Final presentation

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Homework Assignment #04: Chapter 4: 2, 4, 6, 7, 10 due Monday, October 24, 2016

- Final presentation
- Ewald sphere (continued)
- Modulated Structures
- Crystal Truncation Rods
- Diffuse Scattering
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Homework Assignment #04:

Chapter 4: 2, 4, 6, 7, 10

due Monday, October 24, 2016

Homework Assignment #05:

Chapter 5: 1, 3, 7, 9, 10

due Wednesday, November 2, 2016

No class on Wednesday, November 9, 2016

Choose paper for presentation

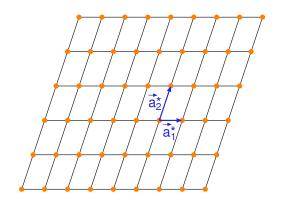
- 1 Choose paper for presentation
- Olear it with me!

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- 2 Clear it with me!
- 3 Do some background research on the technique

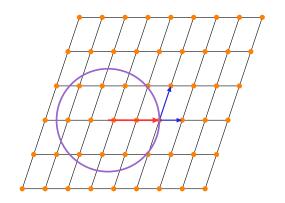
2 / 16

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- 4 Prepare a 15 minute presentation
- **5** Be ready for questions!

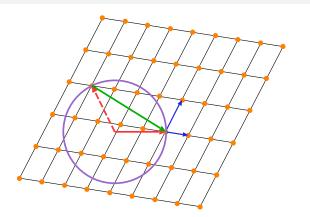


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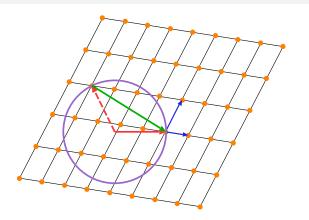
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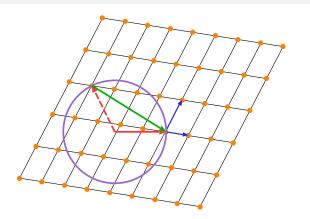


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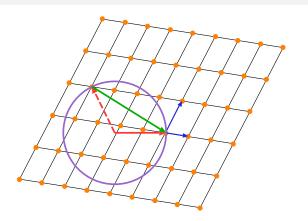


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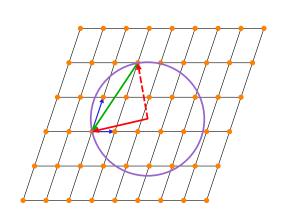
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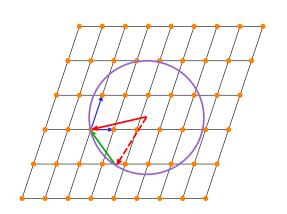
$$\vec{G}_{hlk} = h\vec{a}_1^* + k\vec{a}_2^*$$

Ewald construction



It is often more convenient to visualize the Ewald sphere by keeping the reciprocal lattice fixed and "rotating" the incident beam to visualize the scattering geometry.

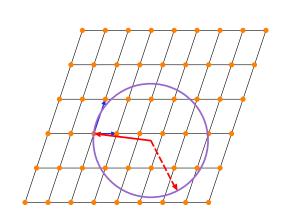
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In directions of \vec{k}' (detector position) where there is no reciprocal lattice point, there can be no diffraction peak.

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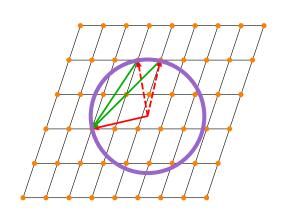


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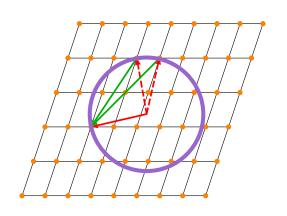
If the crystal is rotated slightly with respect to the incident beam, \vec{k} , there may be no Bragg reflections possible at all.

Polychromatic radiation



If $\Delta \vec{k}$ is large enough, there may be more than one reflection lying on the Ewald sphere.

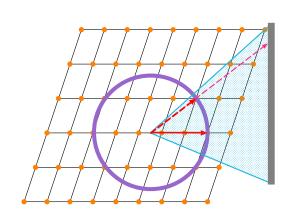
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With an area detector, there may then be multiple reflections appearing for a particular orientation (very common with protein crystals where the unit cell is very large).

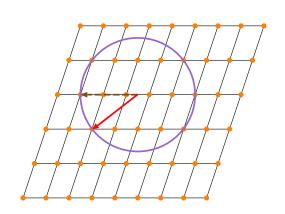
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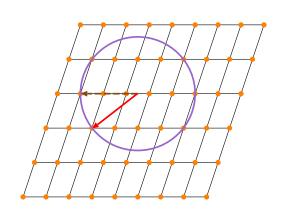
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In protein crystallography, the area detector is in a fixed location with respect to the incident beam and the crystal is rotated on a spindle so that as Laue conditions are met, spots are produced on the detector at the diffraction angle

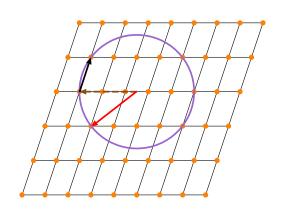


If more than one reciprocal lattice point is on the Ewald sphere, scattering can occur internal to the crystal.



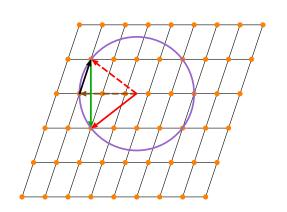
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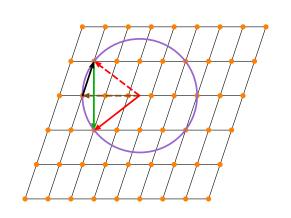
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If more than one reciprocal lattice point is on the Ewald sphere, scattering can occur internal to the crystal.

The xrays are first scattered along \vec{k}_{int} then along the reciprocal lattice vector which connects the two points on the Ewald sphere, \vec{G} and to the detector at \vec{k}' .

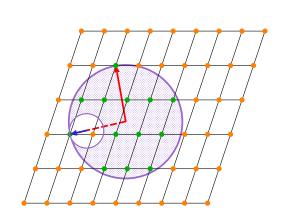


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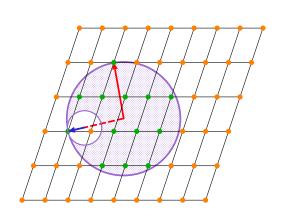
This is the cause of monochromator glitches which sometimes remove intensity but can also add intensity to the reflection the detector is set to measure.

Laue diffraction



The Laue diffraction technique uses a wide range of radiation from \vec{k}_{min} to \vec{k}_{max}

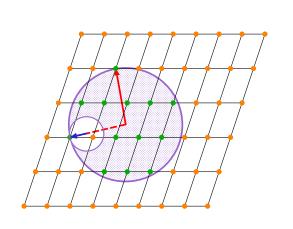
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These define two Ewald spheres and a volume between them such that any reciprocal lattice point which lies in the volume will meet the Laue condition for reflection.

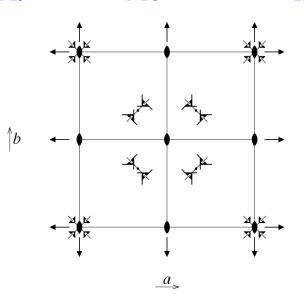
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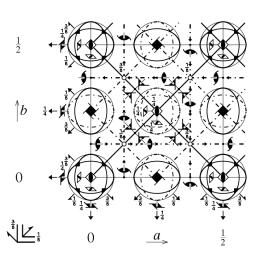
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This technique is useful for taking data on crystals which are changing or may degrade in the beam with a single shot of x-rays on a 2D detector.



- 1 x, y, z
 - 2 $x, \overline{y}, \overline{z}$
 - 3 \overline{x} , y, \overline{z}
 - 4 \bar{x} , \bar{y} , z
 - 5 z, x, y
 - 6 \overline{z} , \overline{x} , y
 - 7 $z, \overline{x}, \overline{y}$
 - 8 \overline{z} , x, \overline{y}
 - 9 y, z, x
 - 10 \overline{y} , z, \overline{x}
 - 11 \overline{y} , \overline{z} , x
 - 12 $y, \overline{z}, \overline{x}$





	- A-
	2KC
1 x, y, z	25 $\frac{1}{4} - x$, $\frac{1}{4} - y$, $\frac{1}{4} - z$
2 $x, \overline{y}, \overline{z}$	$26 \frac{1}{4} - x, \frac{1}{4} + y, \frac{1}{4} + z$
$3 \ \overline{x}, y, \overline{z}$	27 $\frac{1}{4} + x$, $\frac{1}{4} - y$, $\frac{1}{4} + z$
$4 \overline{x}, \overline{y}, z$	28 $\frac{1}{4}$ + x, $\frac{1}{4}$ + y, $\frac{1}{4}$ - z
5 z, x, y	29 $\frac{1}{4}$ - z, $\frac{1}{4}$ - x, $\frac{1}{4}$ - y
6 \overline{z} , \overline{x} , y	30 $\frac{1}{4}$ + z, $\frac{1}{4}$ + x, $\frac{1}{4}$ - y
7 $z, \overline{x}, \overline{y}$	31 $\frac{1}{4}$ - z, $\frac{1}{4}$ + x, $\frac{1}{4}$ + y
8 \overline{z} , x , \overline{y}	32 $\frac{1}{4}$ + z, $\frac{1}{4}$ - x, $\frac{1}{4}$ + y
9 y, z, x	33 $\frac{1}{4} - y$, $\frac{1}{4} - z$, $\frac{1}{4} - x$
10 ȳ, z, x̄	$34 \frac{1}{4} + y, \frac{1}{4} - z, \frac{1}{4} + x$
11 \overline{y} , \overline{z} , x	35 $\frac{1}{4}$ + y, $\frac{1}{4}$ + z, $\frac{1}{4}$ - x
12 y, ₹, ₹	36 $\frac{1}{4}$ - y, $\frac{1}{4}$ + z, $\frac{1}{4}$ + x
13 $\frac{1}{4} + x$, $\frac{1}{4} - z$, $\frac{1}{4} + y$	
$14 \frac{1}{4} + x, \frac{1}{4} + z, \frac{1}{4} - y$	
15 $\frac{1}{4} - x$, $\frac{1}{4} - z$, $\frac{1}{4} - y$	
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23 $\frac{1}{4} - y$, $\frac{1}{4} - x$, $\frac{1}{4} - z$	
$24 \ \frac{1}{4} + y, \frac{1}{4} + x, \frac{1}{4} - z$	48 y, x, z

 $+\,(0,\!\tfrac{1}{2},\!\tfrac{1}{2}),\,(\tfrac{1}{2},\!0,\!\tfrac{1}{2}),\,(\tfrac{1}{2},\!\tfrac{1}{2},\!0)$





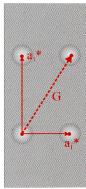
Wyckoff Positions of Group 195 (P23)

Multiplicity	Wyckoff letter	Site symmetry	Coordinates
12	j	1	(x,y,z) (-x,-y,z) (-x,y,-z) (x,-y,-z) (z,x,y) (z,-x,-y) (-z,-x,y) (-z,x,-y) (y,z,x) (-y,z,-x) (y,-z,-x) (-y,-z,x)
6	i	2	(x,1/2,1/2) (-x,1/2,1/2) (1/2,x,1/2) (1/2,-x,1/2) (1/2,1/2,x) (1/2,1/2,-x)
6	h	2	(x,1/2,0) (-x,1/2,0) (0,x,1/2) (0,-x,1/2) (1/2,0,x) (1/2,0,-x)
6	g	2	(x,0,1/2) (-x,0,1/2) (1/2,x,0) (1/2,-x,0) (0,1/2,x) (0,1/2,-x)
6	f	2	(x,0,0) (-x,0,0) (0,x,0) (0,-x,0) (0,0,x) (0,0,-x)
4	е	.3.	(x,x,x) (-x,-x,x) (-x,x,-x) (x,-x,-x)
3	d	222	(1/2,0,0) (0,1/2,0) (0,0,1/2)
3	С	222	(0,1/2,1/2) (1/2,0,1/2) (1/2,1/2,0)
1	b	23.	(1/2,1/2,1/2)
1	а	23.	(0,0,0)

Wyckoff Positions of Group 227 (Fd-3m) [origin choice 1]

Multiplicity	Wyckoff	Site	Coordinates		
Multiplicity	letter	symmetry	(0,0,0) + (0,1/2,1/2) + (1/2,0,1/2) + (1/2,1/2,0) +		
192	i	1	$ \begin{array}{llllllllllllllllllllllllllllllllllll$		
96	h	2	[118 y, y+14] (718 , y+1/2 , y+3/4) (318 y+1/2 y+3/4) (518 , yy+1/4) (y+1/4 , 118 y) (y+3/4 , 718 , y+1/2) (y+3/4 , 318 y+1/2) (y+1/4 , 518 , y) [y+3/4 , 718 , y+1/2 , y+3/4 , 318 , y+1/2 , y+3/4 , y+1/2 ,		
96	g	m	[(X,Z) (X,x×112,Z+12) (x+12,X+12,z) (k+12,x+2x+12) (2(X,X) (2x+12,X+2x+12) (2-2,X+12,X+12) (2-2,X+12,X+12) (2-2,X+12,X+12) (2-2,X+12,X+12) (2-2,X+12,X+12) (2-2,X+12,X+12) (2-2,X+12) (2-2,		
48	f	2.m m	(x,0,0) (-x,1/2,1/2) (0,x,0) (1/2,-x,1/2) (0,0,x) (1/2,-x,1/2) (3/4,x+1/4,3/4) (1/4,-x+1/4,1/4) (x+3/4,1/4,3/4) (-x+3/4,3/4,1/4) (3/4,1/4,-x+3/4) (1/4,3/4) (x+3/4,3/4,1/4)		
32	е	.3m	(x,x) (-x,-x+1/2,x+1/2) (-x+1/2,x+1/2,x) (x+1/2,x,-x+1/2) (x+3/4,x+1/4,-x+3/4) (-x+1/4,-x+1/4,-x+1/4) (x+1/4,-x+3/4,x+3/4) (-x+3/4,x+3/4,x+1/4)		
16	d	3m	(5/8,5/8,5/8) (3/8,7/8,1/8) (7/8,1/8,3/8) (1/8,3/8,7/8)		
16	С	3m	(1/8,1/8,1/8) (7/8,3/8,5/8) (3/8,5/8,7/8) (5/8,7/8,3/8)		
8	b	-43m	(1/2,1/2,1/2) (1/4,3/4,1/4)		
8	а	-43m	(0,0,0) (3/4,1/4,3/4)		

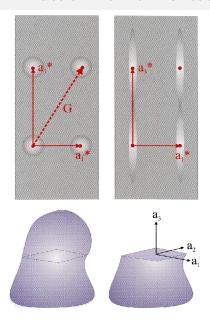
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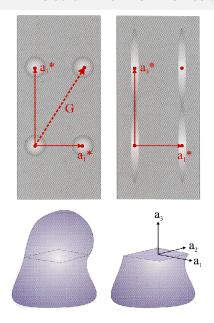
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The scattering intensity can be obtained by treating the charge distribution as a convolution of an infinite sample with a step function in the z-direction.

The scattering amplitude F^{CTR} along a crystal truncation rod is given by summing an infinite stack of atomic layers, each with scattering amplitude $A(\vec{Q})$.

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$$= \frac{A(\vec{Q})}{1 - e^{iQ_z a_3}}$$

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or, in terms of the momentum transfer along the z-axis, $Q_z=2\pi I/a_3$

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$$Q_z=2\pi I/a_3$$

since the intensity is the square of the scattering factor

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$$I^{CTR} = \frac{\left| A(\vec{Q}) \right|^2}{4 \sin^2 \left(Q_z a_3 / 2 \right)}$$

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$$= \frac{\left| A(\vec{Q}) \right|^2}{4 \sin^2 (\pi I + q_z a_3/2)}$$

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$$= \frac{\left| A(\vec{Q}) \right|^2}{4 \sin^2 (q_z a_3/2)}$$

$$\approx \frac{\left| A(\vec{Q}) \right|^2}{4 (q_z a_3/2)^2}$$

$$I^{CTR} = \frac{\left| A(\vec{Q}) \right|^2}{4 \sin^2 (Q_z a_3 / 2)}$$

$$= \frac{\left| A(\vec{Q}) \right|^2}{4 \sin^2 (\pi I + q_z a_3 / 2)}$$

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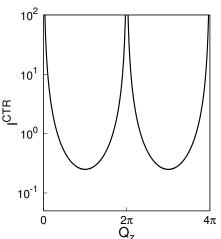
$$\approx \frac{\left| A(\vec{Q}) \right|^2}{4(q_z a_3 / 2)^2} = \frac{\left| A(\vec{Q}) \right|^2}{q_z^2 a_3^2}$$

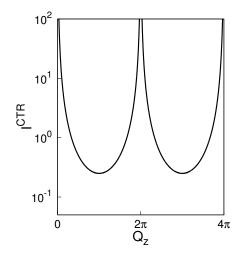
$$I^{CTR} = \frac{\left| A(\vec{Q}) \right|^2}{4 \sin^2 (Q_z a_3 / 2)}$$

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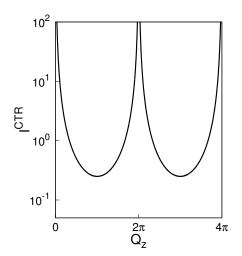
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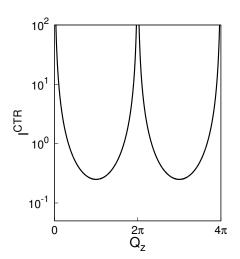




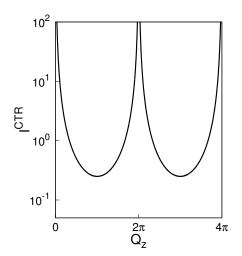
$$F^{CTR} = A(\vec{Q}) \sum_{j=0}^{\infty} e^{iQ_z a_3 j}$$



$$F^{CTR} = A(\vec{Q}) \sum_{j=0}^{\infty} e^{iQ_z a_3 j} e^{-\beta j}$$

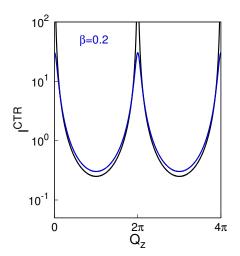


$$F^{CTR} = A(\vec{Q}) \sum_{j=0}^{\infty} e^{iQ_z a_3 j} e^{-\beta j}$$
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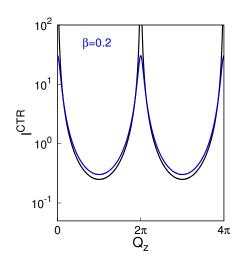


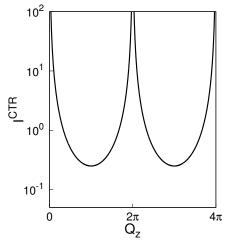
Absorption effects can be included as well by adding a term for each layer penetrated

$$F^{CTR} = A(\vec{Q}) \sum_{j=0}^{\infty} e^{iQ_z a_3 j} e^{-\beta j}$$

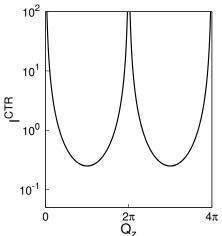
$$= \frac{A(\vec{Q})}{1 - e^{iQ_z a_3} e^{-\beta}}$$

This removes the infinity and increases the scattering profile of the crystal truncation rod



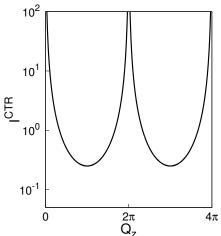


$$F^{total} = F^{CTR} + F^{top\ layer}$$



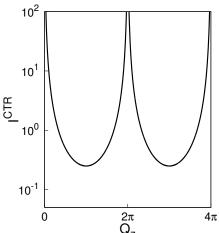
$$F^{total} = F^{CTR} + F^{top \ layer}$$

$$= \frac{A(\vec{Q})}{1 - e^{i2\pi l}}$$



$$F^{total} = F^{CTR} + F^{top\ layer}$$

$$= \frac{A(\vec{Q})}{1 - e^{i2\pi l}} + A(\vec{Q})e^{-i2\pi(1+z_0)l}$$

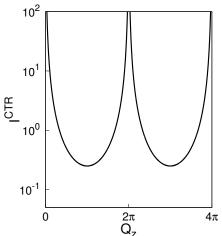


The CTR profile is sensitive to the termination of the surface. This makes it an ideal probe of electron density of adsorbed species or single atom overlayers.

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where z_0 is the relative displacement of the top layer from the bulk lattice spacing a_3

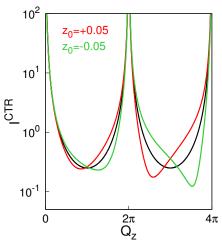


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where z_0 is the relative displacement of the top layer from the bulk lattice spacing a_3 This effect gets larger for larger momentum transfers

