## Today's Outline - September 21, 2016

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Homework Assignment \#03:
Chapter3: 1, 3, 4, 6, 8
due Wednesday, October 05, 2016

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Two days have been set aside for our class to be at Sector 10 MRCAT at the Advanced Photon Source

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I will try to get BioCAT time as well for those interested

## Writing a General User Proposal

(1) Log into the APS site
(2) Start a general user proposal
(3) Add an Abstract
(4) Choose a beam line
(5) Answer the 6 questions

## Parratt's recursive method

Treat the multilayer as a stratified medium on top of an infinitely thick substrate.

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and the wavevector transfer in the $\mathrm{j}^{\text {th }}$ layer

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Q_{j} & =2 k_{j} \sin \alpha_{j}=2 k_{z j}
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Q_{j} & =2 k_{j} \sin \alpha_{j}=2 k_{z j} \\
& =\sqrt{Q^{2}-8 k^{2} \delta_{j}+8 i k^{2} \beta_{j}}
\end{aligned}
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## Parratt reflectivity calculation

The reflectivity from the interface between layer $j$ and $j+1$, not including multiple reflections is

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$$
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to the substrate, where multiple reflections are not present

The reflectivity from the top of the $N^{\text {th }}$ layer, including multiple reflections is

$$
r_{N, \infty}^{\prime}=\frac{Q_{N}-Q_{\infty}}{Q_{N}+Q_{\infty}}
$$ now calculated (note no prime!)

The recursive relation can be seen from the calculation of reflectivity of the next layer up

$$
r_{N-2, N-1}=\frac{r_{N-2, N-1}^{\prime}+r_{N-1, N} p_{N-1}^{2}}{1+r_{N-2, N-1}^{\prime} r_{N-1, N} p_{N-1}^{2}}
$$

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Parratt peaks shifted to slightly higher values of $Q$

Peaks in kinematical calculation are somewhat higher reflectivity than true value.

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The reflectivity can be computed as the superposition of the reflectivity of a series of infinitesmal slabs of thickness $d z$ at a depth $z$.

## Reflectivity of a graded interface

The differential reflectivity from a slab of thickness $d z$ at depth $z$ is:

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Calculating the full reflection coefficient relative to the Fresnel reflection coefficient

$$
\frac{R(Q)}{R_{F}(Q)}=\left|\int_{-\infty}^{\infty}\left(\frac{d f}{d z}\right) e^{i Q z} d z\right|^{2}
$$

## The error function - a specific case

The error function is often chosen as a model for the density gradient

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f(z)=\operatorname{erf}\left(\frac{z}{\sqrt{2} \sigma}\right)=\frac{1}{\sqrt{\pi}} \int_{0}^{z / \sqrt{2} \sigma} e^{-t^{2}} d t
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whose Fourier transform is also a Gaussian, which when squared to obtain the reflection coefficient, gives.

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R(Q)=R_{F}(Q) e^{-Q^{2} \sigma^{2}}
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$$
\begin{gathered}
R(Q)=R_{F}(Q) e^{-Q^{2} \sigma^{2}}=R_{F}(Q) e^{-Q Q^{\prime} \sigma^{2}} \\
Q=k \sin \theta, \quad Q^{\prime}=k^{\prime} \sin \theta^{\prime}
\end{gathered}
$$

## Rough surfaces

When a surface or interface is not perfectly smooth but has some roughness the reflectivity is no longer simply specular but has a non-zero diffuse component which we must include in the model.

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$$ illuminated volume is given by using Gauss' theorem.

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\int_{V}(\vec{\nabla} \cdot \vec{C}) d \vec{r}=\int_{S} \vec{C} \cdot d \vec{S}
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## Conversion to surface integral

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We have

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The actual scattering cross section is the square of this integral

$$
\frac{d \sigma}{d \Omega}=\left(\frac{r_{0} \rho}{Q_{z}}\right)^{2} \int_{S} \int_{S^{\prime}} e^{i Q_{z}\left(h(x, y)-h\left(x^{\prime}, y^{\prime}\right)\right)} e^{i Q_{x}\left(x-x^{\prime}\right)} e^{i Q_{y}\left(y-y^{\prime}\right)} d x d y d x^{\prime} d y^{\prime}
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## Scattering cross section

If we assume that $h(x, y)-h\left(x^{\prime}, y^{\prime}\right)$ depends only on the relative difference in position, $x-x^{\prime}$ and $y-y^{\prime}$ the four dimensional integral collapses to the product of two two dimensional integrals

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where $A_{0} / \sin \theta_{1}$ is just the illuminated surface area and the term in the angled brackets is an ensemble average over all possible choices of the origin within the illuminated area.
Finally, it is assumed that the statistics of the height variation are Gaussian and

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\left(\frac{d \sigma}{d \Omega}\right)=\left(\frac{r_{0} \rho}{Q_{z}}\right)^{2} \frac{A_{0}}{\sin \theta_{1}} \int e^{-Q_{z}^{2}\left\langle[h(0,0)-h(x, y)]^{2}\right\rangle / 2} e^{i Q_{x} x} e^{i Q_{y} y} d x d y
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R\left(Q_{z}\right)=\frac{I_{s c}}{I_{0}}=\left(\frac{Q_{c}^{2} / 8}{Q_{z}}\right)^{2}\left(\frac{1}{Q_{z} / 2}\right)^{2}=\left(\frac{Q_{c}}{2 Q_{z}}\right)^{4}
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where $h$ is a fractal parameter which defines the shape of the surface. jagged surface for $h \ll 1 \quad$ smoother surface for $h \rightarrow 1$ If the resolution in the $y$ direction is very broad (typical for a synchrotron), we can eliminate the $y$-integral and have

## Correlated surfaces

Assume that height fluctuations are isotropically correlated in the $x-y$ plane. Therefore, $g(x, y)=g(r)=g\left(\sqrt{x^{2}+y^{2}}\right)$.
In the limit that the correlations are unbounded as $r \rightarrow \infty, g(x, y)$ is given by

$$
g(x, y)=\mathcal{A} r^{2 h}
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where $h$ is a fractal parameter which defines the shape of the surface. jagged surface for $h \ll 1 \quad$ smoother surface for $h \rightarrow 1$

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$$
\left(\frac{d \sigma}{d \Omega}\right)=\left(\frac{r_{0} \rho}{Q_{z}}\right)^{2} \frac{A_{0}}{\sin \theta_{1}} \int e^{-\mathcal{A} Q_{z}^{2}|x|^{2 h} / 2} \cos \left(Q_{x} x\right) d x
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And the scattering exhibits both a specular peak, reduced by uncorrelated roughness, and diffuse scattering from the correlated portion of the surface

