• Lattice & Basis Functions

- Lattice & Basis Functions
- Reciprocal Lattice for FCC

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Reading assignment: Chapter 5.2

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Reading assignment: Chapter 5.2

Homework Assignment #04: Chapter 4: 2, 4, 6, 7, 10 due Tuesday, March 10, 2015

$$\vec{R}_n = n_1 \vec{a}_1 + n_2 \vec{a}_2$$



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primitive



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non-primitive



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#### primitive

non-primitive

#### non-conventional

#### More about lattice vectors



sometimes conventional axes...

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#### sometimes conventional axes...

... are not primitive

## Miller indices



planes designated (hk), intercept the unit cell axes at

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### Miller indices



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# Miller indices

(20) a<sub>2</sub> planes designated (hk), intercept the unit cell axes at

 $\frac{a_1}{h}, \quad \frac{a_2}{k}$ 

for a lattice with orthogonal unit vectors

$$\frac{1}{d_{hk}^2} = \frac{h^2}{a_1^2} + \frac{k^2}{a_2^2}$$

# Reciprocal lattice



# **Reciprocal lattice**



$$ec{a}_1^* = rac{2\pi}{V_c}ec{a}_2 imes ec{a}_3 \qquad ec{a}_2^* = rac{2\pi}{V_c}ec{a}_3 imes ec{a}_1 \qquad ec{a}_3^* = rac{2\pi}{V_c}ec{a}_1 imes ec{a}_2$$

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$$F^{crystal}(\vec{Q}) = \sum_{l}^{N} f_{l}(\vec{Q}) e^{i\vec{Q}\cdot\vec{r}_{l}}$$

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Since  $F^{crystal}(\vec{Q})$  is simply the Fourier Transform of the crystal function,  $C(x) = \mathcal{L}(x) \star \mathcal{B}(x)$ , it must be the product of the Fourier Transforms of  $\mathcal{L}(x)$  and  $\mathcal{B}(x)$ .

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$$\vec{G}_{hkl} = h\vec{a}_1^* + k\vec{a}_2^* + l\vec{a}_3^*$$
  
$$\vec{G}_{hkl} \cdot \vec{R}_n = (n_1\vec{a}_1 + n_2\vec{a}_2 + n_3\vec{a}_3) \cdot (h\vec{a}_1^* + k\vec{a}_2^* + l\vec{a}_3^*)$$

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$$\therefore \quad \vec{Q} = \vec{G}_{hkl}$$



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The primitive lattice vectors of the face-centered cubic lattice are

$$\vec{a}_1 = \frac{a}{2}(\hat{y} + \hat{z}), \quad \vec{a}_2 = \frac{a}{2}(\hat{z} + \hat{x}), \quad \vec{a}_3 = \frac{a}{2}(\hat{x} + \hat{y})$$

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The volume of the unit cell is  

$$v_c = \vec{a}_1 \cdot \vec{a}_2 \times \vec{a}_3 = \vec{a}_1 \cdot \frac{a^2}{4} (\hat{y} + \hat{z} - \hat{x})$$

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&= \frac{4\pi}{a} \left( \frac{\hat{y}}{2} + \frac{\hat{z}}{2} - \frac{\hat{x}}{2} \right) \\
\vec{a}_2^* &= \frac{4\pi}{a} \left( \frac{\hat{z}}{2} + \frac{\hat{x}}{2} - \frac{\hat{y}}{2} \right) \\
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PHYS 570 - Spring 2015

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which is a body-centered cubic lattice

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PHYS 570 - Spring 2015

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$$|S_N(Q)| = \frac{\sin(N[h+\xi]a^*a/2)}{\sin([h+\xi]a^*a/2)} = \frac{\sin(N[h+\xi]\pi)}{\sin([h+\xi]\pi)}$$
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$$sin(N\pi[h+\xi]) = sin(N\pi h) cos(N\pi\xi) + cos(N\pi h) sin(N\pi\xi)$$
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with half-width  $|S_N(\xi)| 
ightarrow 0, \quad N\pi\xi = \pi, \quad \xi_{1/2} \approx rac{1}{2N}$ 

the peak area can be obtained by integration
$$\int_{-1/2N}^{+1/2N} |S_N(\xi)| \, d\xi = \int_{-1/2N}^{+1/2N} \frac{\sin(N\pi\xi)}{\sin(\pi\xi)} d\xi$$

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 $|S_N(\xi)| \to \delta(\xi)$ 

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$$|S_N(Q)| o a^* \sum_{G_h} \delta(Q - G_h) = \sum_{n=0}^{N-1} e^{iQ_{na}}$$

That is, the lattice sum (scattering factor) is simply proportional to the reciprocal space lattice

C. Segre (IIT)

PHYS 570 - Spring 2015

the 1D modulus squared

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$$|S_N(Q)|^2 
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in 2D, with  $N_1 \times N_2 = N$  unit cells

the 1D modulus squared

$$|S_N(Q)|^2 o Na^* \sum_{G_h} \delta(Q - G_h)$$

$$\left|S_{N}(\vec{Q})\right|^{2} \rightarrow (N_{1}a_{1}^{*})(N_{2}a_{2}^{*})\sum_{\vec{G}}\delta(\vec{Q}-\vec{G}_{h})$$

in 2D, with  $N_1 \times N_2 = N$  unit cells

the 1D modulus squared

in 2D, cells

$$|S_N(Q)|^2 o Na^* \sum_{G_h} \delta(Q - G_h)$$

with 
$$N_1 \times N_2 = N$$
 unit  
$$\begin{vmatrix} S_N(\vec{Q}) \end{vmatrix}^2 \rightarrow (N_1 a_1^*) (N_2 a_2^*) \sum_{\vec{G}} \delta(\vec{Q} - \vec{G}_h) \\ = NA^* \sum_{\vec{G}} \delta(\vec{Q} - \vec{G}_h)$$

the 1D modulus squared

$$|S_N(Q)|^2 \to Na^* \sum_{G_h} \delta(Q - G_h)$$

in 2D, with 
$$N_1 \times N_2 = N$$
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cells  $|S_N(\vec{Q})|^2 \rightarrow (N_1 a_1^*)(N_2 a_2^*) \sum_{\vec{G}} \delta(\vec{Q} - \vec{G}_h)$   
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and similarly in 3D

the 1D modulus squared

$$|S_N(Q)|^2 
ightarrow Na^* \sum_{G_h} \delta(Q - G_h)$$

in 2D, with 
$$N_1 \times N_2 = N$$
 unit cells  
$$\left| S_N(\vec{Q}) \right|^2 \to (N_1 a_1^*) (N_2 a_2^*) \sum_{\vec{G}} \delta(\vec{Q} - \vec{G}_h)$$
$$= NA^* \sum_{\vec{G}} \delta(\vec{Q} - \vec{G}_h)$$

and similarly in 3D

$$\left|S_N(\vec{Q})\right|^2 o NV_c^* \sum_{\vec{G}} \delta(\vec{Q} - \vec{G}_h)$$

$$\int_{-\infty}^{\infty} \mathcal{L}(x) e^{iQx} dx = \int_{-\infty}^{\infty} \sum_{n} \delta(x - na) e^{iQx} dx$$

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$$= \sum_{n} e^{iQna} = a^* \sum_{h} \delta(Q - ha^*) = a^* \sum_{h} \delta(Q - G_h)$$

Consider the Fourier transform of the lattice function,  $\mathcal{L}(x)$ , (in 1-D for simplicity)

$$\int_{-\infty}^{\infty} \mathcal{L}(x)e^{iQx}dx = \int_{-\infty}^{\infty} \sum_{n} \delta(x - na)e^{iQx}dx = \sum_{n} \int_{-\infty}^{\infty} \delta(x - na)e^{iQx}dx$$
$$= \sum_{n} e^{iQna} = a^* \sum_{h} \delta(Q - ha^*) = a^* \sum_{h} \delta(Q - G_h)$$

in general

$$\int_{-\infty}^{\infty} \mathcal{L}(\vec{r}) e^{i\vec{Q}\cdot\vec{r}} dV = V_c^* \sum_{h,k,l} \delta(\vec{Q} - \vec{G}_{hkl})$$





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# Laue condition



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$$\vec{Q} = \vec{G_{hk}}$$



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$$Q = 2k \sin \theta = \frac{2\pi}{d}$$
$$2d \sin \theta = \frac{2\pi}{k} = \lambda$$


The Laue condition states that the scattering vector must be equal to a reciprocal lattice vector

$$\vec{Q} = \vec{G_{hk}}$$
$$Q = 2k \sin \theta = \frac{2\pi}{d}$$
$$2d \sin \theta = \frac{2\pi}{k} = \lambda$$



Thus the Bragg and Laue conditions are equivalent





Must show that for each point in reciprocal space, there exists a set of planes in the real space lattice such that:



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 $\vec{G}_{hkl}$  is perpendicular to the planes with Miller indices (hkl) and

$$|\vec{G}_{hkl}| = \frac{2\pi}{d_{hkl}}$$





The plane with Miller indices (hkl) intersects the three basis vectors of the lattice at  $a_1/h$ ,  $a_2/k$ , and  $a_3/l$ 



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$$ec{v_1} = rac{ec{a_3}}{ec{l}} + rac{ec{a_1}}{h}, \quad ec{v_1} = rac{ec{a_1}}{h} + rac{ec{a_2}}{k}$$



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 $\vec{G}_{hkl} \cdot \vec{v} = \left(h\vec{a}_1^* + k\vec{a}_2^* + l\vec{a}_3^*\right) \cdot \left(\left(\epsilon_2 - \epsilon_1\right)\frac{\vec{a}_1}{h} - \epsilon_2\frac{\vec{a}_2}{k} + \epsilon_1\frac{\vec{a}_3}{l}\right)$ 



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$$= 2\pi(\epsilon_2 - \epsilon_1 - \epsilon_2 + \epsilon_1) = 0$$

Thus  $\vec{G}_{hkl}$  is indeed normal to the plane with Miller indices (hkl)

C. Segre (IIT)





The spacing between planes (hkl) is simply given by the distance from the origin to the plane along a normal vector



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This can be computed as the projection of any vector which connects the origin to the plane onto the unit vector in the  $\vec{G}_{hkl}$  direction. In this case, we choose,  $\vec{a}_1/h$ 



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$$= f(\vec{G}) \left(1 + e^{i\pi(h+k+l)}\right)$$
$$= f(\vec{G}) \times \begin{cases} 2 & h+k+l = 2n \\ 0 & \text{otherwise} \end{cases}$$



In the face centered cubic structure, there are 4 atoms in the conventional, cubic unit cell. These are located at



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=  $f(\vec{G}) \left( 1 + e^{i\pi(h+k)} + e^{i\pi(k+l)} + e^{i\pi(h+l)} \right)$   
=  $f(\vec{G}) \times \begin{cases} 4 & h+k, k+l, h+l = 2n \\ 0 & \text{otherwise} \end{cases}$ 



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This is non-zero when h,k,l all even and h + k + l = 4n or h,k,l all odd





 $\leftarrow \mathsf{bcc}$ 



$$\leftarrow \mathsf{bcc}$$





$$\leftarrow bcc$$
sc  $\rightarrow$ 



 $\leftarrow \mathsf{diamond}$ 





