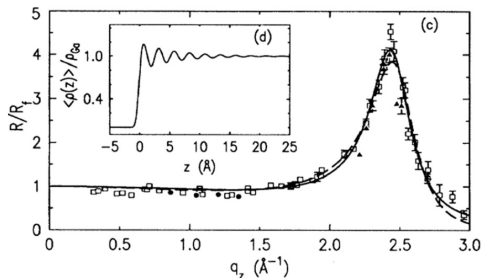
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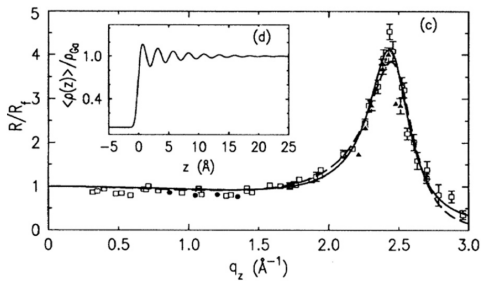
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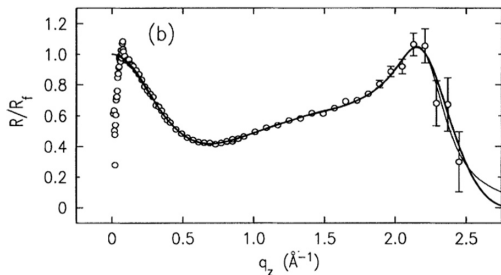
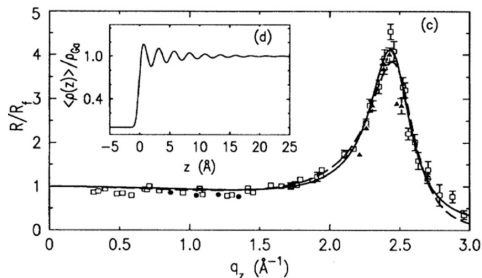
Liquid metal surfaces

X-ray reflectivity using synchrotron radiation has made possible the study of the surface of liquid metals

a liquid can be described as charged ions in a sea of conduction electrons

this leads to a well-defined surface structure as can be seen in liquid gallium

contrast this with the scattering from liquid mercury



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Liquid metal eutectics

High vapor pressure and thermal excitations limit the number of pure metals which can be studied but alloy eutectics provide many possibilities

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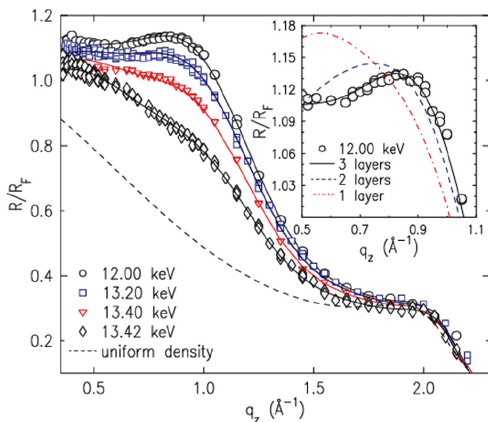
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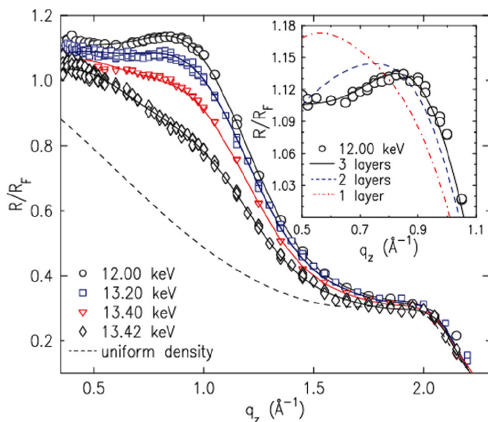
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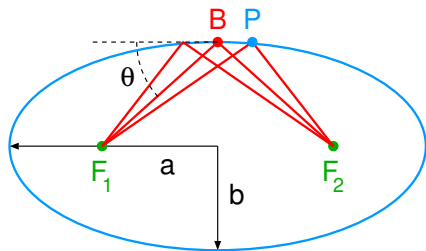
surface layer is rich in Bi (95%), second layer is deficient (25%), and third layer is rich in Bi (53%) once again



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Tangential focusing mirror

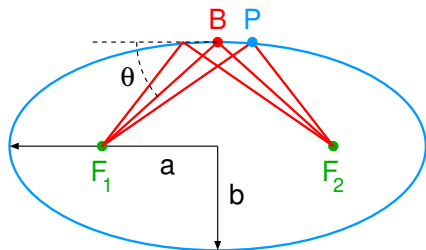
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Tangential focusing mirror

The shape of an ideal mirror is an ellipse, where any ray coming from one focus will be projected to the second focus. Consider a 1:1 focusing mirror. For an ellipse the sum of the distances from any point on the ellipse to the foci is a constant.

$$F_1P + F_2P = 2a$$

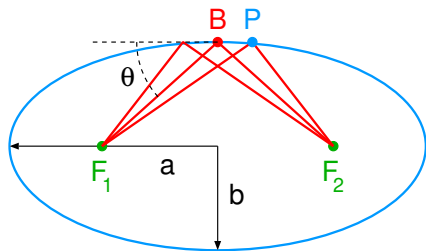


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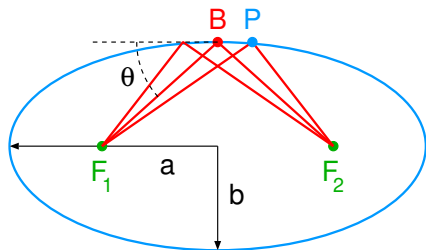
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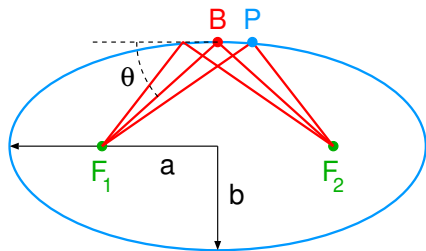
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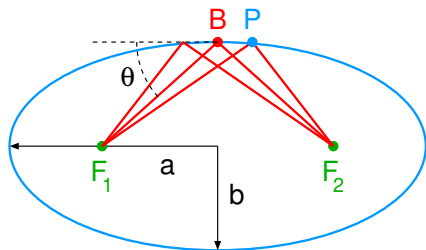
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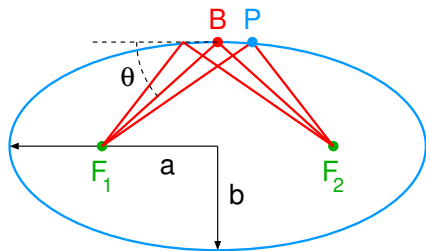
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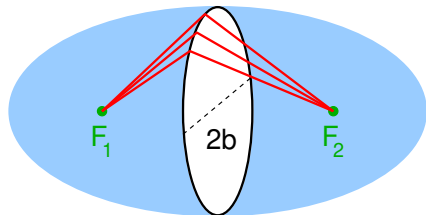


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Saggital focusing mirror

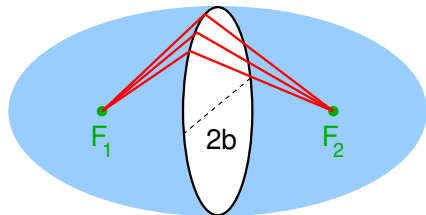
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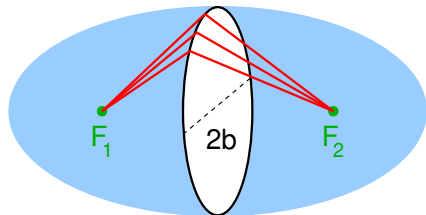


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The tangential focus is also usually approximated by a circular cross-section with radius



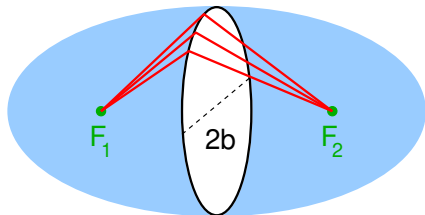
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$$\rho_{sagittal} = b = 2f \sin \theta$$

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$$\rho_{tangential} = a = \frac{2f}{\sin \theta}$$



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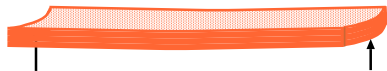


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A cost effective way to focus in both directions is a toroidal mirror which has a fixed bend in the transverse direction but which can be bent longitudinally to change the vertical focus.



Dual focusing options

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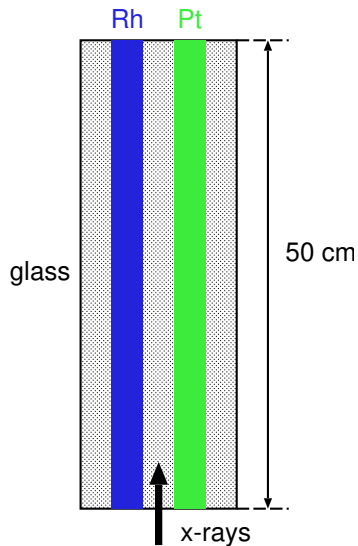
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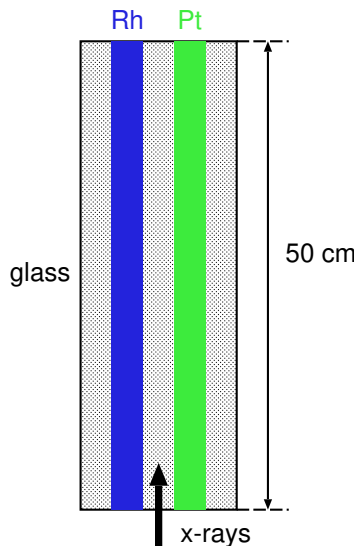
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- **Refractive lenses** — good final focus, focus moves with energy, significant attenuation and hard to change focal length

The MRCAT mirror

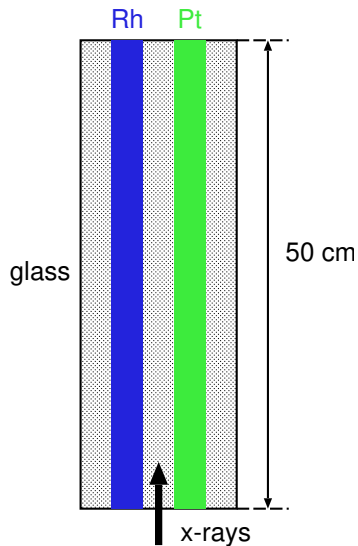


The MRCAT mirror



Ultra low expansion glass polished to a few Å roughness

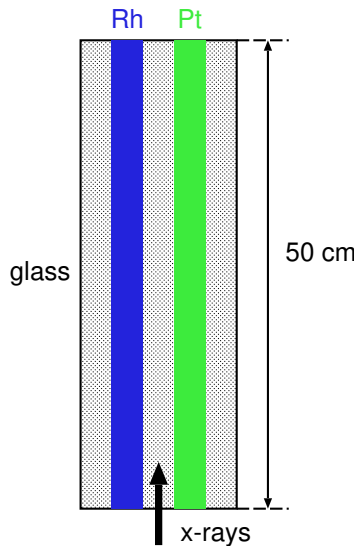
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One platinum stripe and one rhodium stripe deposited along the length of the mirror on top of a chromium buffer layer

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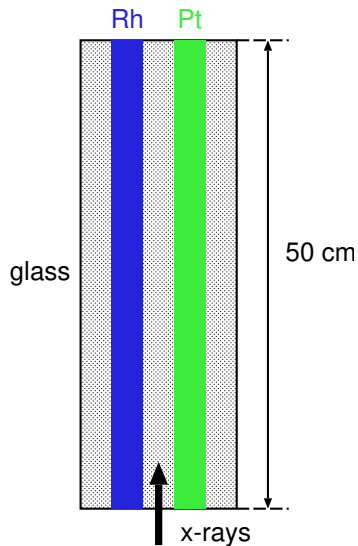


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A mounting system which permits angular positioning to less than 1/100 of a degree as well as horizontal and vertical motions

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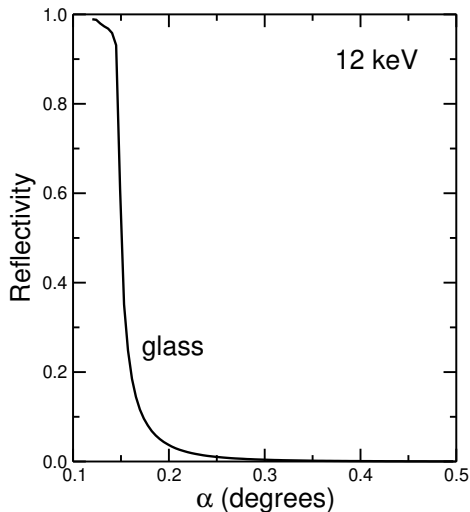
One platinum stripe and one rhodium stripe deposited along the length of the mirror on top of a chromium buffer layer

A mounting system which permits angular positioning to less than 1/100 of a degree as well as horizontal and vertical motions

A bending mechanism to permit vertical focusing of the beam to $\sim 60 \mu\text{m}$

Mirror performance

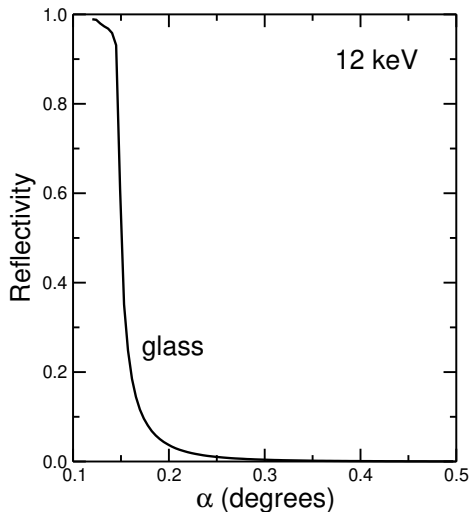
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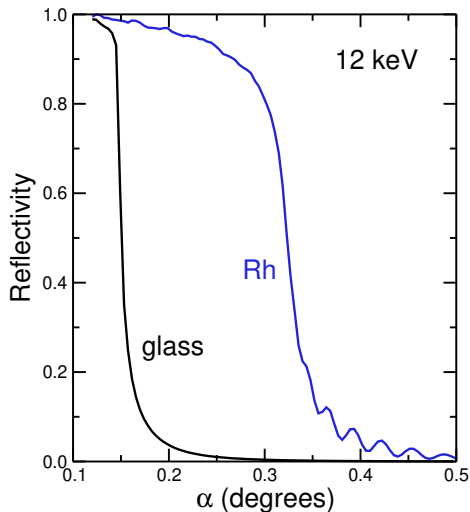
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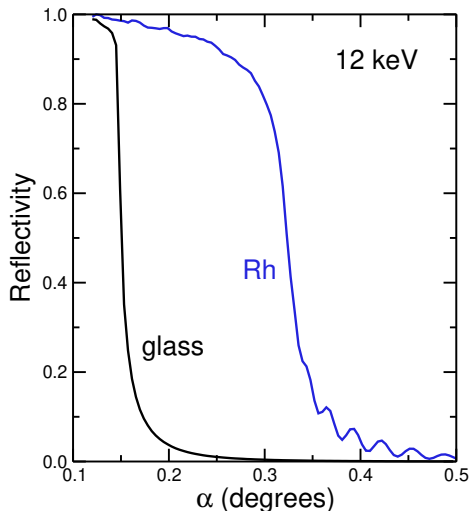


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The Pt stripe gives a higher critical angle still but a lower reflectivity and it looks like an infinite slab.

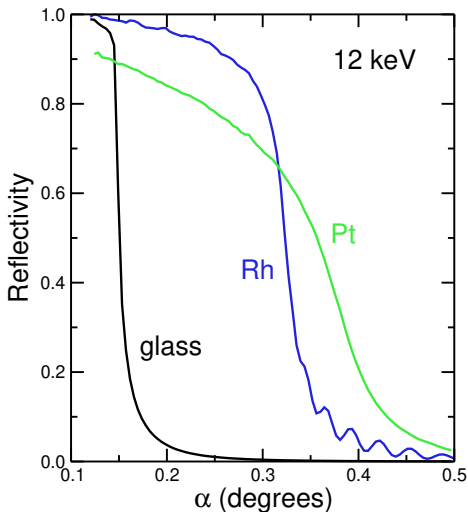


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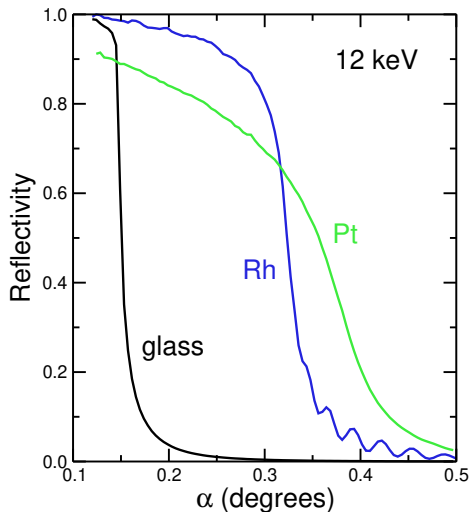


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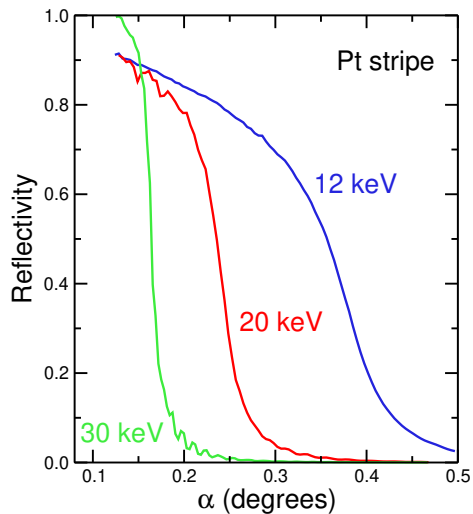
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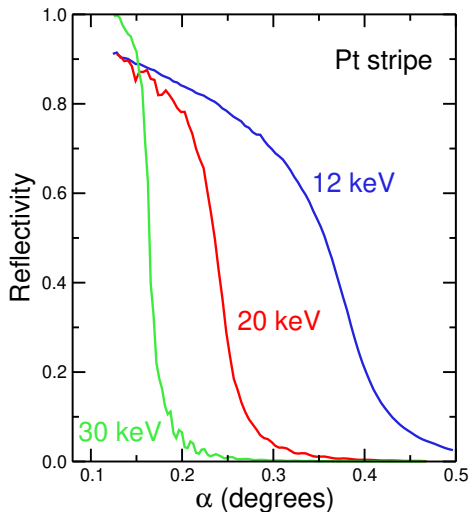
The Pt stripe gives a higher critical angle still but a lower reflectivity and it looks like an infinite slab. Why?



Mirror performance (cont.)

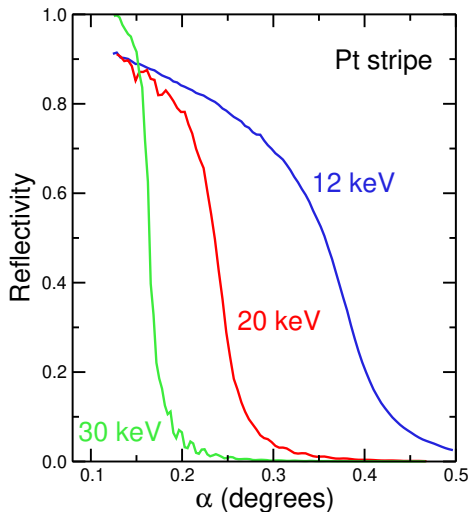


Mirror performance (cont.)



As we move up in energy the critical angle for the Pt stripe drops.

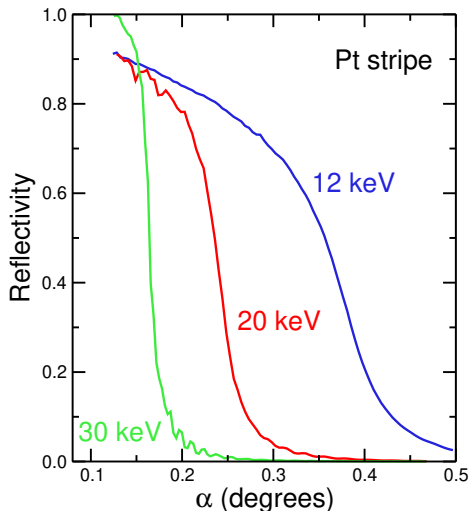
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As energy rises, the Pt layer begins to show the reflectivity of a thin slab.