• Beam time at MRCAT

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Homework Assignment #03: Chapter3: 1, 3, 4, 6, 8 due Wednesday, October 05, 2016

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Friday, October 28, 2016 - bending magnet line

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I will try to get BioCAT time as well for those interested

Writing a General User Proposal

- 1 Log into the APS site
- 2 Start a general user proposal
- 8 Add an Abstract
- 4 Choose a beam line
- **5** Answer the 6 questions

Treat the multilayer as a stratified medium on top of an infinitely thick substrate.

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The recursive relation can be seen from the calculation of reflectivity of the next layer up

$$r_{N-2,N-1} = \frac{r'_{N-2,N-1} + r_{N-1,N}p_{N-1}^2}{1 + r'_{N-2,N-1}r_{N-1,N}p_{N-1}^2}$$

PHYS 570 - Spring 2016





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Kinematical - Parratt comparison



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Peaks in kinematical calculation are somewhat higher reflectivity than true value.

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The reflectivity can be computed as the superposition of the reflectivity of a series of infinitesmal slabs of thickness dz at a depth z.

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Calculating the full reflection coefficient relative to the Fresnel reflection coefficient

$$\frac{R(Q)}{R_F(Q)} = \left| \int_{-\infty}^{\infty} \left(\frac{df}{dz} \right) e^{iQz} dz \right|^2$$

C. Segre (IIT)

PHYS 570 - Spring 2016

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$$Q = k \sin \theta, \quad Q' = k' \sin \theta'$$

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$$r_{S} = -r_{0}\rho \frac{1}{iQ_{z}}\int_{S}e^{i\vec{Q}\cdot\vec{r}}dxdy$$

Evaluation of surface integral

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$$ec{Q}\cdotec{r}=Q_zh(x,y)+Q_xx+Q_yy$$

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The side surfaces of the volume do not contribute to this integral as they are along the \hat{z} direction, but we can also choose the thickness of the slab such that the lower surface will not contribute either.

Thus, the integral need only be evaluated over the top, rough surface whose variation we characterize by the function h(x, y)

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The actual scattering cross section is the square of this integral

$$\frac{d\sigma}{d\Omega} = \left(\frac{r_0\rho}{Q_z}\right)^2 \int_S \int_{S'} e^{iQ_z(h(x,y) - h(x',y'))} e^{iQ_x(x-x')} e^{iQ_y(y-y')} dx dy dx' dy'$$

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Finally, it is assumed that the statistics of the height variation are Gaussian and

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Unbounded correlations - limiting cases

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C. Segre (IIT)

PHYS 570 - Spring 2016

September 21, 2016 18 / 19

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$$\begin{pmatrix} \frac{d\sigma}{d\Omega} \end{pmatrix} = \left(\frac{r_0\rho}{Q_z}\right)^2 \frac{A_0}{\sin\theta_1} e^{-Q_z^2\sigma^2} \int e^{Q_z^2C(x,y)} e^{iQ_xx} e^{iQ_yy} dxdy$$

$$= \left(\frac{r_0\rho}{Q_z}\right)^2 \frac{A_0}{\sin\theta_1} e^{-Q_z^2\sigma^2} \int \left[e^{Q_z^2C(x,y)} - 1 + 1\right] e^{iQ_xx} e^{iQ_yy} dxdy$$

$$= \left(\frac{d\sigma}{d\Omega}\right)_{\text{Fresnel}} e^{-Q_z^2\sigma^2} + \left(\frac{r_0\rho}{Q_z}\right)^2 \frac{A_0}{\sin\theta_1} e^{-Q_z^2\sigma^2} F_{\text{diffuse}}(\vec{Q})$$

If the correlations remain bounded as $r \to \infty$

$$g(x,y) = 2\langle h^2 \rangle - 2\langle h(0,0)h(x,y) \rangle = 2\sigma^2 - 2C(x,y)$$

where

$$C(x,y) = \sigma^2 e^{-(r/\xi)^{2h}}$$

$$\begin{pmatrix} \frac{d\sigma}{d\Omega} \end{pmatrix} = \left(\frac{r_0\rho}{Q_z}\right)^2 \frac{A_0}{\sin\theta_1} e^{-Q_z^2\sigma^2} \int e^{Q_z^2C(x,y)} e^{iQ_x x} e^{iQ_y y} dxdy$$

$$= \left(\frac{r_0\rho}{Q_z}\right)^2 \frac{A_0}{\sin\theta_1} e^{-Q_z^2\sigma^2} \int \left[e^{Q_z^2C(x,y)} - 1 + 1\right] e^{iQ_x x} e^{iQ_y y} dxdy$$

$$= \left(\frac{d\sigma}{d\Omega}\right)_{\text{Fresnel}} e^{-Q_z^2\sigma^2} + \left(\frac{r_0\rho}{Q_z}\right)^2 \frac{A_0}{\sin\theta_1} e^{-Q_z^2\sigma^2} F_{\text{diffuse}}(\vec{Q})$$

And the scattering exhibits both a specular peak, reduced by uncorrelated roughness, and diffuse scattering from the correlated portion of the surface

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