## PHYS 570 - Introduction to Synchrotron Radiation

| Term: | Fall 2016 |
| :--- | :--- |
| Meetings: | Monday \& Wednesday 17:00-18:15 |
| Location: | 212 Stuart Building |
|  |  |
| Instructor: | Carlo Segre |
| Office: | 106A Life Sciences |
| Phone: | 312.567.3498 <br> email: <br> segre@iit.edu |



Book: Elements of Modern X-Ray Physics, $2^{\text {nd }}$ ed., J. Als-Nielsen and D. McMorrow (Wiley, 2011)

Web Site: http://csrri.iit.edu/~segre/phys570/16F

## Course objectives

- Understand the means of production of synchrotron x-ray radiation


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- Understand the function of various components of a synchrotron beamline


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- Be able to make an oral presentation of a synchrotron radiation research topic


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- Understand the physics behind a variety of experimental techniques
- Be able to make an oral presentation of a synchrotron radiation research topic
- Be able to write a General User Proposal in the format used by the Advanced Photon Source


## Course syllabus

- Focus on applications of synchrotron radiation


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- Make proposal and get approval before starting
- Visits to Advanced Photon Source (outside class, not required)
- All students who plan to attend will need to request badges from APS
- Go to the APS User Portal, https://www1.aps.anl.gov/Users-Information and register as a new user.
- Use MRCAT (Sector 10) as location of experiment
- Use Carlo Segre as local contact
- State that your beamtime will be in the second week of October


## Course grading

$33 \%$ - Homework assignments

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Grading scale

$$
\begin{aligned}
& \text { A - } 80 \% \text { to } 100 \% \\
& \text { B - } 65 \% \text { to } 80 \% \\
& \text { C - } 50 \% \text { to } 65 \% \\
& \text { E - } 0 \% \text { to } 50 \%
\end{aligned}
$$

## Topics to be covered (at a minimum)

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- Resonant scattering


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- Resonant scattering
- Imaging


## Resources for the course

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- X-ray Oriented Programs: http://www.esrf.eu/Instrumentation/software/data-analysis/xop2.4


## Today's outline - August 22, 2016

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- Atomic form factor


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Reading Assignment: Chapter 1.1-1.6; 2.1-2.2

## History of x-ray sources



- 1895 x-rays discovered by William Röntgen


## History of x-ray sources



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## History of x-ray sources



- 1895 x-rays discovered by William Röntgen
- $1^{\text {st }}$ generation synchrotrons initially used in parasitic mode (SSRL, CHESS)
- $2^{\text {nd }}$ generation were dedicated sources (NSLS, SRC, CAMD)
- $3^{r d}$ generation featured insertion devices (APS, ESRF, ALS)
- $4^{\text {th }}$ generation are free electron lasers (LCLS, XFEL)


## The classical x-ray

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where $\hat{\epsilon}$ is a unit vector in the direction of the electric field, $\mathbf{k}$ is the wavevector of the radiation along the propagation direction, and $\omega$ is the angular frequency of oscillation of the radiation.

If the energy, $\mathcal{E}$ is in keV , the relationship among these quantities is given by:

$$
\hbar \omega=h \nu=\mathcal{E}, \lambda \nu=c
$$

$$
\begin{aligned}
\lambda & =h c / \mathcal{E} \\
& =\left(4.1357 \times 10^{-15} \mathrm{eV} \cdot \mathrm{~s}\right)\left(2.9979 \times 10^{8} \mathrm{~m} / \mathrm{s}\right) / \mathcal{E} \\
& =\left(4.1357 \times 10^{-18} \mathrm{keV} \cdot \mathrm{~s}\right)\left(2.9979 \times 10^{18} \AA / \mathrm{s}\right) / \mathcal{E} \\
& =12.398 \AA \cdot \mathrm{keV} / \mathcal{E} \quad \text { to give units of } \AA
\end{aligned}
$$

## Interactions of x-rays with matter

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4 Pair production
We will only discuss the first three.

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or in terms of momentum transfer: $\hbar \mathbf{Q}=\hbar \mathbf{k}-\hbar \mathbf{k}^{\prime}$
Start with the scattering from a single electron, then build up to more complexity

## Thomson scattering

Assumptions:


## Thomson scattering

Assumptions:
incident x-ray plane wave


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incident x-ray plane wave electron is a point charge


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scattering is elastic
scattered intensity $\propto 1 / R^{2}$


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The observer at $R$ "sees" a scattered electric field $E_{r a d}(R, t)$ at a later time $t=t^{\prime}+R / c$

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The observer at $R$ "sees" a scattered electric field $E_{r a d}(R, t)$ at a later time $t=t^{\prime}+R / c$

Using this, calculate the elastic scattering cross-section

## Thomson scattering



$$
E_{r a d}(R, t)=-\frac{-e}{4 \pi \epsilon_{0} c^{2} R} a_{x}\left(t^{\prime}\right) \sin \Psi
$$

## Thomson scattering



$$
E_{r a d}(R, t)=-\frac{-e}{4 \pi \epsilon_{0} c^{2} R} a_{x}\left(t^{\prime}\right) \sin \psi \quad \text { where } \quad t^{\prime}=t-R / c
$$

## Thomson scattering



$$
\begin{aligned}
E_{r a d}(R, t) & =-\frac{-e}{4 \pi \epsilon_{0} c^{2} R} a_{x}\left(t^{\prime}\right) \sin \psi \quad \text { where } \quad t^{\prime}=t-R / c \\
a_{x}\left(t^{\prime}\right) & =-\frac{e}{m} E_{\times 0} e^{-i \omega t^{\prime}}
\end{aligned}
$$

## Thomson scattering



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\begin{aligned}
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a_{x}\left(t^{\prime}\right) & =-\frac{e}{m} E_{x 0} e^{-i \omega t^{\prime}}=-\frac{e}{m} E_{x 0} e^{-i \omega t} e^{i \omega R / c}
\end{aligned}
$$

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a_{x}\left(t^{\prime}\right) & =-\frac{e}{m} E_{x 0} e^{-i \omega t^{\prime}}=-\frac{e}{m} E_{x 0} e^{-i \omega t} e^{i \omega R / c} \\
a_{x}\left(t^{\prime}\right) & =-\frac{e}{m} E_{i n} e^{i \omega R / c}
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E_{r a d}(R, t) & =-\frac{-e}{4 \pi \epsilon_{0} c^{2} R} \frac{-e}{m} E_{i n} e^{i \omega R / c} \sin \psi \\
\frac{E_{r a d}(R, t)}{E_{i n}} & =-\frac{e^{2}}{4 \pi \epsilon_{0} m c^{2}} \frac{e^{i \omega R / c}}{R} \sin \psi
\end{aligned}
$$

## Thomson scattering



$$
\begin{aligned}
& E_{r a d}(R, t)=-\frac{-e}{4 \pi \epsilon_{0} c^{2} R} \frac{-e}{m} E_{i n} e^{i \omega R / c} \sin \psi \\
& \frac{E_{r a d}(R, t)}{E_{i n}}=-\frac{e^{2}}{4 \pi \epsilon_{0} m c^{2}} \frac{e^{i \omega R / c}}{R} \sin \psi \quad \text { but } k=\omega / c
\end{aligned}
$$

## Thomson scattering



$$
\frac{E_{r a d}(R, t)}{E_{i n}}=-\frac{e^{2}}{4 \pi \epsilon_{0} m c^{2}} \frac{e^{i k R}}{R} \sin \psi=-r_{0} \frac{e^{i k R}}{R} \sin \psi
$$

## Thomson scattering



$$
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\frac{E_{r a d}(R, t)}{E_{i n}} & =-\frac{e^{2}}{4 \pi \epsilon_{0} m c^{2}} \frac{e^{i k R}}{R} \sin \psi=-r_{0} \frac{e^{i k R}}{R} \sin \psi \\
r_{0} & =\frac{e^{2}}{4 \pi \epsilon_{0} m c^{2}}=2.82 \times 10^{-5} \AA
\end{aligned}
$$

## Scattering cross-section



## Scattering cross-section



Detector of solid angle $\Delta \Omega$ at a distance $R$ from electron

## Scattering cross-section



Detector of solid angle $\Delta \Omega$ at a distance $R$ from electron Cross-section of incoming beam is $A_{0}$

## Scattering cross-section



Detector of solid angle $\Delta \Omega$ at a distance $R$ from electron Cross-section of incoming beam is

$$
\Phi_{0} \equiv \frac{l_{0}}{A_{0}}=c \frac{\left|E_{i n}\right|^{2}}{\hbar \omega}
$$

$A_{0}$

## Scattering cross-section



Detector of solid angle $\Delta \Omega$ at a distance $R$ from electron Cross-section of incoming beam is $A_{0}$

$$
\begin{aligned}
& \Phi_{0} \equiv \frac{I_{0}}{A_{0}}=c \frac{\left|E_{i n}\right|^{2}}{\hbar \omega} \\
& I_{s c} \propto c\left(R^{2} \Delta \Omega\right) \frac{\left|E_{r a d}\right|^{2}}{\hbar \omega}
\end{aligned}
$$

## Scattering cross-section



Detector of solid angle $\Delta \Omega$ at a distance $R$ from electron Cross-section of incoming beam is $A_{0}$
Cross section of scattered beam

$$
\Phi_{0} \equiv \frac{I_{0}}{A_{0}}=c \frac{\left|E_{i n}\right|^{2}}{\hbar \omega}
$$

$$
I_{s c} \propto c\left(R^{2} \Delta \Omega\right) \frac{\left|E_{r a d}\right|^{2}}{\hbar \omega}
$$ (into detector) is $R^{2} \Delta \Omega$

## Scattering cross-section



Detector of solid angle $\Delta \Omega$ at a distance $R$ from electron Cross-section of incoming beam is $A_{0}$
Cross section of scattered beam (into detector) is $R^{2} \Delta \Omega$

$$
\begin{aligned}
& \Phi_{0} \equiv \frac{I_{0}}{A_{0}}=c \frac{\left|E_{i n}\right|^{2}}{\hbar \omega} \\
& I_{s c} \propto c\left(R^{2} \Delta \Omega\right) \frac{\left|E_{r a d}\right|^{2}}{\hbar \omega} \\
& \frac{I_{s c}}{I_{0}}=\frac{\left|E_{r a d}\right|^{2}}{\left|E_{i n}\right|^{2}} R^{2} \Delta \Omega
\end{aligned}
$$

## Scattering cross-section



Differential cross-section is obtained by normalizing

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$$
\frac{d \sigma}{d \Omega}
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Differential cross-section is obtained by normalizing

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\frac{d \sigma}{d \Omega}=\frac{I_{s c}}{\Phi_{0} \Delta \Omega}
$$

## Scattering cross-section



Differential cross-section is obtained by normalizing

$$
\frac{d \sigma}{d \Omega}=\frac{I_{s c}}{\Phi_{0} \Delta \Omega}=\frac{I_{s c}}{\left(I_{0} / A_{0}\right) \Delta \Omega}
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\frac{d \sigma}{d \Omega}=\frac{I_{s c}}{\Phi_{0} \Delta \Omega}=\frac{I_{s c}}{\left(I_{0} / A_{0}\right) \Delta \Omega}=\frac{\left|E_{r a d}\right|^{2}}{\left|E_{i n}\right|^{2}} R^{2}
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\frac{d \sigma}{d \Omega} & =\frac{I_{s c}}{\Phi_{0} \Delta \Omega}=\frac{I_{s c}}{\left(I_{0} / A_{0}\right) \Delta \Omega}=\frac{\left|E_{r a d}\right|^{2}}{\left|E_{i n}\right|^{2}} R^{2} \\
\frac{E_{r a d}}{E_{i n}} & =-r_{0} \frac{e^{i k R}}{R}\left|\hat{\epsilon} \cdot \hat{\epsilon}^{\prime}\right|
\end{aligned}
$$

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\end{aligned}
$$

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\begin{aligned}
\frac{d \sigma}{d \Omega} & =\frac{I_{s c}}{\Phi_{0} \Delta \Omega}=\frac{I_{s c}}{\left(I_{0} / A_{0}\right) \Delta \Omega}=\frac{\left|E_{r a d}\right|^{2}}{\left|E_{i n}\right|^{2}} R^{2}=r_{0}^{2} \sin ^{2} \psi \\
\frac{E_{r a d}}{E_{i n}} & =-r_{0} \frac{e^{i k R}}{R}\left|\hat{\epsilon} \cdot \hat{\epsilon}^{\prime}\right|=-r_{0} \frac{e^{i k R}}{R}\left|\cos \left(\frac{\pi}{2}-\Psi\right)\right|=-r_{0} \frac{e^{i k R}}{R} \sin \psi
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## Total cross-section



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|\mathbf{Q}|=2|\mathbf{k}| \sin \theta=\frac{4 \pi}{\lambda} \sin \theta
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f(\mathbf{Q}, \hbar \omega)=f^{0}(\mathbf{Q})+f^{\prime}(\hbar \omega)+i f^{\prime \prime}(\hbar \omega)
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