

Today's outline - January 21, 2020

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Homework Assignment #01:

Chapter 2: 2,3,5,6,8

due Thursday, January 30, 2020

Refraction of x-rays

X-rays can be treated like light when interaction with a medium. However, unlike visible light, the index of refraction of x-rays in matter is very close to unity:

Refraction of x-rays

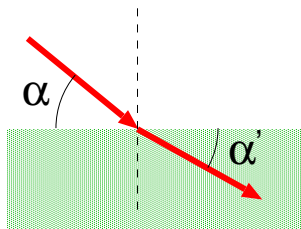
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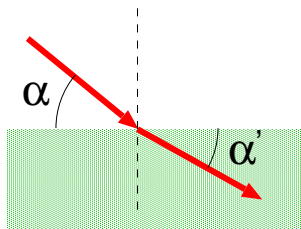


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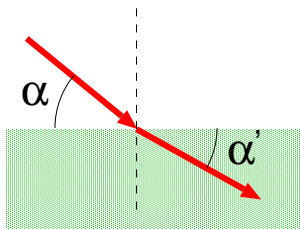
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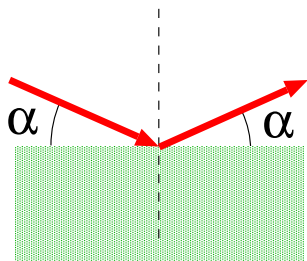
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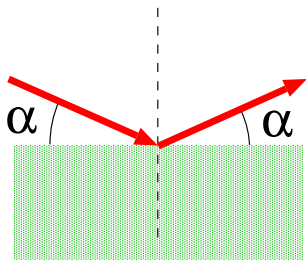
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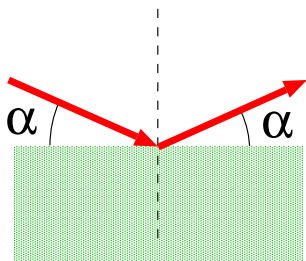


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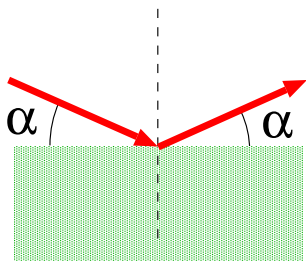
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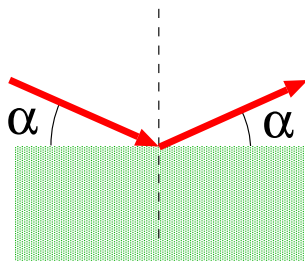
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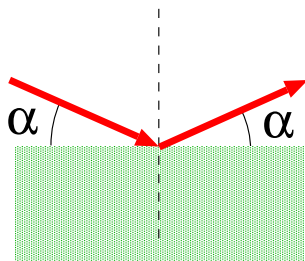
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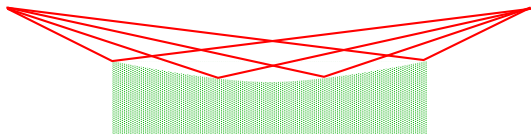
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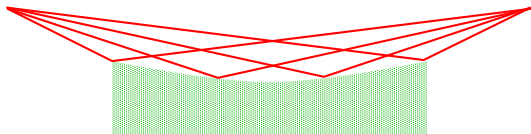
$$\delta = \frac{\alpha_c^2}{2} \quad \longrightarrow \quad \alpha_c = \sqrt{2\delta}$$

Uses of total external reflection



X-ray mirrors

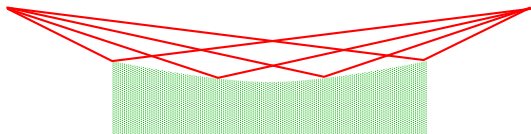
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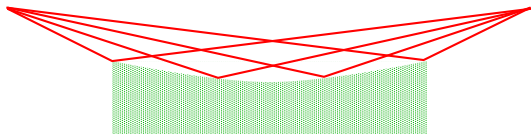
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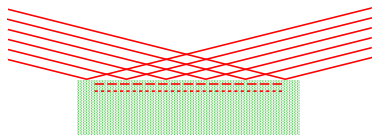
- harmonic rejection
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Uses of total external reflection



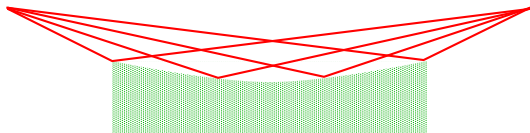
X-ray mirrors

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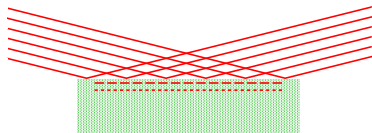
Evanscent wave experiments

Uses of total external reflection



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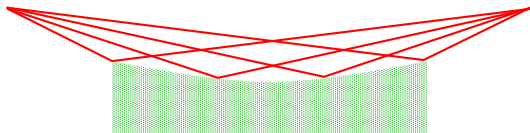
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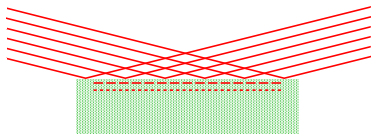
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Uses of total external reflection



X-ray mirrors

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Coherence: what is it?

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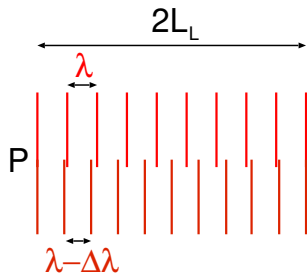
Because of these imperfections the “coherence length” of an x-ray beam is finite and we can calculate it.

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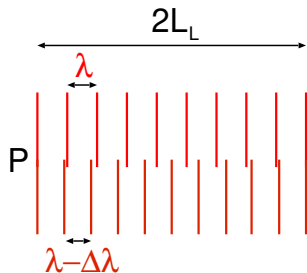
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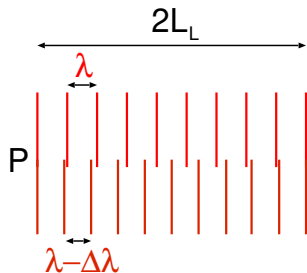


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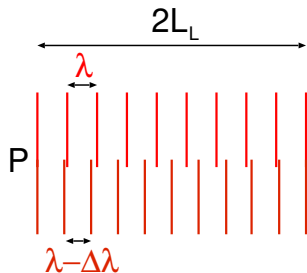
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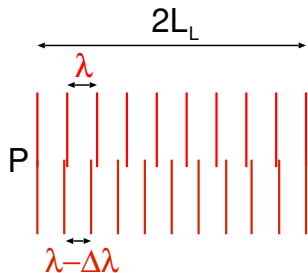
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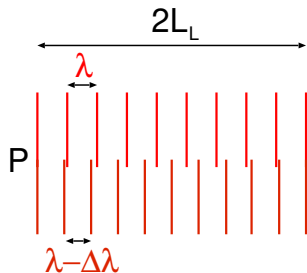
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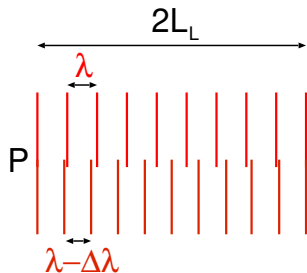
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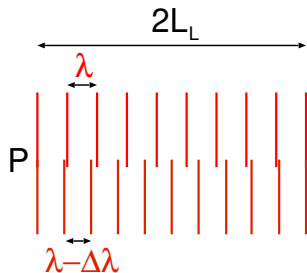
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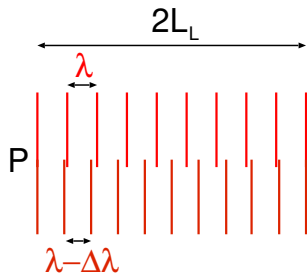
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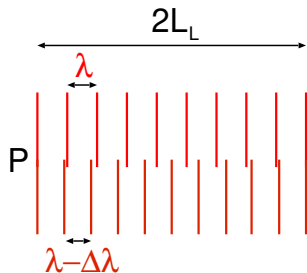
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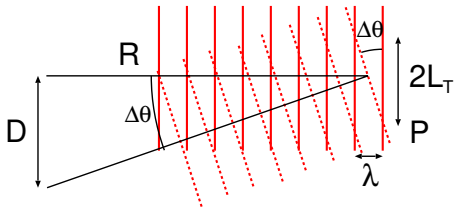
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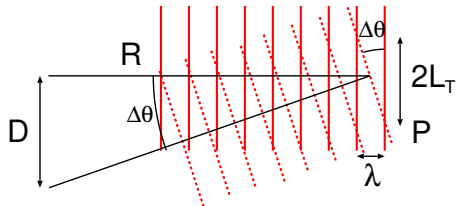
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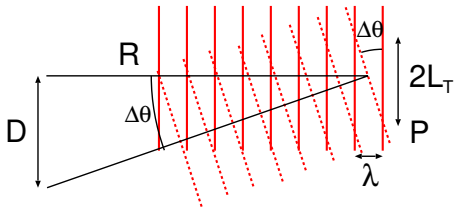
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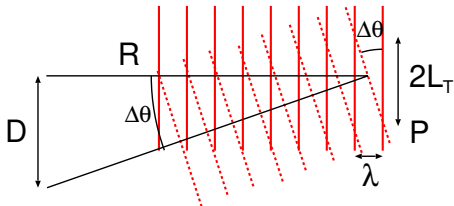
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$$L_T = \frac{\lambda R}{2D}$$

Coherence lengths at the APS

For a typical 3rd generation undulator source, such as at the Advanced Photon Source the vertical source size is $D_v = 10\mu\text{m}$ and we are typically $R = 50\text{m}$ away with our experiment. If we assume a typical wavelength of $\lambda = 1\text{\AA}$, and a monochromator resolution of $\Delta\lambda/\lambda = 10^{-5}$ we have for the vertical direction:

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$$L_L = \frac{\lambda^2}{2\Delta\lambda} = \frac{\lambda}{2} \cdot \frac{\lambda}{\Delta\lambda} = \frac{1 \times 10^{-10}}{2 \cdot 10^{-5}}$$

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For a typical 3rd generation undulator source, such as at the Advanced Photon Source the vertical source size is $D_v = 10\mu\text{m}$ and we are typically $R = 50\text{m}$ away with our experiment. If we assume a typical wavelength of $\lambda = 1\text{\AA}$, and a monochromator resolution of $\Delta\lambda/\lambda = 10^{-5}$ we have for the vertical direction:

$$L_L = \frac{\lambda^2}{2\Delta\lambda} = \frac{\lambda}{2} \cdot \frac{\lambda}{\Delta\lambda} = \frac{1 \times 10^{-10}}{2 \cdot 10^{-5}} = 5\mu\text{m}$$

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$$L_T = \frac{\lambda R}{2D}$$

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$$L_T = \frac{\lambda R}{2D} = \frac{(1 \times 10^{-10}) \cdot 50}{2 \cdot (10 \times 10^{-6})}$$

Coherence lengths at the APS

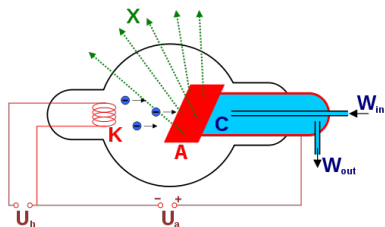
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$$L_T = \frac{\lambda R}{2D} = \frac{(1 \times 10^{-10}) \cdot 50}{2 \cdot (10 \times 10^{-6})} = 250\mu\text{m}$$

X-ray tube schematics

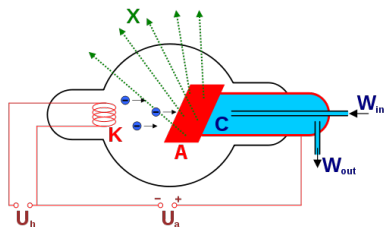
Fixed anode tube



- low power
- low maintenance

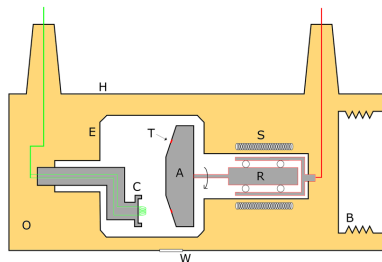
X-ray tube schematics

Fixed anode tube



- low power
- low maintenance

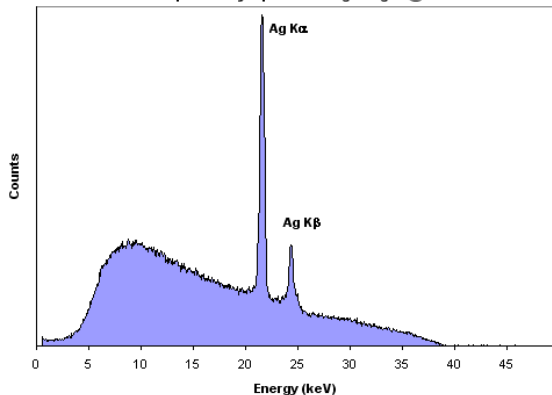
Rotating anode tube



- high power
- high maintenance

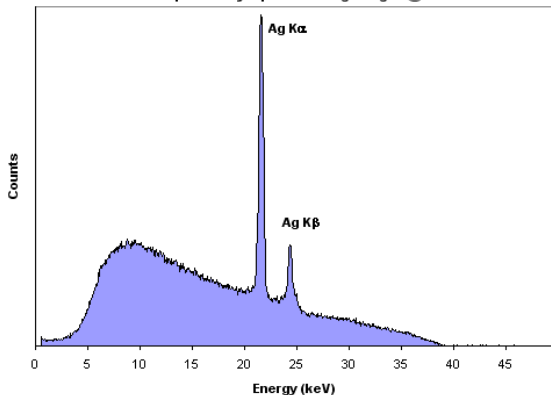
X-ray tube spectrum

Mini-X Output X-Ray Spectrum: Ag Target @ 40 kV



X-ray tube spectrum

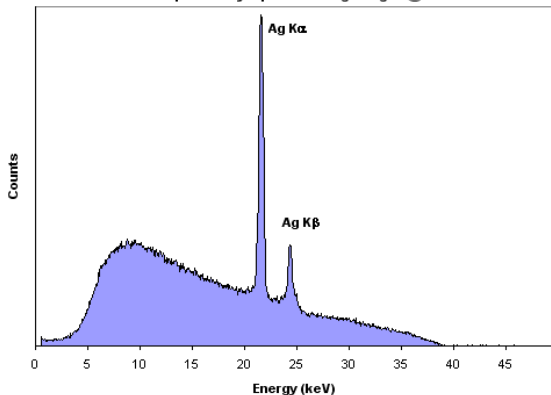
Mini-X Output X-Ray Spectrum: Ag Target @ 40 kV



- Minimum wavelength (maximum energy) set by accelerating potential

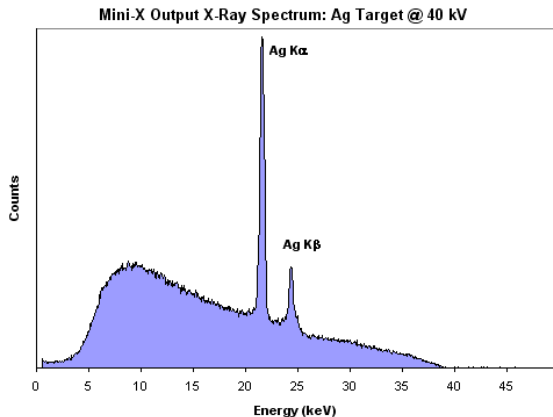
X-ray tube spectrum

Mini-X Output X-Ray Spectrum: Ag Target @ 40 kV



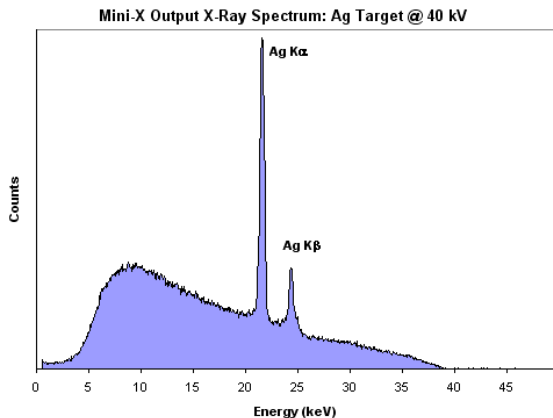
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X-ray tube spectrum



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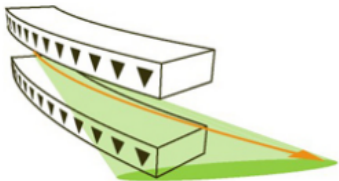
X-ray tube spectrum



- Minimum wavelength (maximum energy) set by accelerating potential
- Bremsstrahlung radiation provides smooth background (charged particle deceleration)
- Highest intensity at the characteristic fluorescence emission energy of the anode material
- Unpolarized, incoherent x-rays emitted in all directions from anode surface, must be collimated with slits

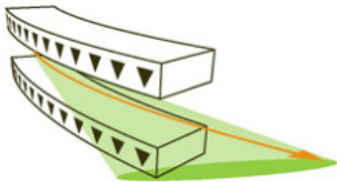
Synchrotron sources

Bending magnet



Synchrotron sources

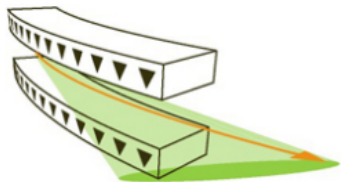
Bending magnet



- Wide horizontal beam

Synchrotron sources

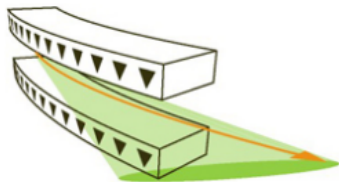
Bending magnet



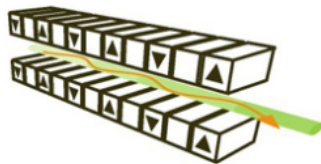
- Wide horizontal beam
- Broad spectrum to high energies

Synchrotron sources

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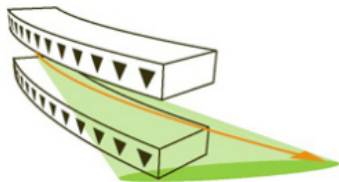
Undulator



- Wide horizontal beam
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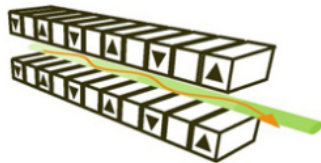
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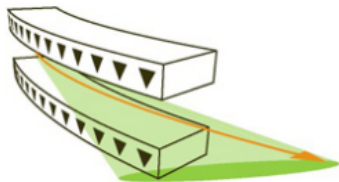
Undulator



- Highly collimated beam

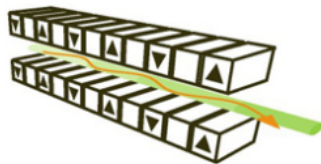
Synchrotron sources

Bending magnet



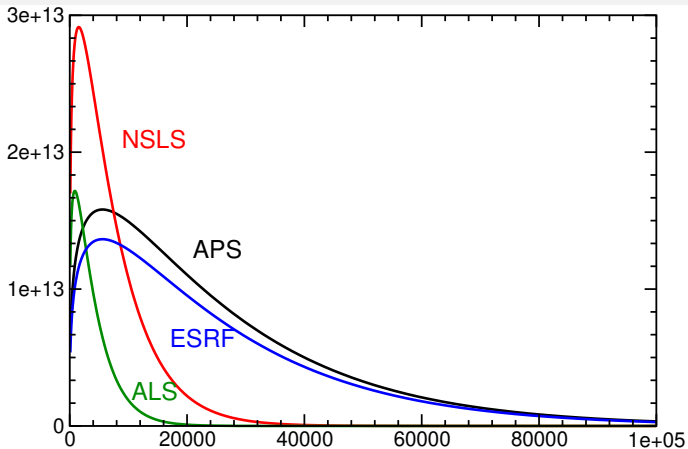
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Undulator

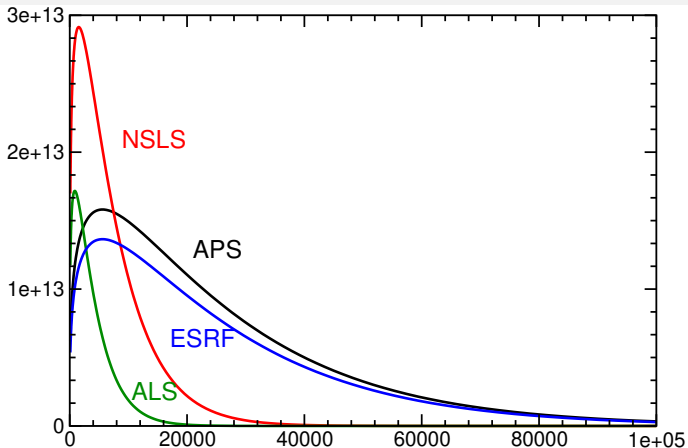


- Highly collimated beam
- Highly peaked spectrum with harmonics

Bending magnet spectra

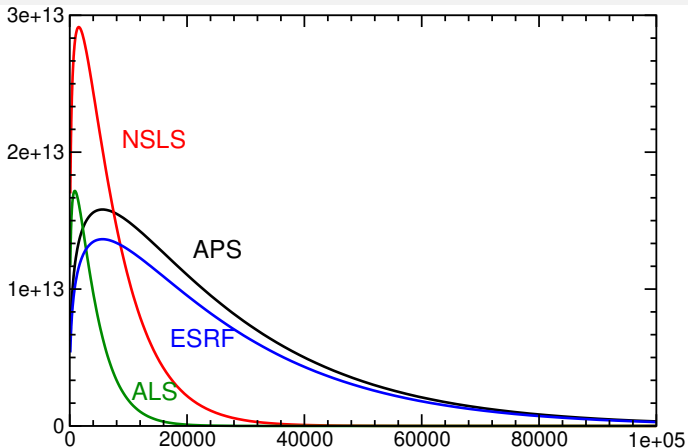


Bending magnet spectra



Lower energy sources, such as NSLS have lower peak energy and higher intensity at the peak.

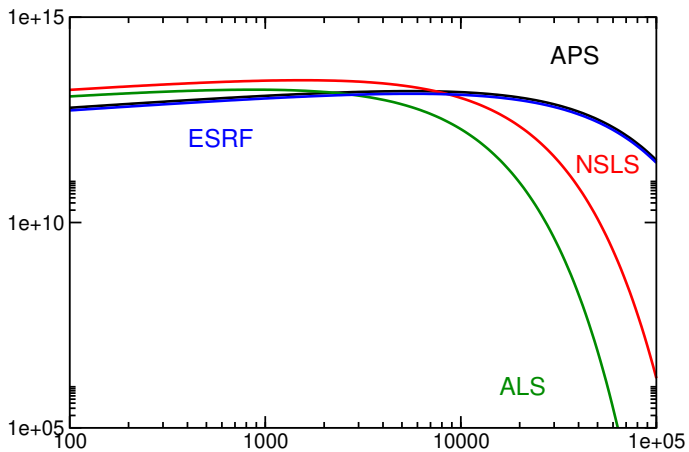
Bending magnet spectra



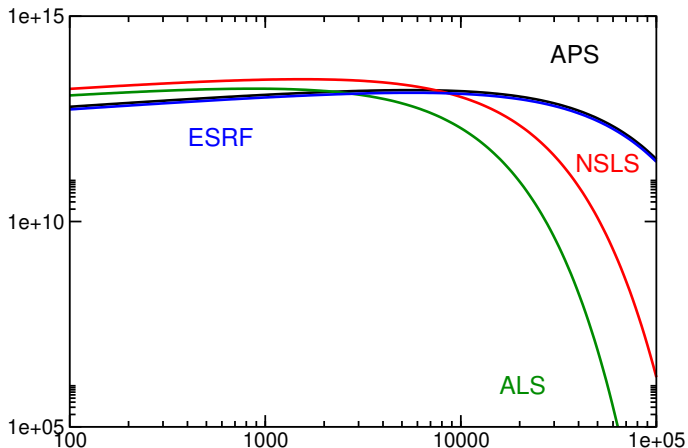
Lower energy sources, such as NSLS have lower peak energy and higher intensity at the peak.

Higher energy sources, such as APS have higher energy spectrum and are only off by a factor of 2 intensity at low energy.

Bending magnet spectra

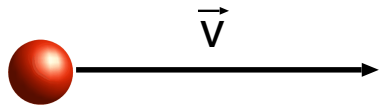


Bending magnet spectra

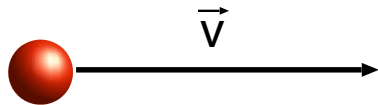


Logarithmic scale shows clearly how much more energetic and intense the bending magnet sources at the 6 GeV and 7 GeV sources are.

Review of special relativity

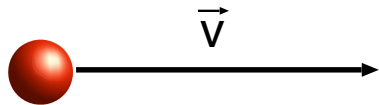


Review of special relativity



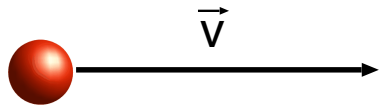
$$\beta = \frac{v}{c}$$

Review of special relativity



$$\beta = \frac{v}{c} \quad \gamma = \sqrt{\frac{1}{1 - \beta^2}}$$

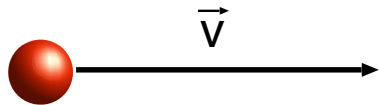
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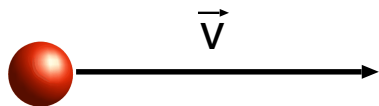


$$\beta = \frac{v}{c} \quad \gamma = \sqrt{\frac{1}{1 - \beta^2}}$$

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$$\beta = \sqrt{1 - \frac{1}{\gamma^2}}$$

Review of special relativity



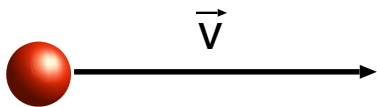
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$$\beta = \sqrt{1 - \frac{1}{\gamma^2}} \rightarrow \beta \approx 1 - \frac{1}{2\gamma^2}$$

use binomial expansion since $1/\gamma^2 \ll 1$

Review of special relativity



Let's calculate these quantities for an electron at NSLS and APS

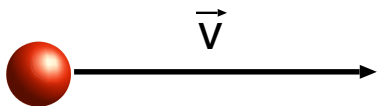
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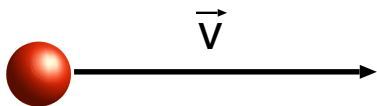
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$$m_e = 0.511 \text{ MeV}/c^2$$

Review of special relativity



$$\beta = \frac{v}{c} \quad \gamma = \sqrt{\frac{1}{1 - \beta^2}}$$

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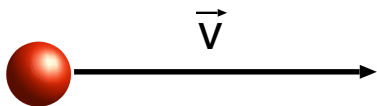
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$$\text{NSLS: } E = 1.5 \text{ GeV}$$

$$\gamma = \frac{1.5 \times 10^9}{0.511 \times 10^6} = 2935$$

Review of special relativity



$$\beta = \frac{v}{c} \quad \gamma = \sqrt{\frac{1}{1 - \beta^2}}$$

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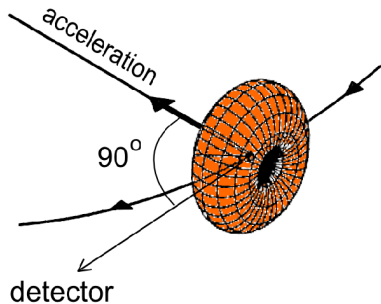
$$\gamma = \frac{1.5 \times 10^9}{0.511 \times 10^6} = 2935$$

APS: $E = 7.0 \text{ GeV}$

$$\gamma = \frac{7.0 \times 10^9}{0.511 \times 10^6} = 13700$$

“Headlight” effect

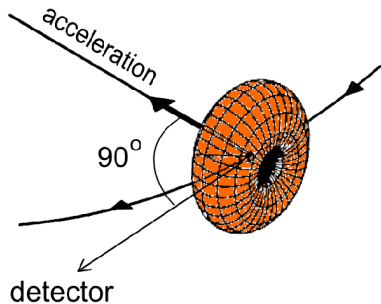
In electron rest frame:



emission is symmetric about the axis of the acceleration vector

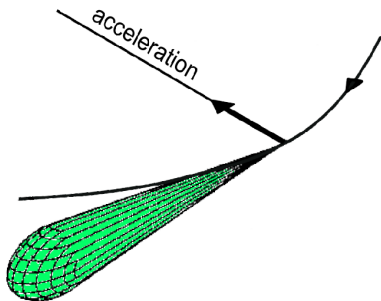
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In electron rest frame:



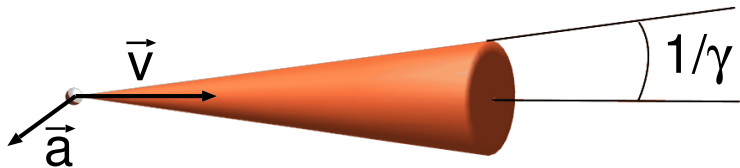
emission is symmetric about the axis of the acceleration vector

In lab frame:



emission is pushed into the direction of motion of the electron

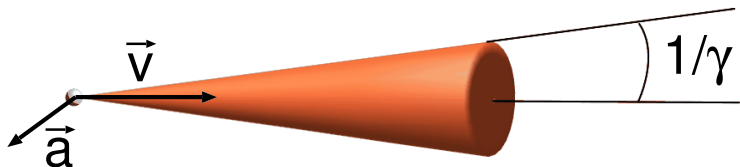
Relativistic emission



the electron is in constant transverse acceleration due to the Lorentz force from the magnetic field of the bending magnet

$$\vec{F} = e\vec{v} \times \vec{B} = m_e\vec{a}$$

Relativistic emission

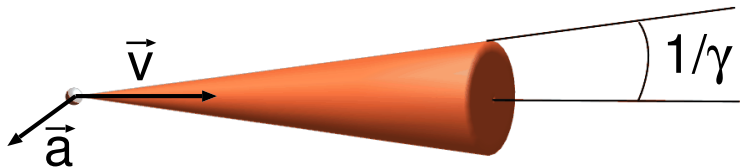


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Relativistic emission



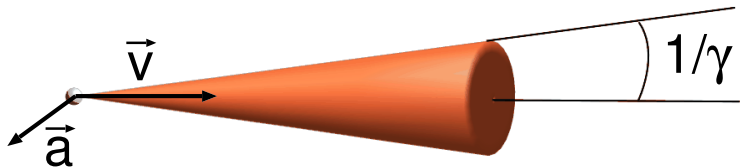
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Relativistic emission



the electron is in constant transverse acceleration due to the Lorentz force from the magnetic field of the bending magnet

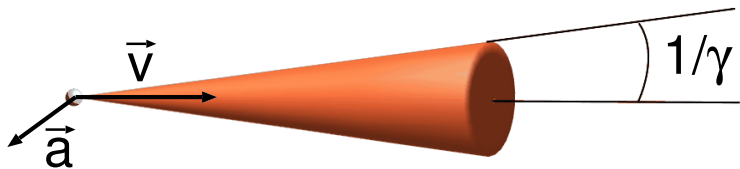
$$\vec{F} = e\vec{v} \times \vec{B} = m_e\vec{a}$$

the aperture angle of the radiation cone is $1/\gamma$

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$$E_{max} \approx \gamma^3 \omega_0$$

Relativistic emission



the electron is in constant transverse acceleration due to the Lorentz force from the magnetic field of the bending magnet

$$\vec{F} = e\vec{v} \times \vec{B} = m_e \vec{a}$$

the aperture angle of the radiation cone is $1/\gamma$

the angular frequency of the electron in the ring is $\omega_0 \approx 10^6$ and the cutoff energy for emission is

$$E_{max} \approx \gamma^3 \omega_0$$

for the APS, with $\gamma \approx 10^4$ we have

$$E_{max} \approx (10^4)^3 \cdot 10^6 = 10^{18}$$

Flux and brilliance

There are a number of important quantities which are relevant to the quality of an x-ray source:

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photon flux

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There are a number of important quantities which are relevant to the quality of an x-ray source:

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$$brilliance = \frac{flux \text{ [photons/s]}}{divergence \text{ [mrad}^2\text{]}}$$

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$$\textit{brilliance} = \frac{\textit{flux} [\text{photons/s}]}{\textit{divergence} [\text{mrad}^2] \cdot \textit{source size} [\text{mm}^2]}$$

Flux and brilliance

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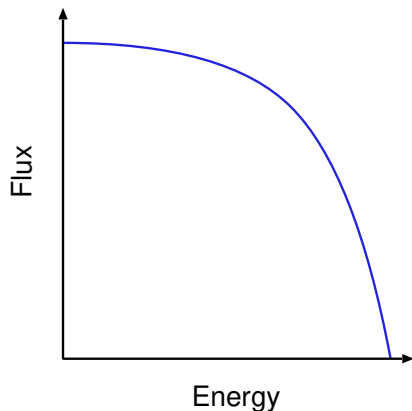
$$\mathit{brilliance} = \frac{\mathit{flux} \text{ [photons/s]}}{\mathit{divergence} \text{ [mrad}^2\text{]} \cdot \mathit{source size} \text{ [mm}^2\text{]} \text{ [0.1\% bandwidth]}}$$

Computing brilliance

$$brilliance = \frac{flux \text{ [photons/s]}}{divergence \text{ [mrad}^2\text{]} \cdot source \text{ size [mm}^2\text{]} \cdot [0.1\% \text{ bandwidth}]}$$

Computing brilliance

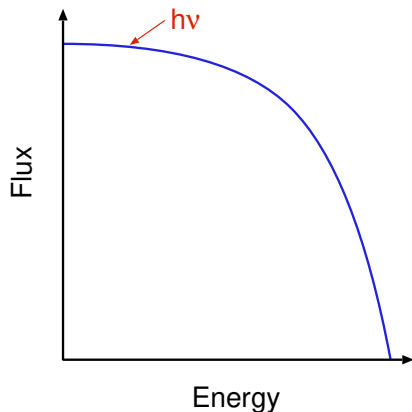
$$\text{brilliance} = \frac{\text{flux} [\text{photons/s}]}{\text{divergence} [\text{mrad}^2] \cdot \text{source size} [\text{mm}^2] \cdot [0.1\% \text{ bandwidth}]}$$



For a specific photon flux distribution, we would normally integrate to get the total flux.

Computing brilliance

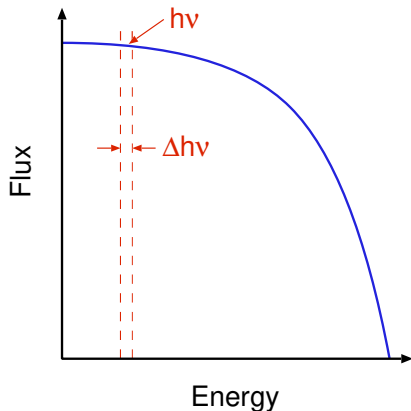
$$\text{brilliance} = \frac{\text{flux [photons/s]}}{\text{divergence [mrad}^2\text{]} \cdot \text{source size [mm}^2\text{]} \cdot [0.1\% \text{ bandwidth}]}$$



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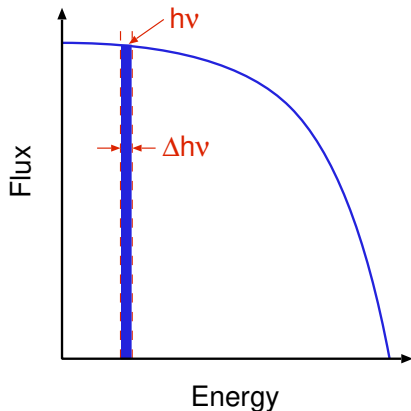


For a specific photon flux distribution, we would normally integrate to get the total flux. But this ignores that most experiments are only interested in a specific energy $h\nu$.

Take a bandwidth $\Delta h\nu = h\nu/1000$, which is about 10 times wider than the bandwidth of the typical monochromator.

Computing brilliance

$$\text{brilliance} = \frac{\text{flux} [\text{photons/s}]}{\text{divergence} [\text{mrad}^2] \cdot \text{source size} [\text{mm}^2] \cdot [0.1\% \text{ bandwidth}]}$$



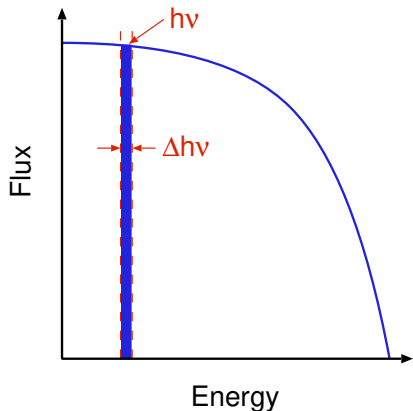
For a specific photon flux distribution, we would normally integrate to get the total flux. But this ignores that most experiments are only interested in a specific energy $h\nu$.

Take a bandwidth $\Delta h\nu = h\nu/1000$, which is about 10 times wider than the bandwidth of the typical monochromator.

Compute the **integrated photon flux in that bandwidth**.

Computing brilliance

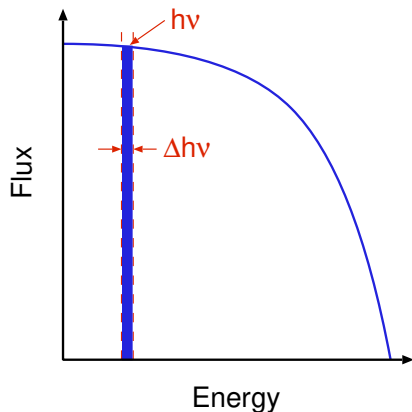
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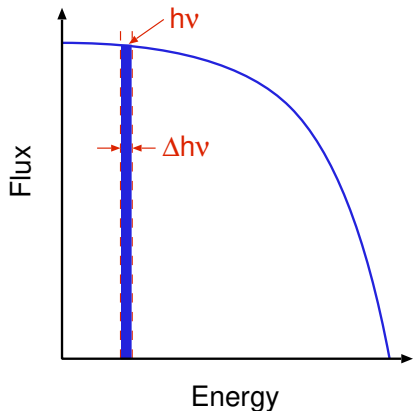


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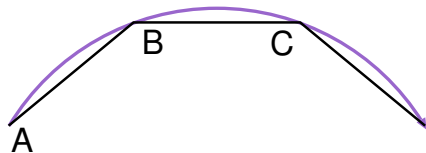
$$\alpha \approx x/z \quad \beta \approx y/z,$$

where z is the distance from the source over which there is a lateral spread x and y in each direction

Segmented arc approximation

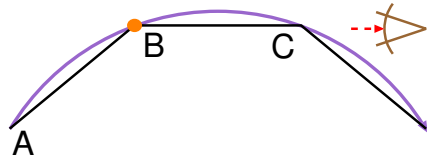


Segmented arc approximation



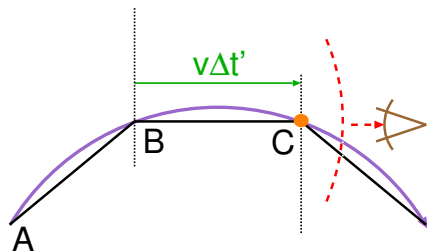
- Approximate the electron's path as a series of segments

Segmented arc approximation



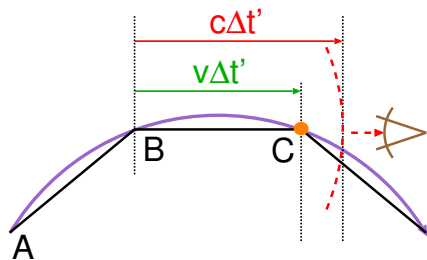
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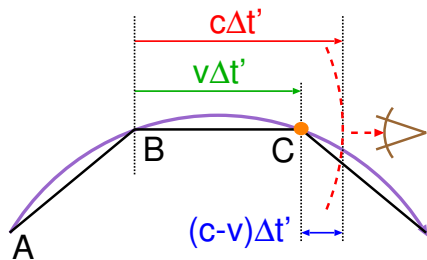
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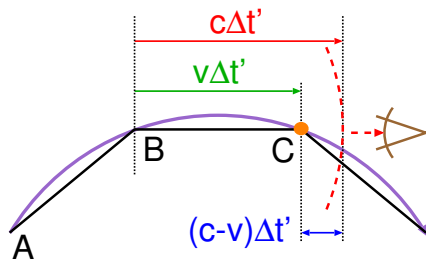
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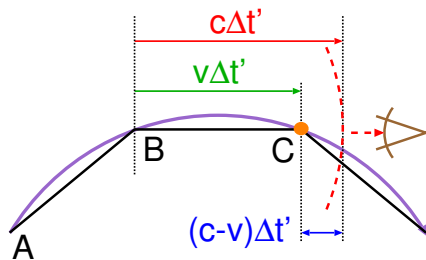


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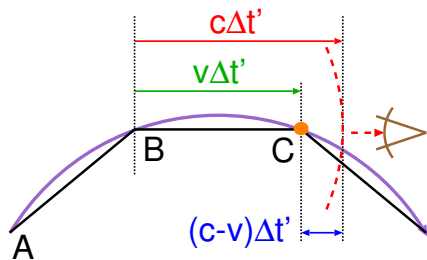
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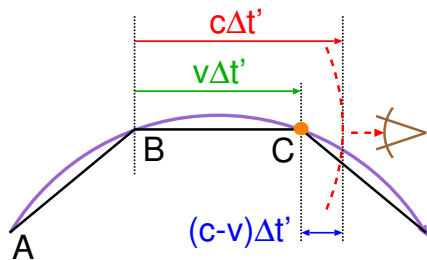
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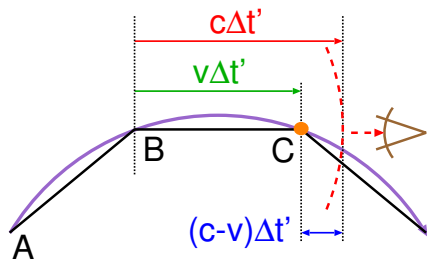
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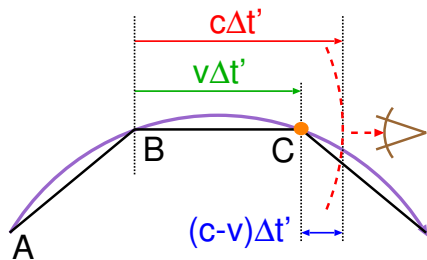
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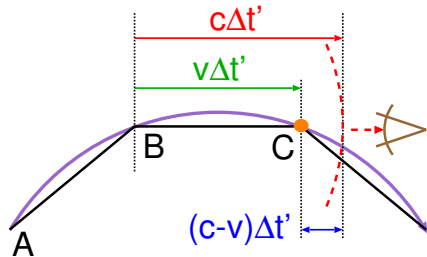
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Doppler compression



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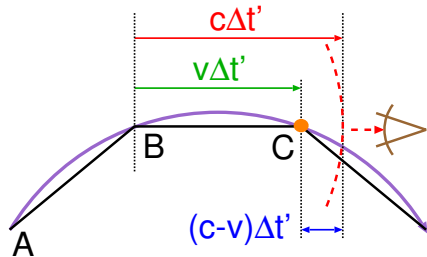
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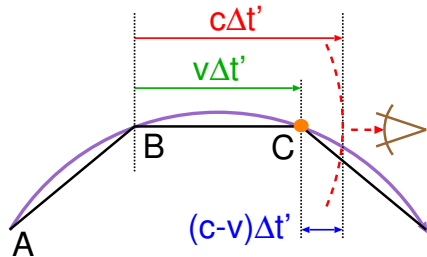
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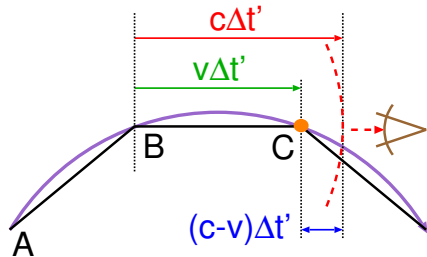
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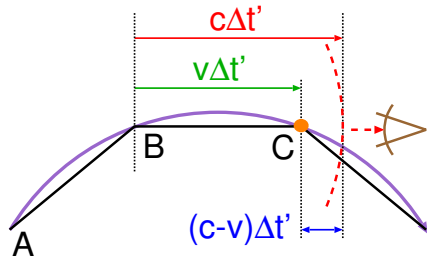
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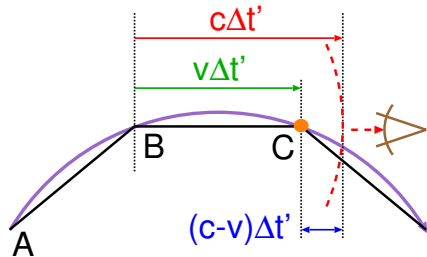
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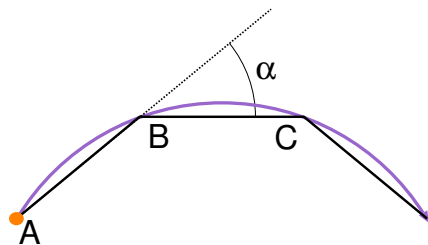
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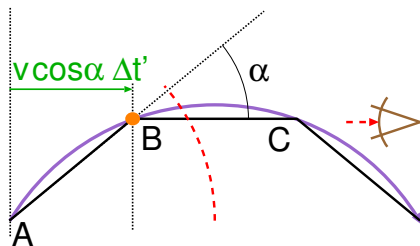
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Off-axis emission



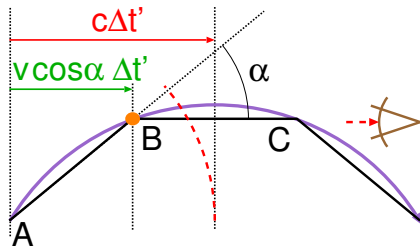
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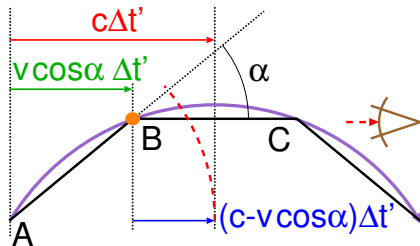
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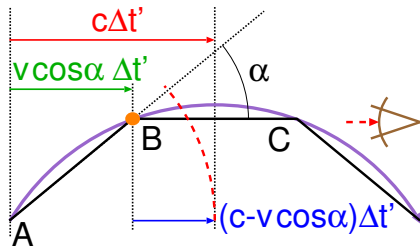


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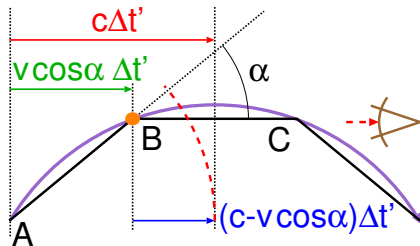


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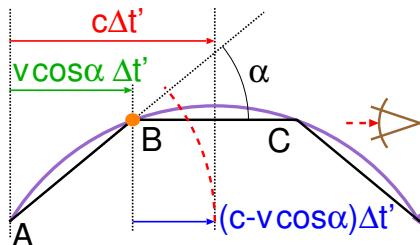
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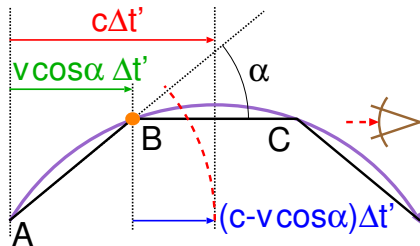
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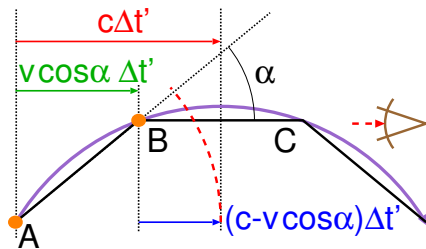
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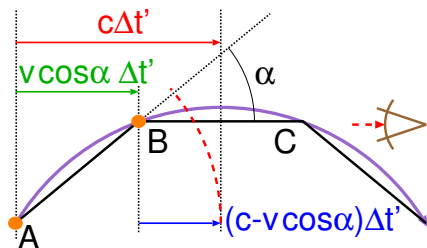
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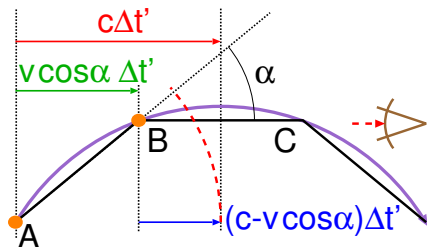
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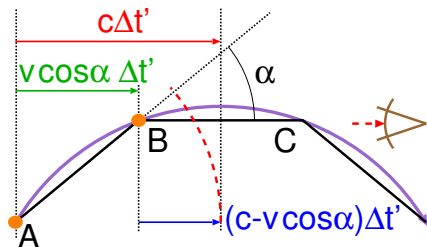


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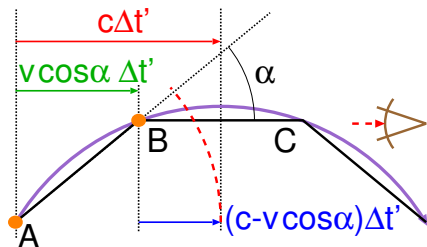
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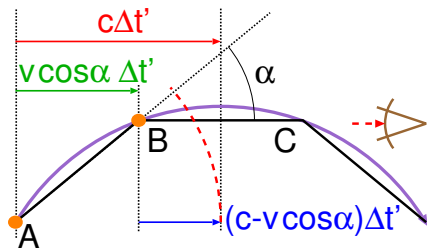
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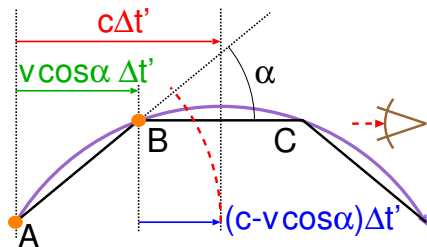
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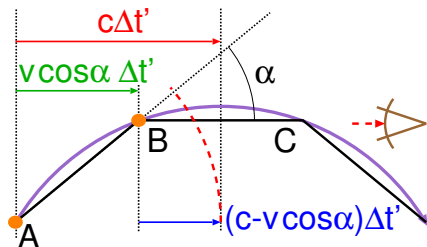
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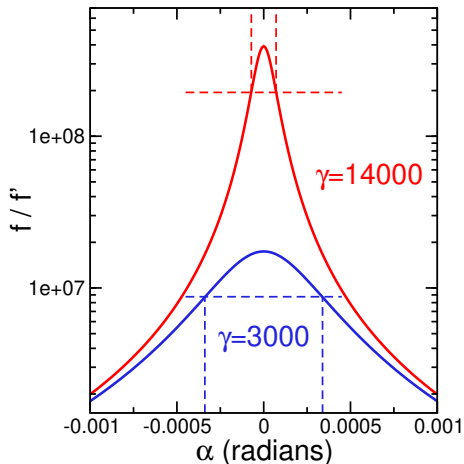
$$\frac{\Delta t}{\Delta t'} \approx \frac{\alpha^2}{2} + \frac{1}{2\gamma^2} = \frac{1 + \alpha^2 \gamma^2}{2\gamma^2}$$

called the time compression ratio.

Radiation opening angle

The Doppler shift is defined in terms of the time compression ratio

$$\frac{f}{f'} = \frac{\Delta t'}{\Delta t} = \frac{2\gamma^2}{1 + \alpha^2\gamma^2}$$

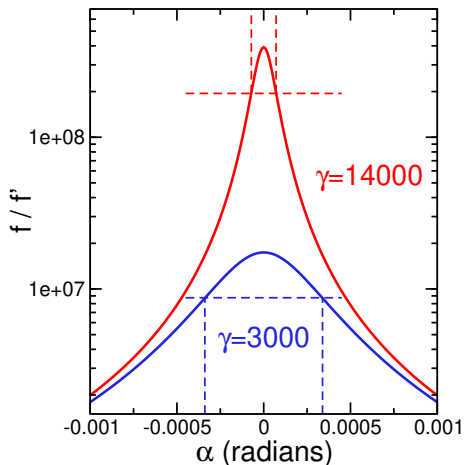


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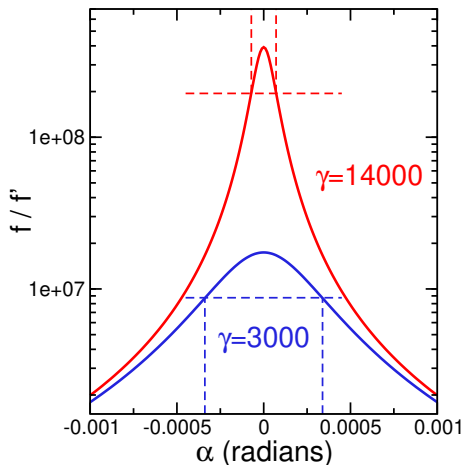


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- The highest energy emitted radiation appears within a cone of half angle $1/\gamma$