## Today's Outline - April 23, 2015

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- Imaging


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- Imaging
- Radiography and Tomography


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Final Exam, Tuesday, May 5, 2015, Life Sciences 240
4 sessions: 09:00-11:00; 11:00-13:00; 14:00-16:00; 16:00-18:00

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Send me your presentation in Powerpoint or PDF format before before your session

## Problem 5.9

Consider a simple model of the surface roughness of a crystal in which all of the lattice sites of the $z=0$ layer are fully occupied by atoms, but the next layer out $(z=-1)$ has a site occupancy of $\eta$, with $\eta \leq 1$, the $z=-2$ layer an occupancy of $\eta^{2}$, etc. Show that midway between the Bragg points, the so-called anti-Bragg points, the intensity of the crystal truncation rods is given by

$$
I^{C T R}=\frac{(1-\eta)^{2}}{4(1+\eta)^{2}}
$$

What effect does a small, but finite, value of the roughness parameter $\eta$ have on I ${ }^{C T R}$ ?

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The general expression for the scattering factor of the crystal truncation rod in the $a_{3}$ direction of the crystal, excluding absorption effects is

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F^{C T R}=A(\vec{Q}) \sum_{j=0}^{\infty} e^{i Q_{z} a_{3} j}
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\begin{aligned}
F^{C T R} & =A(\vec{Q}) \sum_{j=0}^{\infty} e^{i Q_{z} a_{3} j} \\
& =A(\vec{Q})\left[\sum_{j=0}^{\infty} e^{i Q_{z} a_{3} j}+\sum_{j=0}^{\infty} \eta^{j} e^{-i Q_{z} a_{3} j}\right]
\end{aligned}
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The rougness model adds terms with $j<0$ of the form $\eta^{j} e^{-i Q_{z} a_{3} j}$ so a second sum is added
evaluating these geometric sums in the usual way

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& =A(\vec{Q})\left[\frac{1}{1-e^{i 2 \pi l}}+\frac{\eta e^{-i 2 \pi l}}{1-\eta e^{-i 2 \pi l}}\right]
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at the anti-Bragg points, $I=0.5$, and the exponentials all become -1

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& =A(\vec{Q})\left[\frac{1}{2}-\frac{\eta}{1+\eta}\right]
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& =A(\vec{Q})\left[\frac{1}{2}-\frac{\eta}{1+\eta}\right] \\
& =A(\vec{Q}) \frac{(1+\eta)-2 \eta}{2(1+\eta)}
\end{aligned}
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& =A(\vec{Q})\left[\frac{1}{1-e^{i 2 \pi l}}+\frac{\eta e^{-i 2 \pi I}}{1-\eta e^{-i 2 \pi l}}\right] \\
& =A(\vec{Q})\left[\frac{1}{2}-\frac{\eta}{1+\eta}\right] \\
& =A(\vec{Q}) \frac{(1+\eta)-2 \eta}{2(1+\eta)}=A(\vec{Q}) \frac{(1-\eta)}{2(1+\eta)}
\end{aligned}
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## Problem 5.9

The scattering factor thus has a dependence on the coverage factor, $\eta$, of

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$$
F^{C T R}=A(\vec{Q}) \frac{(1-\eta)}{2(1+\eta)}
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The scattering factor thus has a dependence on the coverage factor, $\eta$, of and the observed intensity

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F^{C T R}=A(\vec{Q}) \frac{(1-\eta)}{2(1+\eta)}
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The scattering factor thus has a dependence on the coverage factor, $\eta$, of
and the observed intensity has a dependence of
if the coverage is very low, this can be approximated using the first term of a binomial expansion

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$$
\begin{aligned}
F^{C T R} & =A(\vec{Q}) \frac{(1-\eta)}{2(1+\eta)} \\
I^{C T R} & =|A|^{2} \frac{(1-\eta)^{2}}{4(1+\eta)^{2}} \\
& \approx \frac{1}{4}|A|^{2}\left(\frac{1-2 \eta}{1+2 \eta}\right)
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\end{aligned}
$$ mial expansion

Thus any small value of the roughness parameter will dampen the intensity of the peaks

## Phase difference in scattering



## Phase difference in scattering



All imaging can be broken into a three step process

## Phase difference in scattering



All imaging can be broken into a three step process
(1) x-ray interaction with sample

## Phase difference in scattering



All imaging can be broken into a three step process
(1) $x$-ray interaction with sample
(2) scattered x-ray propagation

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The spherical waves scattered off the two points will travel different distances

## Phase difference in scattering



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(1) x-ray interaction with sample
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The spherical waves scattered off the two points will travel different distances

In the far field, the phase difference is $\vec{Q} \cdot \vec{r}$ with $\vec{Q}=\vec{k}-\overrightarrow{k^{\prime}}$

## Franuhofer, Fresnel, and contact regimes



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In the far field (Fraunhofer) regime, the phase difference is $\approx \overrightarrow{k^{\prime}} \cdot \vec{r}$

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In the far field (Fraunhofer) regime, the phase difference is $\approx \overrightarrow{k^{\prime}} \cdot \vec{r}$ The error in difference computed with the far field approximation is

$$
\Delta=R-R \cos \psi
$$

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In the far field (Fraunhofer) regime, the phase difference is $\approx \overrightarrow{k^{\prime}} \cdot \vec{r}$
The error in difference computed with the far field approximation is

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\begin{aligned}
\Delta & =R-R \cos \psi \\
& \approx R\left(1-\left(1-\psi^{2} / 2\right)\right)
\end{aligned}
$$

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In the far field (Fraunhofer) regime, the phase difference is $\approx \overrightarrow{k^{\prime}} \cdot \vec{r}$ The error in difference computed with the far field approximation is $\quad R \gg a^{2} / \lambda \quad$ Fraunhofer

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## Franuhofer, Fresnel, and contact regimes



In the far field (Fraunhofer) regime, the phase difference is $\approx \overrightarrow{k^{\prime}} \cdot \vec{r}$ The error in difference computed with the far field approximation is

$$
\begin{array}{rlrr}
\Delta & =R-R \cos \psi & & R \approx a^{2} / \lambda \\
& \approx R\left(1-\left(1-\psi^{2} / 2\right)\right) & & R \ll a^{2} / \lambda \\
& & \text { Fontact } \\
& =a^{2} /(2 R) & &
\end{array}
$$

$$
R \gg a^{2} / \lambda \quad \text { Fraunhofer }
$$

## Contact to far-field imaging



## Fourier slice theorem

$$
I=I_{0} e^{-\int \mu(x, y) d y^{\prime}}
$$



## Fourier slice theorem

$$
\begin{aligned}
I & =I_{0} e^{-\int \mu(x, y) d y^{\prime}} \\
\log _{10}\left(\frac{I_{0}}{I}\right) & =\int \mu(x, y) d y^{\prime}
\end{aligned}
$$



## Fourier slice theorem

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p(x) & =\int f(x, y) d y
\end{aligned}
$$



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\log _{10}\left(\frac{I_{0}}{I}\right) & =\int \mu(x, y) d y^{\prime} \\
p(x) & =\int f(x, y) d y \\
P\left(q_{x}\right) & =\int p(x) e^{i q_{x} x} d x
\end{aligned}
$$



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\end{aligned}
$$



$$
F\left(q_{x}, q_{y}\right)=\iint f(x, y) e^{i q_{x} x+q_{y} y} d x d y
$$

## Fourier slice theorem

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$$

$$
F\left(q_{x}, q_{y}=0\right)=\int\left[\int f(x, y) d y\right] e^{i q_{x} x} d x
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$$
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P\left(q_{x}\right) & =\int p(x) e^{i q_{x} x} d x
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$$



$$
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F\left(q_{x}, q_{y}=0\right)=\int\left[\int f(x, y) d y\right] e^{i q_{x} x} d x=\int p(x) e^{i q_{x} x} d x=P\left(q_{x}\right)
$$

## Fourier transform reconstruction



## Sinograms


(c) Model $f(x, y)$


(d) Sinogram


(e) Reconstructed $f(x, y)$


## Medical tomography



## Microscopy



## Microscopy

## Zone Plate

 Lens

## Sample

Translations $\downarrow$

## Microscopy



## Microscopy



## Angular deviation from refraction

When x-rays cross an interface that is not normal to their direction, there is refraction

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The angle of refraction $\alpha$ can be calculated

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$$
\lambda_{n}=\frac{\lambda}{n}
$$

## Angular deviation from refraction

When x-rays cross an interface that is not normal to their direction, there is refraction

The angle of refraction $\alpha$ can be calculated

$$
\lambda_{n}=\frac{\lambda}{n}=\frac{\lambda}{1-\delta}
$$

## Angular deviation from refraction

When x-rays cross an interface that is not normal to their direction, there is refraction

The angle of refraction $\alpha$ can be calculated

$$
\begin{aligned}
\lambda_{n} & =\frac{\lambda}{n}=\frac{\lambda}{1-\delta} \\
& \approx \lambda(1+\delta)
\end{aligned}
$$

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\begin{aligned}
\lambda_{n} & =\frac{\lambda}{n}=\frac{\lambda}{1-\delta} \\
& \approx \lambda(1+\delta) \\
\alpha & =\frac{\lambda(1+\delta)-\lambda}{\Delta x}
\end{aligned}
$$

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& =\delta \frac{\lambda}{\Delta x}
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& \approx \lambda(1+\delta) \\
\alpha & =\frac{\lambda(1+\delta)-\lambda}{\Delta x} \\
& =\delta \frac{\lambda}{\Delta x} \approx \delta \tan \omega
\end{aligned}
$$

## Angular deviation from graded density

In a similar way, there is an angular deviation when the material density varies normal to the propagation direction

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## Angular deviation from graded density



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The angle of refraction $\alpha$ can be calculated

$$
\alpha=\frac{\lambda(1+\delta(x+\Delta x))-\lambda(1+\delta(x))}{\Delta x}
$$

## Angular deviation from graded density



In a similar way, there is an angular deviation when the material density varies normal to the propagation direction

The angle of refraction $\alpha$ can be calculated

$$
\begin{array}{r}
\alpha=\frac{\lambda(1+\delta(x+\Delta x))-\lambda(1+\delta(x))}{\Delta x} \\
\delta(x+\Delta x) \approx \delta(x)+\Delta x \frac{\partial \delta(x)}{\partial x}
\end{array}
$$

## Angular deviation from graded density



In a similar way, there is an angular deviation when the material density varies normal to the propagation direction

The angle of refraction $\alpha$ can be calculated

$$
\begin{gathered}
\alpha=\frac{\lambda(1+\delta(x+\Delta x))-\lambda(1+\delta(x))}{\Delta x}=\frac{\lambda \Delta x \frac{\partial \delta(x)}{\partial x}}{\Delta x} \\
\delta(x+\Delta x) \approx \delta(x)+\Delta x \frac{\partial \delta(x)}{\partial x}
\end{gathered}
$$

## Angular deviation from graded density



In a similar way, there is an angular deviation when the material density varies normal to the propagation direction

The angle of refraction $\alpha$ can be calculated

$$
\begin{aligned}
\alpha= & \frac{\lambda(1+\delta(x+\Delta x))-\lambda(1+\delta(x))}{\Delta x}=\frac{\lambda \Delta x \frac{\partial \delta(x)}{\partial x}}{\Delta x} \\
& \delta(x+\Delta x) \approx \delta(x)+\Delta x \frac{\partial \delta(x)}{\partial x} \\
\alpha_{\text {gradient }}= & \lambda \frac{\partial \delta(x)}{\partial x}
\end{aligned}
$$

## Angular deviation from graded density



In a similar way, there is an angular deviation when the material density varies normal to the propagation direction

The angle of refraction $\alpha$ can be calculated

$$
\begin{gathered}
\alpha=\frac{\lambda(1+\delta(x+\Delta x))-\lambda(1+\delta(x))}{\Delta x}=\frac{\lambda \Delta x \frac{\partial \delta(x)}{\partial x}}{\Delta x} \\
\delta(x+\Delta x) \approx \delta(x)+\Delta x \frac{\partial \delta(x)}{\partial x} \\
\alpha_{\text {gradient }}=\lambda \frac{\partial \delta(x)}{\partial x} \quad \text { compare to } \quad \alpha_{\text {refrac }}=\lambda \frac{\delta}{\Delta x}
\end{gathered}
$$

## Phase shift from angular deviation

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\begin{aligned}
\hat{n} & =\frac{\vec{k}^{\prime}}{k^{\prime}}=\frac{\lambda}{2 \pi} \nabla \phi(\vec{r}) \\
\alpha_{x} & =\frac{\lambda}{2 \pi} \frac{\partial \phi(x, y)}{\partial x} \\
\alpha_{y} & =\frac{\lambda}{2 \pi} \frac{\partial \phi(x, y)}{\partial y}
\end{aligned}
$$

Thus the angular deviation, in each of the $x$ and $y$ directions in the plane perpendicular to the original propagation direction becomes
By measuring the angular deviation as a function of position in a sample, one can reconstruct the phase shift $\phi(x, y)$ due to the sample by integration.

## Phase contrast experiment



## Phase contrast experiment



## Imaging a silicon trough



## Imaging blood cells



