

Today's outline - March 31, 2022



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- Properties of density operators

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- Geometry of mixed states

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Reading assignment: 10.2 – 10.3

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Homework Assignment #06:

See Blackboard

Due Tuesday, April 05, 2022

Quantum circuit simulator <https://algassert.com/quirk>

More properties of density operators



Any density operator, ρ_x^A , is Hermitian and $\text{Tr}(\rho_x^A) = \sum_j \overline{x_{ij}} x_{ij} \equiv 1$ since $|x\rangle$ is a unit vector

More properties of density operators



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$$\langle v | \rho_x^A | v \rangle = \sum_{i=0}^{M-1} \sum_{k=0}^{M-1} \sum_{j=0}^{L-1} \langle v | (\overline{x_{ij}} x_{kj} | \alpha_k \rangle \langle \alpha_i |) | v \rangle$$

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$$\sum_{i=0}^{M-1} \lambda_i \equiv 1, \quad \rho_x^A = \sum_{i=0}^{M-1} \lambda_i |v_i\rangle \langle v_i|$$

More properties of density operators



Any operator which satisfies the three conditions above must be a density operator

More properties of density operators



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If ρ is an operator acting on A of dimension $M = 2^m$ with eigenbasis $\{|\psi_0\rangle, \dots, |\psi_{M-1}\rangle\}$ which satisfies all three conditions



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Let B be a quantum system with a vector space of dimension $2^n > M$ and let $\{|\phi_0\rangle, \dots, |\phi_{M-1}\rangle\}$ be the first M elements of an orthonormal basis for B , define



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$|x\rangle \in A \otimes B$ is a so-called pure state which satisfies $\rho_x^A = \rho$



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The density operator $\rho_x^X = |x\rangle \langle x|$ for a pure state $|x\rangle$ is all zeros except for a 1 in the i^{th} diagonal element where $|x\rangle$ is the i^{th} element in a basis

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The density operator of a pure state is a projection operator, such that $\rho_x^X \rho_x^X = \rho_x^X$

More properties of density operators



Projection operators of pure states also eliminate the issue associated with global phase differences

More properties of density operators



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The density matrix for **mixed** and **pure** states are very different

More properties of density operators



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For the pure state $|+\rangle$ and the evenly mixed ensemble of $|0\rangle$ and $|1\rangle$ we have

$$\rho_{mixed} = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix},$$

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Geometry of mixed states



It is possible to use the Bloch sphere to visualize single-qubit mixed states which are linear combinations of pure states with non-negative coefficients that sum to 1

Geometry of mixed states



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$$\rho = \begin{pmatrix} a & c - id \\ c + id & b \end{pmatrix}$$

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$$= \frac{1}{2} \begin{pmatrix} 1 + z & x - iy \\ x + iy & 1 - z \end{pmatrix} = \frac{1}{2}(I + x\sigma_x + y\sigma_y + z\sigma_z)$$

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$$\det(\rho) = \begin{vmatrix} \frac{1+z}{2} & \frac{x-iy}{2} \\ \frac{x+iy}{2} & \frac{1-z}{2} \end{vmatrix}$$

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$$= \frac{1}{2} \begin{pmatrix} 1+z & x-iy \\ x+iy & 1-z \end{pmatrix} = \frac{1}{2}(I + x\sigma_x + y\sigma_y + z\sigma_z)$$

$$\sigma_x = X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_y = -iY = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_z = Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

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Geometry of mixed states



It is possible to use the Bloch sphere to visualize single-qubit mixed states which are linear combinations of pure states with non-negative coefficients that sum to 1

A density operator for a single qubit must be Hermitian and self-adjoint with trace 1 and the general form

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Density matrix and the Bloch sphere



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A maximally uncertain state for an n -qubit system has all the diagonal elements equal to 2^{-n} so $S(\rho) = n$

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The von Neumann entropy for a single qubit system is just a function of the distance of the state from the center of the Bloch sphere

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$$\rho_{ME} = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

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$$\rho_{ME} = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

Bipartite entanglement



It is useful to find a good measure of entanglement for bipartite systems such as $X = A \otimes B$

The 2-qubit system is the simplest bipartite system with a maximally entangled state

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For each of the two qubits, the density matrix ρ_{ME} has maximal von Neumann entropy

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It makes sense to use the von Neumann entropy of the partial trace as a measure of the entanglement if it can be assumed that the partial trace is the same for each of the two subsystems, A and B

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If $|\psi\rangle$ is a pure state of the system $A \otimes B$, there exists orthonormal sets of states $\{|\psi_i^A\rangle\}$ and $\{|\psi_i^B\rangle\}$

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The amount of entanglement between the two parts of a pure state $|\psi\rangle \in X = A \otimes B$ with density operator $\rho = |\psi\rangle\langle\psi|$ is defined to be $S(\mathrm{Tr}_A(\rho))$ or $S(\mathrm{Tr}_B(\rho))$

Example 10.2.1



Given a 2-qubit system in the maximally entangled Bell state



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Example 10.2.1 (cont.)



$$(\rho_y^A)_{01} = \frac{1}{2}(\langle 00|01\rangle + \langle 00|10\rangle)(\langle 01|10\rangle + \langle 10|10\rangle) + \frac{1}{2}(\langle 01|01\rangle + \langle 01|10\rangle)(\langle 01|11\rangle + \langle 10|11\rangle)$$

Example 10.2.1 (cont.)



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Example 10.2.1 (cont.)



$$\begin{aligned}(\rho_y^A)_{01} &= \frac{1}{2}(\langle \cancel{00|01} \rangle + \langle \cancel{00|10} \rangle)(\langle \cancel{01|10} \rangle + \langle 10|10 \rangle) + \frac{1}{2}(\langle 01|01 \rangle + \langle \cancel{01|10} \rangle)(\langle \cancel{01|11} \rangle + \langle \cancel{10|11} \rangle) \\ &= 0\end{aligned}$$

$$\begin{aligned}(\rho_y^A)_{10} &= \frac{1}{2}(\langle \cancel{10|01} \rangle + \langle 10|10 \rangle)(\langle \cancel{01|00} \rangle + \langle \cancel{10|00} \rangle) + \frac{1}{2}(\langle \cancel{11|01} \rangle + \langle \cancel{11|10} \rangle)(\langle 01|01 \rangle + \langle \cancel{10|01} \rangle) \\ &= 0\end{aligned}$$

Example 10.2.1 (cont.)



$$\begin{aligned}(\rho_y^A)_{01} &= \frac{1}{2}(\langle \cancel{00|01} \rangle + \langle \cancel{00|10} \rangle)(\langle \cancel{01|10} \rangle + \langle 10|10 \rangle) + \frac{1}{2}(\langle 01|01 \rangle + \langle \cancel{01|10} \rangle)(\langle \cancel{01|11} \rangle + \langle \cancel{10|11} \rangle) \\ &= 0\end{aligned}$$

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$$(\rho_y^A)_{11} = \frac{1}{2}(\langle 10|01 \rangle + \langle 10|10 \rangle)(\langle 01|10 \rangle + \langle 10|10 \rangle) + \frac{1}{2}(\langle 11|01 \rangle + \langle 11|10 \rangle)(\langle 01|11 \rangle + \langle 10|11 \rangle)$$

Example 10.2.1 (cont.)



$$\begin{aligned}(\rho_y^A)_{01} &= \frac{1}{2}(\langle \cancel{00|01} \rangle + \langle \cancel{00|10} \rangle)(\langle \cancel{01|10} \rangle + \langle 10|10 \rangle) + \frac{1}{2}(\langle 01|01 \rangle + \langle \cancel{01|10} \rangle)(\langle \cancel{01|11} \rangle + \langle \cancel{10|11} \rangle) \\ &= 0\end{aligned}$$

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Example 10.2.1 (cont.)



$$\begin{aligned}(\rho_y^A)_{01} &= \frac{1}{2}(\langle \cancel{00|01} \rangle + \langle \cancel{00|10} \rangle)(\langle \cancel{01|10} \rangle + \langle 10|10 \rangle) + \frac{1}{2}(\langle 01|01 \rangle + \langle \cancel{01|10} \rangle)(\langle \cancel{01|11} \rangle + \langle \cancel{10|11} \rangle) \\ &= 0\end{aligned}$$

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Thus $\rho_y^A = \frac{1}{2}I = \rho_y^B = \rho_{ME}$ and the entropy, $S(\rho_{ME}) = 1$ for this state also

Example 10.2.1 (cont.)



$$(\rho_y^A)_{01} = \frac{1}{2}(\langle \cancel{00|01} \rangle + \langle \cancel{00|10} \rangle)(\langle \cancel{01|10} \rangle + \langle 10|10 \rangle) + \frac{1}{2}(\langle 01|01 \rangle + \langle \cancel{01|10} \rangle)(\langle \cancel{01|11} \rangle + \langle \cancel{10|11} \rangle) \\ = 0$$

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Thus $\rho_y^A = \frac{1}{2}I = \rho_y^B = \rho_{ME}$ and the entropy, $S(\rho_{ME}) = 1$ for this state also

Any other 2-qubit maximally entangled state will give the same results

Example 10-2-2



Compute the partial density operator for the first qubit of the state

Example 10-2-2



Compute the partial density operator for the first qubit of the state

$$|x\rangle = \frac{7}{10}|00\rangle + \frac{1}{10}|01\rangle + \frac{1}{10}|10\rangle + \frac{7}{10}|11\rangle$$

Example 10-2-2



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The partial density matrix is defined as $\rho_x^A = \text{Tr}_B(|x\rangle\langle x|)$, a matrix which has 4 elements

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$$\sum_{j=0}^1 \langle 0j|x\rangle \langle x|0j\rangle |0\rangle\langle 0| = \left[\left(\frac{7}{10}\right)^2 + \left(\frac{1}{10}\right)^2 \right] |0\rangle\langle 0|$$

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$$\sum_{j=0}^1 \langle 0j|x\rangle\langle x|1j\rangle|0\rangle\langle 1| = \left[\frac{7}{10} \frac{1}{10} + \frac{1}{10} \frac{7}{10} \right] |0\rangle\langle 1|$$

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$$\sum_{j=0}^1 \langle 1j|x\rangle \langle x|0j\rangle |1\rangle\langle 0| = \left[\frac{1}{10} \frac{7}{10} + \frac{7}{10} \frac{1}{10} \right] |1\rangle\langle 0|$$

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Example 10-2-2 (cont.)



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$$\rho_x^A = \frac{1}{2}|0\rangle\langle 0| + \frac{7}{50}|0\rangle\langle 1| + \frac{1}{10}|1\rangle\langle 0| + \frac{1}{2}|1\rangle\langle 1|$$

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Example 10-2-2 (cont.)



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Example 10-2-2 (cont.)



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This corresponds to the point $(0.28, 0, 0)$ in the Bloch sphere

Example 10-2-2 (cont.)



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$$\rho_x^A = \frac{1}{2}|0\rangle\langle 0| + \frac{7}{50}|0\rangle\langle 1| + \frac{1}{10}|1\rangle\langle 0| + \frac{1}{2}|1\rangle\langle 1| \frac{1}{100} \begin{pmatrix} 50 & 14 \\ 14 & 50 \end{pmatrix} = \frac{1}{2} \left(I + \frac{14}{50} X \right)$$

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To get the entropy, diagonalize the ρ_x^A matrix

Example 10-2-2 (cont.)



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$$0 = \det \begin{vmatrix} \frac{1}{2} - \lambda & \frac{7}{50} \\ \frac{7}{50} & \frac{1}{2} - \lambda \end{vmatrix}$$

Example 10-2-2 (cont.)



Thus the partial density matrix becomes

$$\rho_x^A = \frac{1}{2}|0\rangle\langle 0| + \frac{7}{50}|0\rangle\langle 1| + \frac{1}{10}|1\rangle\langle 0| + \frac{1}{2}|1\rangle\langle 1| \frac{1}{100} \begin{pmatrix} 50 & 14 \\ 14 & 50 \end{pmatrix} = \frac{1}{2} \left(I + \frac{14}{50} X \right)$$

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Example 10-2-2 (cont.)



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$$\rho_x^A = \frac{1}{2}|0\rangle\langle 0| + \frac{7}{50}|0\rangle\langle 1| + \frac{1}{10}|1\rangle\langle 0| + \frac{1}{2}|1\rangle\langle 1| \frac{1}{100} \begin{pmatrix} 50 & 14 \\ 14 & 50 \end{pmatrix} = \frac{1}{2} \left(I + \frac{14}{50} X \right)$$

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$$\lambda = \frac{1}{2} \pm \frac{1}{2} \sqrt{1 - 4 \frac{1}{4} \left[1 - \left(\frac{7}{50} \right)^2 \right]}$$

Example 10-2-2 (cont.)



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$$\rho_x^A = \frac{1}{2}|0\rangle\langle 0| + \frac{7}{50}|0\rangle\langle 1| + \frac{1}{10}|1\rangle\langle 0| + \frac{1}{2}|1\rangle\langle 1| \frac{1}{100} \begin{pmatrix} 50 & 14 \\ 14 & 50 \end{pmatrix} = \frac{1}{2} (I + \frac{14}{50}X)$$

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$$\lambda = \frac{1}{2} \pm \frac{1}{2} \sqrt{1 - 4 \frac{1}{4} \left[1 - \left(\frac{7}{50} \right)^2 \right]} = \frac{16}{25}, \frac{9}{25}$$

Example 10-2-2 (cont.)



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$$S(\rho_x^A) = -\frac{16}{25} \log_2 \frac{16}{25} - \frac{9}{25} \log_2 \frac{9}{25} = 0.942$$

Example 10.2.4



Determine the amount of entanglement in the 4-qubit state in the 2,4 and 1,2 subsystems

Example 10.2.4



Determine the amount of entanglement in the 4-qubit state in the 2,4 and 1,2 subsystems

$$|\psi\rangle = \frac{1}{2}(|00\rangle + |11\rangle + |22\rangle + |33\rangle) = \frac{1}{2}(|0000\rangle + |0101\rangle + |1010\rangle + |1111\rangle)$$

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$$|\psi\rangle = \frac{1}{\sqrt{2}}(|0\rangle_1|0\rangle_3 + |1\rangle_1|1\rangle_3) \otimes \frac{1}{\sqrt{2}}(|0\rangle_2|0\rangle_4 + |1\rangle_2|1\rangle_4)$$

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Since the state is unentangled, it is a pure state in the 2,4 subsystem and $S(\rho_\psi^{2,4}) \equiv 0$

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Since the state is unentangled, it is a pure state in the 2,4 subsystem and $S(\rho_\psi^{2,4}) \equiv 0$

In the 1,2 and 3,4 decomposition, the partial density operator becomes

$$\rho_\psi^{1,2} = \text{Tr}_{3,4}(|\psi\rangle\langle\psi|)$$

Example 10.2.4



Determine the amount of entanglement in the 4-qubit state in the 2,4 and 1,2 subsystems

$$|\psi\rangle = \frac{1}{2}(|00\rangle + |11\rangle + |22\rangle + |33\rangle) = \frac{1}{2}(|0000\rangle + |0101\rangle + |1010\rangle + |1111\rangle)$$

In the 2,4 decomposition, this state is unentangled

$$|\psi\rangle = \frac{1}{\sqrt{2}}(|0\rangle_1|0\rangle_3 + |1\rangle_1|1\rangle_3) \otimes \frac{1}{\sqrt{2}}(|0\rangle_2|0\rangle_4 + |1\rangle_2|1\rangle_4)$$

Since the state is unentangled, it is a pure state in the 2,4 subsystem and $S(\rho_\psi^{2,4}) \equiv 0$

In the 1,2 and 3,4 decomposition, the partial density operator becomes

$$\rho_\psi^{1,2} = \text{Tr}_{3,4}(|\psi\rangle\langle\psi|) = \sum_{i,j=0}^3 \sum_{k=0}^3 \langle j_3 | \langle k_4 | |\psi\rangle\langle\psi| | i_3 \rangle | k_4 \rangle | j \rangle \langle i |$$

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In this decomposition the state is maximally entangled

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