PHYS 570 - Introduction to Synchrotron Radiation

Term: Fall 2016

Meetings: Monday & Wednesday 17:00-18:15

Location: 212 Stuart Building

Instructor: Carlo Segre

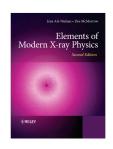
Office: 106A Life Sciences

Phone: 312.567.3498 email: segre@iit.edu

Book: Elements of Modern X-Ray Physics, 2nd ed.,

J. Als-Nielsen and D. McMorrow (Wiley, 2011)

Web Site: http://csrri.iit.edu/~segre/phys570/16F



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- Understand the function of various components of a synchrotron beamline

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- Be able to write a General User Proposal in the format used by the Advanced Photon Source

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- Visits to Advanced Photon Source (outside class, not required)
 - All students who plan to attend will need to request badges from APS
 - Go to the APS User Portal, https://www1.aps.anl.gov/Users-Information and register as a new user.
 - Use MRCAT (Sector 10) as location of experiment
 - Use Carlo Segre as local contact
 - State that your beamtime will be in the second week of October

33% – Homework assignments

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Grading scale

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A - 80% to 100%
B - 65% to 80%
C - 50% to 65%
E - 0% to 50%
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• X-rays and their interaction with matter

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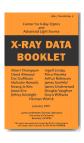
- X-rays and their interaction with matter
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- X-ray Oriented Programs: http://www.esrf.eu/Instrumentation/software/data-analysis/xop2.4



Today's outline - August 22, 2016

- History of x-ray sources
- X-ray interactions with matter

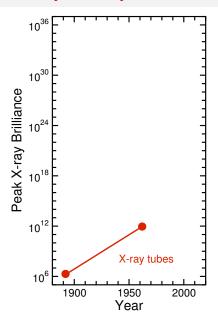
- History of x-ray sources
- X-ray interactions with matter
- Thomson scattering

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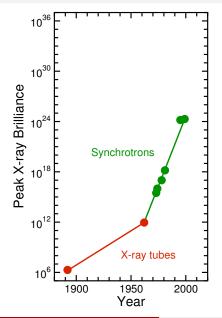
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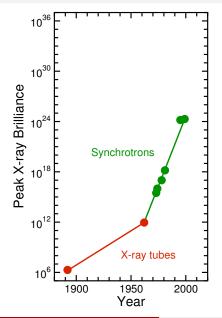
Reading Assignment: Chapter 1.1–1.6; 2.1–2.2



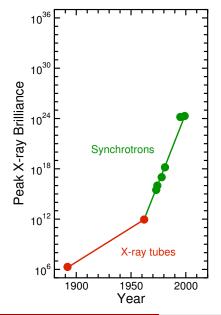
 1895 x-rays discovered by William Röntgen



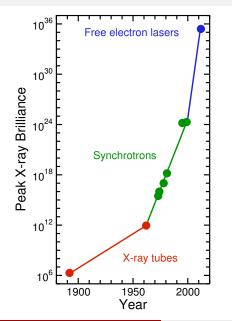
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$$\begin{array}{lll} \lambda & = & hc/\mathcal{E} \\ & = & (4.1357 \times 10^{-15} \, \mathrm{eV \cdot s})(2.9979 \times 10^8 \, \mathrm{m/s})/\mathcal{E} \\ & = & (4.1357 \times 10^{-18} \, \mathrm{keV \cdot s})(2.9979 \times 10^{18} \, \text{Å/s})/\mathcal{E} \\ & = & 12.398 \, \text{Å} \cdot \mathrm{keV}/\mathcal{E} \quad \text{to give units of Å} \end{array}$$

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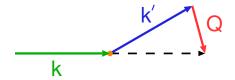
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We will only discuss the first three.

Most of the phenomena we will discuss can be treated classically as elastic scattering of electromagnetic waves (x-rays)

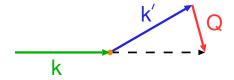
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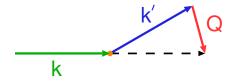
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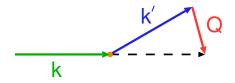
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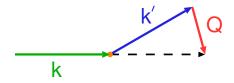
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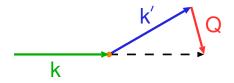
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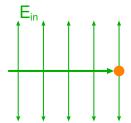
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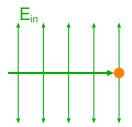
Start with the scattering from a single electron, then build up to more complexity

Assumptions:



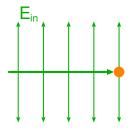
Assumptions:

incident x-ray plane wave



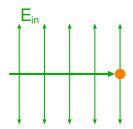
Assumptions:

incident x-ray plane wave electron is a point charge



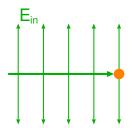
Assumptions:

incident x-ray plane wave electron is a point charge scattering is elastic



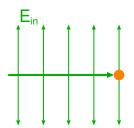
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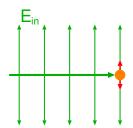
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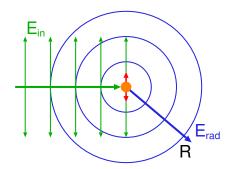


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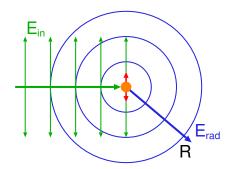
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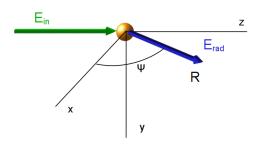


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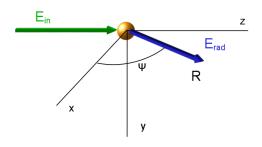
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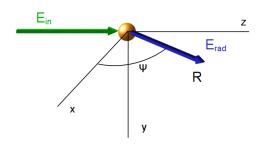
Using this, calculate the elastic scattering cross-section



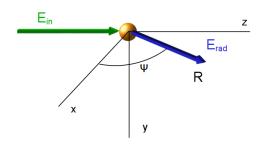
$$E_{rad}(R,t) = -rac{-e}{4\pi\epsilon_0c^2R} rac{a_{\scriptscriptstyle X}(t')}{\sin\Psi}$$



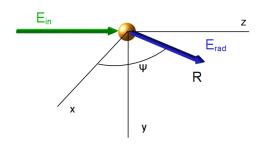
$$E_{rad}(R,t) = -\frac{-e}{4\pi\epsilon_0 c^2 R} \frac{a_x(t')}{a_x(t')} \sin \Psi$$
 where $t' = t - R/c$



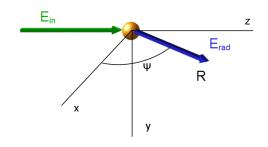
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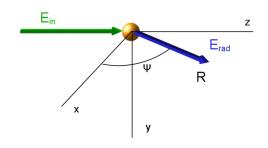
$$\begin{split} E_{rad}(R,t) &= -\frac{-e}{4\pi\epsilon_0 c^2 R} a_{x}(t') \sin \Psi & \text{where} \quad t' = t - R/c \\ a_{x}(t') &= -\frac{e}{m} E_{x0} e^{-i\omega t'} = -\frac{e}{m} E_{x0} e^{-i\omega t} e^{i\omega R/c} \end{split}$$



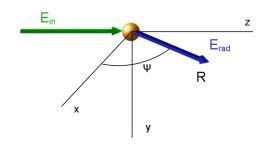
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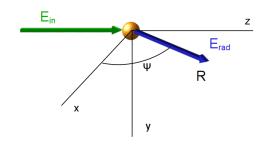
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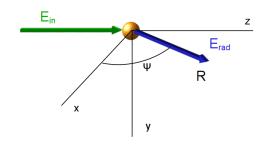
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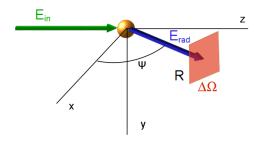
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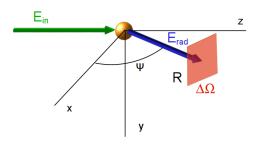


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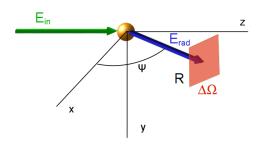


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$$r_0 = \frac{e^2}{4\pi\epsilon_0 mc^2} = 2.82 \times 10^{-5} \text{Å}$$

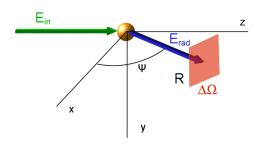




Detector of solid angle $\Delta\Omega$ at a distance R from electron

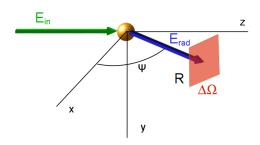


Detector of solid angle $\Delta\Omega$ at a distance R from electron Cross-section of incoming beam is A_0



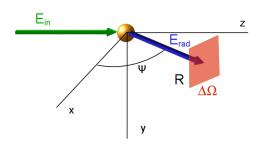
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$$\Phi_0 \equiv \frac{I_0}{A_0} = c \frac{|E_{in}|^2}{\hbar \omega}$$



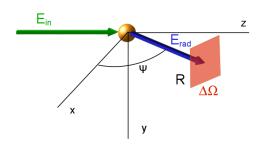
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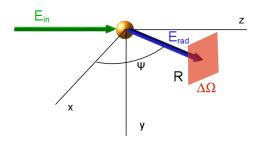
Detector of solid angle $\Delta\Omega$ at a distance R from electron Cross-section of incoming beam is A_0 Cross section of scattered beam (into detector) is $R^2\Delta\Omega$

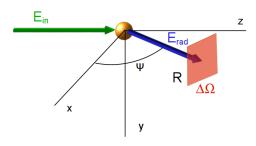
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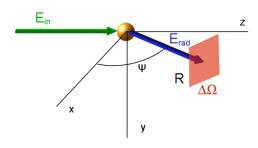
$$\begin{split} \Phi_0 &\equiv \frac{I_0}{A_0} = c \frac{\left|E_{in}\right|^2}{\hbar \omega} \\ I_{sc} &\propto c (R^2 \Delta \Omega) \frac{\left|E_{rad}\right|^2}{\hbar \omega} \\ \frac{I_{sc}}{I_0} &= \frac{\left|E_{rad}\right|^2}{\left|E_{in}\right|^2} R^2 \Delta \Omega \end{split}$$



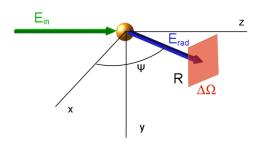


Differential cross-section is obtained by normalizing

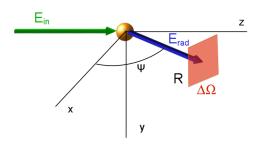
 $\frac{d\sigma}{d\Omega}$



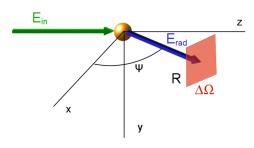
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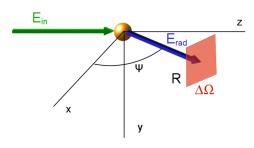


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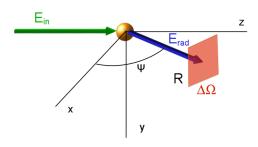


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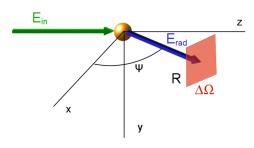
$$\frac{E_{rad}}{E_{in}} = -r_0 \frac{e^{ikR}}{R} |\hat{\epsilon} \cdot \hat{\epsilon}'|$$



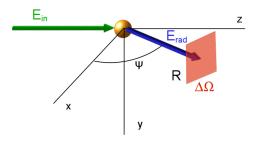
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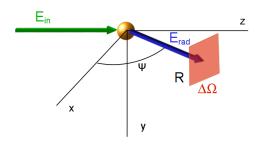
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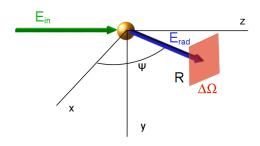


Integrate to obtain the total Thomson scattering cross-section from an electron.



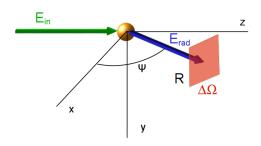
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= 0.665 × 10⁻²⁴ cm²
= 0.665 barn

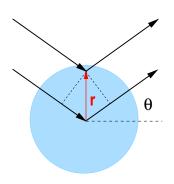


Integrate to obtain the total Thomson scattering cross-section from an electron. If displacement is in vertical direction, $\sin \Psi$ term is replaced by unity and if the source is unpolarized, it is a combination.

$$\begin{split} \sigma &= \frac{8\pi}{3} r_0^2 \\ &= 0.665 \times 10^{-24} \ cm^2 \\ &= 0.665 \ \textit{barn} \end{split} \qquad \begin{array}{l} \text{Polarization factor} &= \begin{cases} 1 \\ \sin^2 \Psi \\ \frac{1}{2} \left(1 + \sin^2 \Psi \right) \end{cases} \end{split}$$

Atomic scattering

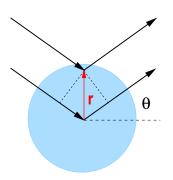
If we have a charge distribution instead of a single electron, the scattering is more complex



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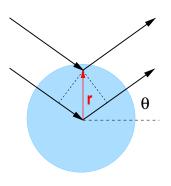
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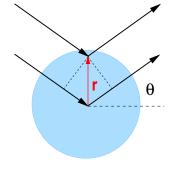
$$\Delta\phi(\mathbf{r}) = (\mathbf{k} - \mathbf{k}') \cdot \mathbf{r} = \mathbf{Q} \cdot \mathbf{r}$$

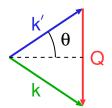


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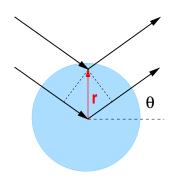


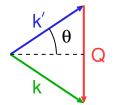


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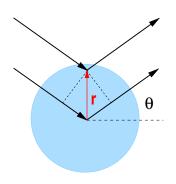


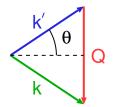
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The volume element at **r** contributes $-r_0\rho(\mathbf{r})d^3r$ with phase factor $e^{i\mathbf{Q}\cdot\mathbf{r}}$

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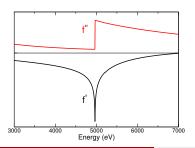
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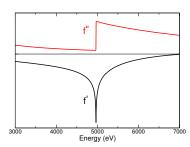
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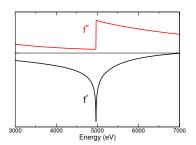


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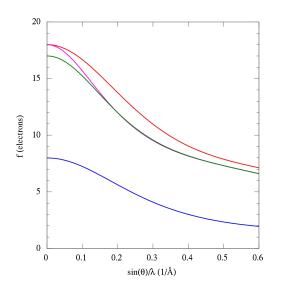
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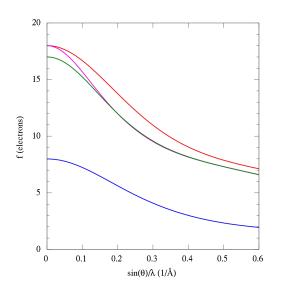


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$$f(\mathbf{Q},\hbar\omega) = f^0(\mathbf{Q}) + f'(\hbar\omega) + if''(\hbar\omega)$$

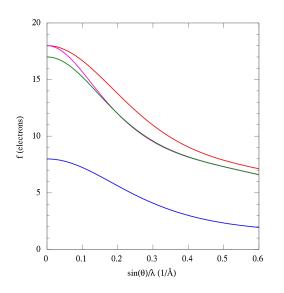


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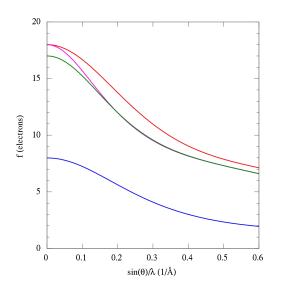
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The atomic form factor has an angular dependence

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lighter atoms have a broader form factor

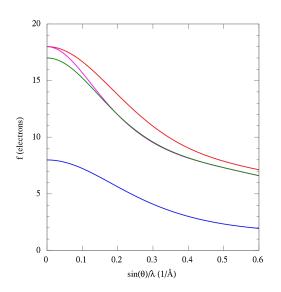


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Scattering from an atom is built up from component quantities:

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Thomson scattering from a single electron $-r_0 = -\frac{e^2}{4\pi\epsilon_0 mc^2}$

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polarization factor

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