

Today's outline - March 31, 2022





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- Properties of density operators



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- Geometry of mixed states



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Reading assignment: 10.2 – 10.3



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Homework Assignment #06:

See Blackboard

Due Tuesday, April 05, 2022

Quantum circuit simulator <https://algassert.com/quirk>



More properties of density operators

Any density operator, ρ_x^A , is Hermitian and $\text{Tr}(\rho_x^A) = \sum_j \overline{x_{ij}} x_{ij} \equiv 1$ since $|x\rangle$ is a unit vector



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$$\sum_{i=0}^{M-1} \lambda_i \equiv 1, \quad \rho_x^A = \sum_{i=0}^{M-1} \lambda_i |v_i\rangle \langle v_i|$$



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Let B be a quantum system with a vector space of dimension $2^n > M$ and let $\{|\phi_0\rangle, \dots, |\phi_{M-1}\rangle\}$ be the first M elements of an orthonormal basis for B , define



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The density operator of a pure state is a projection operator, such that $\rho_x^X \rho_x^X = \rho_x^X$



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$$\det(\rho) = \begin{vmatrix} \frac{1+z}{2} & \frac{x-iy}{2} \\ \frac{x+iy}{2} & \frac{1-z}{2} \end{vmatrix}$$



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Geometry of mixed states

It is possible to use the Bloch sphere to visualize single-qubit mixed states which are linear combinations of pure states with non-negative coefficients that sum to 1

A density operator for a single qubit must be Hermitian and self-adjoint with trace 1 and the general form

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A maximally uncertain state for an n -qubit system has all the diagonal elements equal to 2^{-n} so $S(\rho) = n$



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The von Neumann entropy for a single qubit system is just a function of the distance of the state from the center of the Bloch sphere

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Bipartite entanglement

It is useful to find a good measure of entanglement for bipartite systems such as $X = A \otimes B$

The 2-qubit system is the simplest bipartite system with a maximally entangled state

For each of the two qubits, the density matrix ρ_{ME} has maximal von Neumann entropy

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It makes sense to use the von Neumann entropy of the partial trace as a measure of the entanglement if it can be assumed that the partial trace is the same for each of the two subsystems, A and B



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$$\text{Tr}_B(\rho) = \sum_{k=0}^{K-1} \langle \psi_k^B | \rho | \psi_k^B \rangle = \sum_{k=0}^{K-1} \lambda_k^2 |\psi_k^A\rangle \langle \psi_k^A|$$

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Clearly, $S(\text{Tr}_A(\rho)) = S(\text{Tr}_B(\rho))$ which means that the von Neumann entropy of the partial trace of a bipartite system is consistent when measured on either subsystem

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The amount of entanglement between the two parts of a pure state $|\psi\rangle \in X = A \otimes B$ with density operator $\rho = |\psi\rangle\langle\psi|$ is defined to be $S(\text{Tr}_A(\rho))$ or $S(\text{Tr}_B(\rho))$



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$$(\rho_y^A)_{01} = \frac{1}{2}(\langle 00|01\rangle + \langle 00|10\rangle)(\langle 01|10\rangle + \langle 10|10\rangle) + \frac{1}{2}(\langle 01|01\rangle + \langle 01|10\rangle)(\langle 01|11\rangle + \langle 10|11\rangle)$$



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$$(\rho_y^A)_{10} = \frac{1}{2}(\langle 10|01\rangle + \langle 10|10\rangle)(\langle 01|00\rangle + \langle 10|00\rangle) + \frac{1}{2}(\langle 11|01\rangle + \langle 11|10\rangle)(\langle 01|01\rangle + \langle 10|01\rangle)$$



Example 10.2.1 (cont.)

$$(\rho_y^A)_{01} = \frac{1}{2}(\cancel{\langle 00|01\rangle} + \cancel{\langle 00|10\rangle})(\cancel{\langle 01|10\rangle} + \langle 10|10\rangle) + \frac{1}{2}(\langle 01|01\rangle + \cancel{\langle 01|10\rangle})(\cancel{\langle 01|11\rangle} + \cancel{\langle 10|11\rangle}) \\ = 0$$

$$(\rho_y^A)_{10} = \frac{1}{2}(\cancel{\langle 10|01\rangle} + \langle 10|10\rangle)(\cancel{\langle 01|00\rangle} + \cancel{\langle 10|00\rangle}) + \frac{1}{2}(\cancel{\langle 11|01\rangle} + \cancel{\langle 11|10\rangle})(\cancel{\langle 01|01\rangle} + \cancel{\langle 10|01\rangle}) \\ = 0$$



Example 10.2.1 (cont.)

$$(\rho_y^A)_{01} = \frac{1}{2}(\cancel{\langle 00|01\rangle} + \cancel{\langle 00|10\rangle})(\cancel{\langle 01|10\rangle} + \langle 10|10\rangle) + \frac{1}{2}(\langle 01|01\rangle + \cancel{\langle 01|10\rangle})(\cancel{\langle 01|11\rangle} + \cancel{\langle 10|11\rangle}) \\ = 0$$

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$$(\rho_y^A)_{11} = \frac{1}{2}(\langle 10|01\rangle + \langle 10|10\rangle)(\langle 01|10\rangle + \langle 10|10\rangle) + \frac{1}{2}(\langle 11|01\rangle + \langle 11|10\rangle)(\langle 01|11\rangle + \langle 10|11\rangle)$$



Example 10.2.1 (cont.)

$$(\rho_y^A)_{01} = \frac{1}{2}(\cancel{\langle 00|01\rangle} + \cancel{\langle 00|10\rangle})(\cancel{\langle 01|10\rangle} + \langle 10|10\rangle) + \frac{1}{2}(\cancel{\langle 01|01\rangle} + \cancel{\langle 01|10\rangle})(\cancel{\langle 01|11\rangle} + \cancel{\langle 10|11\rangle}) \\ = 0$$

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Example 10.2.1 (cont.)

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Thus $\rho_y^A = \frac{1}{2}I = \rho_y^B = \rho_{ME}$ and the entropy, $S(\rho_{ME}) = 1$ for this state also

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Any other 2-qubit maximally entangled state will give the same results



Example 10-2-2

Compute the partial density operator for the first qubit of the state



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$$|x\rangle = \frac{7}{10}|00\rangle + \frac{1}{10}|01\rangle + \frac{1}{10}|10\rangle + \frac{7}{10}|11\rangle$$



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$$\sum_{j=0}^1 \langle 0j|x\rangle\langle x|0j\rangle|0\rangle\langle 0| = \left[\left(\frac{7}{10}\right)^2 + \left(\frac{1}{10}\right)^2 \right] |0\rangle\langle 0|$$



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Example 10-2-2 (cont.)

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$$0 = \det \begin{vmatrix} \frac{1}{2} - \lambda & \frac{7}{50} \\ \frac{7}{50} & \frac{1}{2} - \lambda \end{vmatrix}$$

Example 10-2-2 (cont.)

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$$\lambda = \frac{1}{2} \pm \frac{1}{2} \sqrt{1 - 4 \frac{1}{4} \left[1 - \left(\frac{7}{50} \right)^2 \right]}$$

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$$\lambda = \frac{1}{2} \pm \frac{1}{2} \sqrt{1 - 4 \frac{1}{4} \left[1 - \left(\frac{7}{50} \right)^2 \right]} = \frac{16}{25}, \frac{9}{25}$$

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$$S(\rho_x^A) = -\frac{16}{25} \log_2 \frac{16}{25} - \frac{9}{25} \log_2 \frac{9}{25} = 0.942$$



Example 10.2.4

Determine the amount of entanglement in the 4-qubit state in the 2,4 and 1,2 subsystems



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$$|\psi\rangle = \frac{1}{2}(|00\rangle + |11\rangle + |22\rangle + |33\rangle) = \frac{1}{2}(|0000\rangle + |0101\rangle + |1010\rangle + |1111\rangle)$$



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$$|\psi\rangle = \frac{1}{\sqrt{2}}(|0\rangle_1|0\rangle_3 + |1\rangle_1|1\rangle_3) \otimes \frac{1}{\sqrt{2}}(|0\rangle_2|0\rangle_4 + |1\rangle_2|1\rangle_4)$$



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$$\rho_{\psi}^{1,2} = \text{Tr}_{3,4}(|\psi\rangle\langle\psi|)$$

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In the 1,2 and 3,4 decomposition, the partial density operator becomes

$$\rho_{\psi}^{1,2} = \text{Tr}_{3,4}(|\psi\rangle\langle\psi|) = \sum_{i,j=0}^3 \sum_{k=0}^3 \langle j_3| \langle k_4 | |\psi\rangle\langle\psi| | i_3 \rangle | k_4 \rangle | j \rangle \langle i |$$

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Determine the amount of entanglement in the 4-qubit state in the 2,4 and 1,2 subsystems

$$|\psi\rangle = \frac{1}{2}(|00\rangle + |11\rangle + |22\rangle + |33\rangle) = \frac{1}{2}(|0000\rangle + |0101\rangle + |1010\rangle + |1111\rangle)$$

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In this decomposition the state is
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