

Today's outline - February 15, 2022



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- Deutch's problem

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- Deutch's problem
- Quantum subroutines

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- State-dependent phase shift

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- State-dependent amplitude shift

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Reading Assignment: Chapter 7.5-7.7

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Reading Assignment: Chapter 7.5-7.7

Homework Assignment #05:

Chapter 7:1,3,4

due Thursday, February 24, 2022



Deutsch's algorithm

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For $f(x)$ **constant**, both terms pick up the same phase shift and the state is $|+\rangle|-\rangle$



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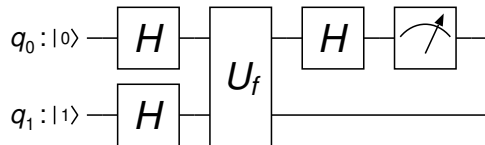
For $f(x)$ **constant**, both terms pick up the same phase shift and the state is $|+\rangle|-\rangle$

For $f(x)$ **balanced**, only one term picks up a phase shift, giving a result of $|-\rangle|-\rangle$

Deutsch's algorithm



As a quantum circuit, Deutsch's algorithm is implemented by

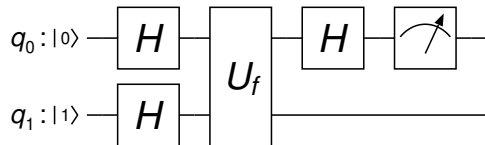


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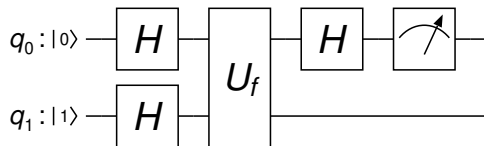


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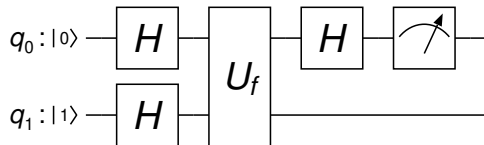


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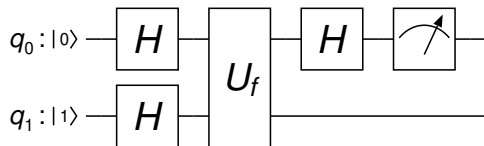


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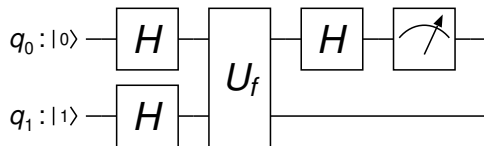


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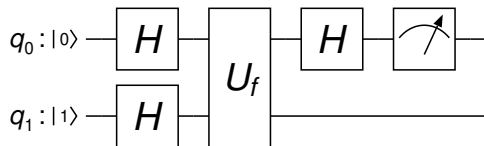


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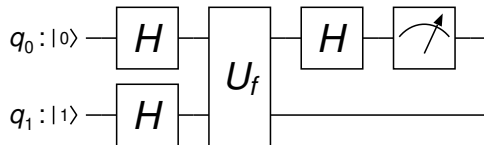
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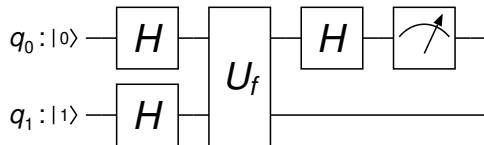


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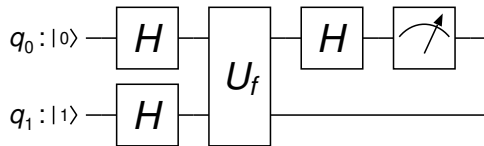
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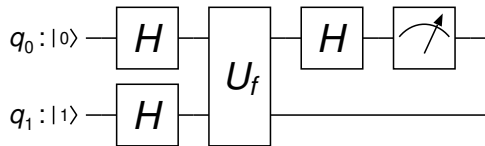
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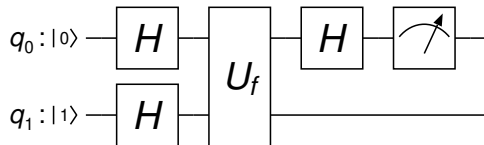
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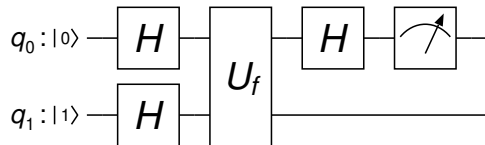
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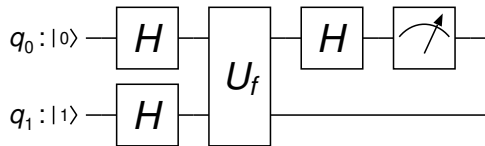
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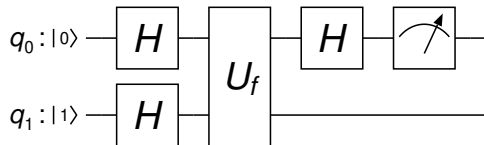
$f(x)$	$U_f q_0\rangle q_1\rangle$	q_0
$f(0) \rightarrow 0; f(1) \rightarrow 0$	$+ +\rangle -\rangle$	$ +\rangle \rightarrow 0\rangle$
$f(0) \rightarrow 1; f(1) \rightarrow 1$	$- +\rangle -\rangle$	$ +\rangle \rightarrow 0\rangle$
$f(0) \rightarrow 0; f(1) \rightarrow 1$	$+ -\rangle -\rangle$	$ -\rangle \rightarrow 1\rangle$

Deutsch's algorithm



As a quantum circuit, Deutsch's algorithm is implemented by

1. prepare two qubits: $q_0 = |0\rangle$ and $q_1 = |1\rangle$
2. apply the Hadamard transform to each qubit
3. apply the black box algorithm
4. apply the Hadamard transform to q_0
5. measure q_0 and interpret



(a) $q_0 = 0 \rightarrow f(x)$ is **constant**

(b) $q_0 = 1 \rightarrow f(x)$ is **balanced**

$$q_0 : |0\rangle \rightarrow |+\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$$

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$f(0) \rightarrow 0; f(1) \rightarrow 0$	$+\textcolor{red}{ +\rangle} -\rangle$	$\textcolor{red}{ +\rangle} \rightarrow \textcolor{red}{ 0\rangle}$
$f(0) \rightarrow 1; f(1) \rightarrow 1$	$-\textcolor{red}{ +\rangle} -\rangle$	$\textcolor{red}{ +\rangle} \rightarrow \textcolor{red}{ 0\rangle}$
$f(0) \rightarrow 0; f(1) \rightarrow 1$	$+\textcolor{blue}{ -\rangle} -\rangle$	$\textcolor{blue}{ -\rangle} \rightarrow \textcolor{blue}{ 1\rangle}$
$f(0) \rightarrow 1; f(1) \rightarrow 0$	$-\textcolor{blue}{ -\rangle} -\rangle$	$\textcolor{blue}{ -\rangle} \rightarrow \textcolor{blue}{ 1\rangle}$

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For this reason, it is essential to uncompute $|y_i\rangle$ inside the subroutine before the output qubits are transmitted

Entanglement problems

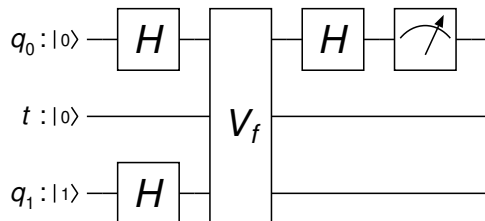


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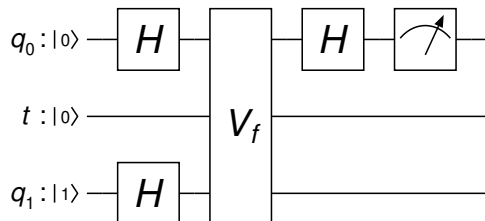


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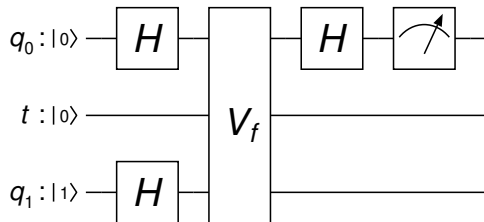
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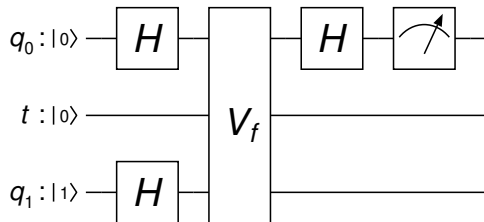


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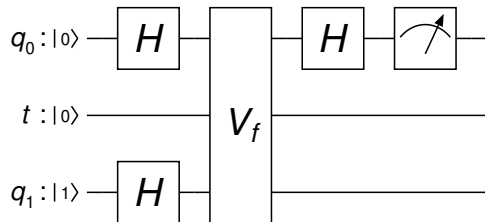
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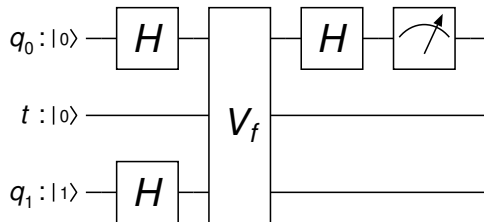
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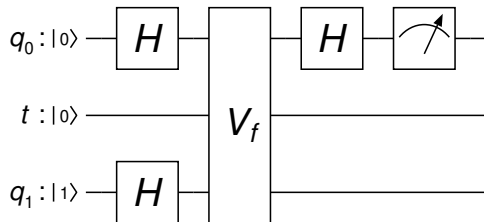
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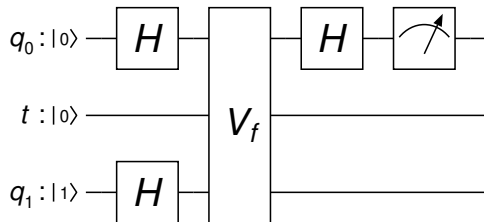
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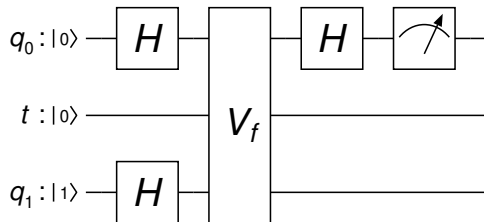
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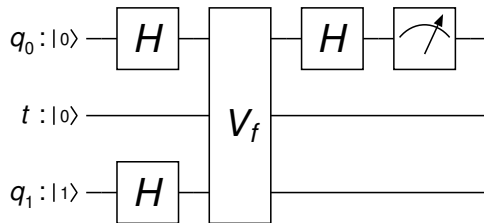


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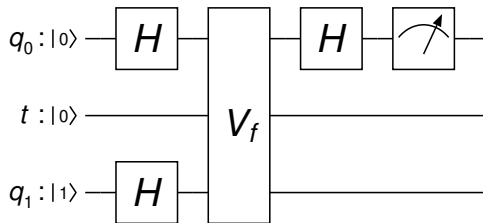
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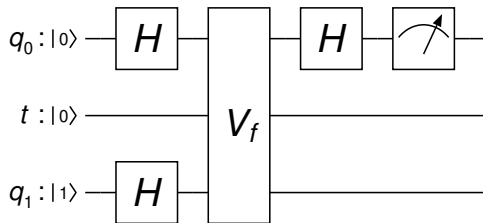


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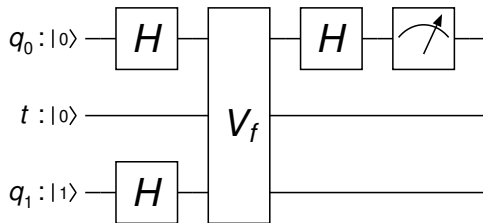
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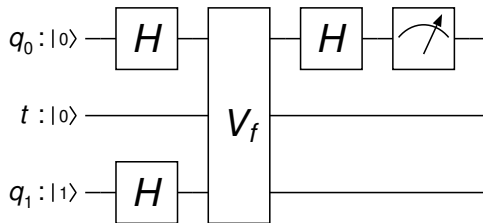
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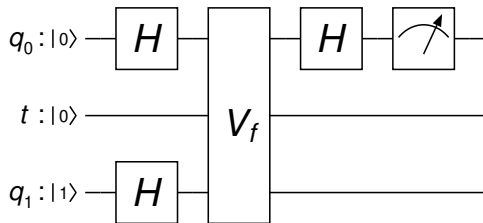
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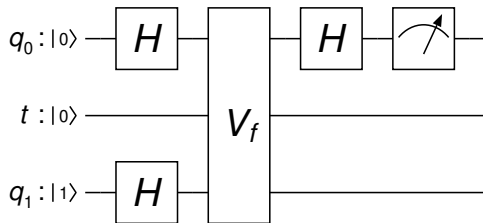
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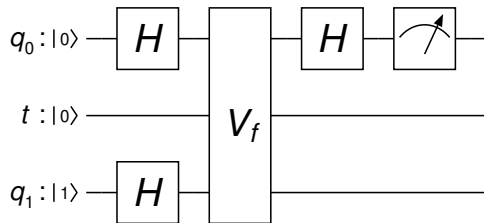
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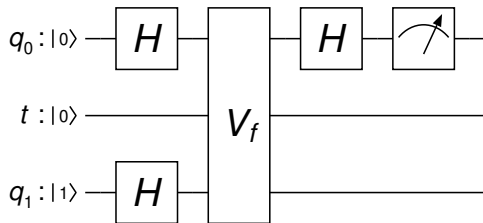
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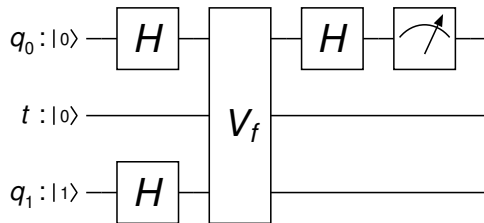
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The final step of applying the Hadamard transformation to q_0 uses the $H \otimes I \otimes I$ transformation which yields



$$f(x) \text{ constant} \quad H \otimes I \otimes I(|00\rangle + |11\rangle)|-\rangle = \frac{1}{2}(|\textcolor{red}{0}0\rangle + |\textcolor{blue}{1}0\rangle + |\textcolor{red}{0}1\rangle - |\textcolor{blue}{1}1\rangle)|-\rangle$$

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The q_0 qubit can be measured to be $|\textcolor{red}{0}\rangle$ or $|\textcolor{blue}{1}\rangle$ with equal probability in both cases since two of the terms do not cancel but by uncomputing the t qubit, the states would be left as

$$f(x) \text{ constant} \quad H \otimes I \otimes I(|00\rangle + |10\rangle)|-\rangle = \frac{1}{2}(|\textcolor{red}{0}0\rangle + \cancel{|\textcolor{blue}{1}0\rangle} + |\textcolor{red}{0}0\rangle - \cancel{|\textcolor{blue}{1}0\rangle})|-\rangle = |\textcolor{red}{0}\rangle|0\rangle|-\rangle$$

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The depends on being able to compute membership in X efficiently but if this is possible with a quantum transformation U_f then through the use of a temporary qubit, it is possible to compute S_X^ϕ

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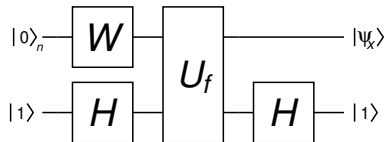
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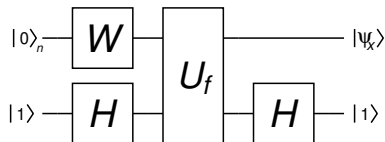


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The circuit starts with a uniform superposition of an n -qubit register and an ancilla qubit in the $|1\rangle$ state to create the superposition $|\psi_X\rangle = \sum (-1)^{f(x)} |x\rangle$

State-dependent phase changes



Suppose we wish to apply a phase shift that depends on the state of a specific qubit, $|x\rangle \rightarrow e^{i\phi(x)}|x\rangle$ where there is an associated function $f : \mathbf{Z}_n \rightarrow \mathbf{Z}_s$ that is efficiently computable

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Given a transformation U_f that is efficient, it is possible to perform the state-dependent phase shift in $O(s)$ steps plus 2 invocations of U_f

Suppose that $f(x) = x$, we want a subroutine that changes the phase of an s -qubit standard basis state $|x\rangle$ by $\phi(x) = 2\pi x/2^s$ using the transformation

$$P(\phi) = T\left(-\frac{\phi}{2}\right) K\left(\frac{\phi}{2}\right) = \begin{pmatrix} 1 & 0 \\ 0 & e^{i\phi} \end{pmatrix}$$

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State-dependent phase changes



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State-dependent phase changes



Using the subroutine $Phase : |a\rangle \rightarrow e^{i2\pi s/2^s}$ it is now possible to write a program that implements the n -qubit transformation $Phase_f : |x\rangle \rightarrow e^{i2\pi f(x)/2^s}$

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Step 4 unentangles $|a\rangle$ from $|x\rangle$ leaving it in the desired state

State-dependent amplitude shifts



We wish to rotate each term in a superposition by a single qubit rotation $R(\beta(x))$ where $\beta(x)$ is state-dependent such that $|x\rangle \otimes |b\rangle \rightarrow |x\rangle \otimes (R(\beta(x))|b\rangle)$

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