### The variational method





• The variational theorem



- The variational theorem
- Proof of the variational theorem



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- Harmonic oscillator ground state energy



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- Delta function well ground state energy



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 $\psi = \sum_{n} c_{n} \psi_{n}$ 

where  $\psi_n$  are the (unknown) eigenfunctions of H, which form a complete set



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This is the so-called variational principle which allows us to compute an upper bound on the ground state energy and, if we make a judicious choice of arbitrary wave function,  $\psi$  (ie. not arbitrary at all!) we can get very close to the actual ground state energy.

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In practice this means we minimize the value of the Hamiltonian expectation value to achieve this upper bound.



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 $\langle H \rangle = \langle T \rangle + \langle V \rangle$ 

# Example 8.1 (cont.)



The kinetic and potential energies are instances of the Gaussian integral


$$\langle T \rangle = -\frac{\hbar^2}{2m} |A|^2 \int_{-\infty}^{\infty} e^{-bx^2} \frac{d^2}{dx^2} \left( e^{-bx^2} \right) dx$$



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$$=\frac{1}{r_1r_2}\left[\sqrt{r_1^2+r_2^2+2r_1r_2}-\sqrt{r_1^2+r_2^2-2r_1r_2}\right] =\frac{1}{r_1r_2}\left[(r_1+r_2)-|r_1-r_2|\right]$$

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$$\langle V_{ee} \rangle = \left(\frac{e^2}{4\pi\epsilon_0}\right) \left(\frac{8}{\pi a^3}\right)^2 \int \frac{e^{-4r_1/a} e^{-4r_2/a}}{\sqrt{r_1^2 + r_2^2 - 2r_1r_2\cos\theta_2}} d^3\vec{r_1} d^3\vec{r_2}$$

the  $\vec{r_2}$  integral then can be evaluated

$$I_2 \equiv \int \frac{e^{-4r_2/a} r_2^2 \sin \theta_2}{\sqrt{r_1^2 + r_2^2 - 2r_1 r_2 \cos \theta_2}} \, dr_2 \, d\theta_2 \, d\phi_2$$

the  $\phi_2$  integral gives  $2\pi$  and the  $\theta_2$  integrand is a perfect derivative

$$\begin{split} &I_{\theta} = \int_{0}^{\pi} \frac{\sin \theta_{2}}{\sqrt{r_{1}^{2} + r_{2}^{2} - 2r_{1}r_{2}\cos\theta_{2}}} \, d\theta_{2} = \left. \frac{\sqrt{r_{1}^{2} + r_{2}^{2} - 2r_{1}r_{2}\cos\theta_{2}}}{r_{1}r_{2}} \right|_{0}^{\pi} \\ &= \frac{1}{r_{1}r_{2}} \left[ \sqrt{r_{1}^{2} + r_{2}^{2} + 2r_{1}r_{2}} - \sqrt{r_{1}^{2} + r_{2}^{2} - 2r_{1}r_{2}} \right] = \frac{1}{r_{1}r_{2}} \Big[ (r_{1} + r_{2}) - |r_{1} - r_{2}| \Big] = \left\{ \end{split}$$

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$$\langle V_{ee} \rangle = \left(\frac{e^2}{4\pi\epsilon_0}\right) \left(\frac{8}{\pi a^3}\right)^2 \int \frac{e^{-4r_1/a} e^{-4r_2/a}}{\sqrt{r_1^2 + r_2^2 - 2r_1r_2\cos\theta_2}} d^3\vec{r_1} d^3\vec{r_2}$$

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the  $\phi_2$  integral gives  $2\pi$  and the  $\theta_2$  integrand is a perfect derivative

$$\begin{split} & l_{\theta} = \int_{0}^{\pi} \frac{\sin \theta_{2}}{\sqrt{r_{1}^{2} + r_{2}^{2} - 2r_{1}r_{2}\cos\theta_{2}}} \, d\theta_{2} = \left. \frac{\sqrt{r_{1}^{2} + r_{2}^{2} - 2r_{1}r_{2}\cos\theta_{2}}}{r_{1}r_{2}} \right|_{0}^{\pi} \\ & = \frac{1}{r_{1}r_{2}} \left[ \sqrt{r_{1}^{2} + r_{2}^{2} + 2r_{1}r_{2}} - \sqrt{r_{1}^{2} + r_{2}^{2} - 2r_{1}r_{2}} \right] = \frac{1}{r_{1}r_{2}} \Big[ (r_{1} + r_{2}) - |r_{1} - r_{2}| \Big] = \begin{cases} 2/r_{1}, & r_{2} < r_{1} \\ r_{2} < r_{1} \end{cases} \end{split}$$

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$$\langle V_{ee} \rangle = \left(\frac{e^2}{4\pi\epsilon_0}\right) \left(\frac{8}{\pi a^3}\right)^2 \int \frac{e^{-4r_1/a} e^{-4r_2/a}}{\sqrt{r_1^2 + r_2^2 - 2r_1r_2\cos\theta_2}} d^3\vec{r_1} d^3\vec{r_2}$$

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the  $\phi_2$  integral gives  $2\pi$  and the  $\theta_2$  integrand is a perfect derivative

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# V

#### The $r_2$ integral

$$I_2 = 4\pi \left( + \right)$$

# V

#### The $r_2$ integral

$$I_2 = 4\pi \left( \frac{1}{r_1} \int_0^{r_1} r_2^2 e^{-4r_2/a} \, dr_2 + \right)$$



$$I_{2} = 4\pi \left( \frac{1}{r_{1}} \int_{0}^{r_{1}} r_{2}^{2} e^{-4r_{2}/a} dr_{2} + \int_{r_{1}}^{\infty} r_{2} e^{-4r_{2}/a} dr_{2} \right)$$



The  $r_2$  integral is thus split into two parts

$$I_2 = 4\pi \left( \frac{1}{r_1} \int_0^{r_1} r_2^2 e^{-4r_2/a} dr_2 + \int_{r_1}^{\infty} r_2 e^{-4r_2/a} dr_2 \right)$$

these can be solved using integration by parts with  $dv = e^{-4r_2/a} dr_2$


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$$I_{2} = 4\pi \left( \frac{1}{r_{1}} \int_{0}^{r_{1}} r_{2}^{2} e^{-4r_{2}/a} dr_{2} + \int_{r_{1}}^{\infty} r_{2} e^{-4r_{2}/a} dr_{2} \right)$$

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 $I_{red} =$ 



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$$I_{red} = -\left(rac{a}{4}
ight) r_2^2 e^{-4r_2/a} \Big|_0^{r_1}$$



The  $r_2$  integral is thus split into two parts

$$I_{2} = 4\pi \left( \frac{1}{r_{1}} \int_{0}^{r_{1}} r_{2}^{2} e^{-4r_{2}/a} dr_{2} + \int_{r_{1}}^{\infty} r_{2} e^{-4r_{2}/a} dr_{2} \right)$$

$$I_{red} = -\left(\frac{a}{4}\right) r_2^2 e^{-4r_2/a} \Big|_0^{r_1} + 2\left(\frac{a}{4}\right) \int_0^{r_1} r_2 e^{-4r_2/a} dr_2$$



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$$I_{red} = -\left(\frac{a}{4}\right) r_2^2 e^{-4r_2/a} \Big|_0^{r_1} + 2\left(\frac{a}{4}\right) \int_0^{r_1} r_2 e^{-4r_2/a} dr_2$$
$$= -\left(\frac{a}{4}\right) r_1^2 e^{-4r_1/a}$$



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$$I_{2} = 4\pi \left( \frac{1}{r_{1}} \int_{0}^{r_{1}} r_{2}^{2} e^{-4r_{2}/a} dr_{2} + \int_{r_{1}}^{\infty} r_{2} e^{-4r_{2}/a} dr_{2} \right)$$

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$$= -\left(\frac{a}{4}\right) r_1^2 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^2 r_2 e^{-4r_2/a} \Big|_0^{r_1}$$



The  $r_2$  integral is thus split into two parts

$$I_{2} = 4\pi \left( \frac{1}{r_{1}} \int_{0}^{r_{1}} r_{2}^{2} e^{-4r_{2}/a} dr_{2} + \int_{r_{1}}^{\infty} r_{2} e^{-4r_{2}/a} dr_{2} \right)$$

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$$= -\left(\frac{a}{4}\right) r_1^2 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^2 r_2 e^{-4r_2/a} \Big|_0^{r_1} + 2\left(\frac{a}{4}\right)^2 \int_0^{r_1} e^{-4r_2/a} dr_2$$

I



The  $r_2$  integral is thus split into two parts

$$I_2 = 4\pi \left( \frac{1}{r_1} \int_0^{r_1} r_2^2 e^{-4r_2/a} dr_2 + \int_{r_1}^{\infty} r_2 e^{-4r_2/a} dr_2 \right)$$

these can be solved using integration by parts with  $dv = e^{-4r_2/a} dr_2$ 

$$\begin{aligned} red &= -\left(\frac{a}{4}\right) r_2^2 e^{-4r_2/a} \Big|_0^{r_1} + 2\left(\frac{a}{4}\right) \int_0^{r_1} r_2 e^{-4r_2/a} \, dr_2 \\ &= -\left(\frac{a}{4}\right) r_1^2 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^2 r_2 e^{-4r_2/a} \Big|_0^{r_1} + 2\left(\frac{a}{4}\right)^2 \int_0^{r_1} e^{-4r_2/a} \, dr_2 \\ &= -\left(\frac{a}{4}\right) r_1^2 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^2 r_1 e^{-4r_1/a} \end{aligned}$$

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I



The  $r_2$  integral is thus split into two parts

$$I_{2} = 4\pi \left( \frac{1}{r_{1}} \int_{0}^{r_{1}} r_{2}^{2} e^{-4r_{2}/a} dr_{2} + \int_{r_{1}}^{\infty} r_{2} e^{-4r_{2}/a} dr_{2} \right)$$

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The  $r_2$  integral is thus split into two parts

$$I_{2} = 4\pi \left( \frac{1}{r_{1}} \int_{0}^{r_{1}} r_{2}^{2} e^{-4r_{2}/a} dr_{2} + \int_{r_{1}}^{\infty} r_{2} e^{-4r_{2}/a} dr_{2} \right)$$

these can be solved using integration by parts with  $dv = e^{-4r_2/a} dr_2$ 

$$I_{red} = -\left(\frac{a}{4}\right) r_2^2 e^{-4r_2/a} \Big|_0^{r_1} + 2\left(\frac{a}{4}\right) \int_0^{r_1} r_2 e^{-4r_2/a} dr_2$$
  
=  $-\left(\frac{a}{4}\right) r_1^2 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^2 r_2 e^{-4r_2/a} \Big|_0^{r_1} + 2\left(\frac{a}{4}\right)^2 \int_0^{r_1} e^{-4r_2/a} dr_2$   
=  $-\left(\frac{a}{4}\right) r_1^2 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^2 r_1 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^3 e^{-4r_1/a} + 2\left(\frac{a}{4}\right)^3$ 

 $I_{blue} =$ 

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The  $r_2$  integral is thus split into two parts

$$I_{2} = 4\pi \left( \frac{1}{r_{1}} \int_{0}^{r_{1}} r_{2}^{2} e^{-4r_{2}/a} dr_{2} + \int_{r_{1}}^{\infty} r_{2} e^{-4r_{2}/a} dr_{2} \right)$$

these can be solved using integration by parts with  $dv = e^{-4r_2/a} dr_2$ 

$$I_{red} = -\left(\frac{a}{4}\right) r_2^2 e^{-4r_2/a} \Big|_0^{r_1} + 2\left(\frac{a}{4}\right) \int_0^{r_1} r_2 e^{-4r_2/a} dr_2$$
  
$$= -\left(\frac{a}{4}\right) r_1^2 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^2 r_2 e^{-4r_2/a} \Big|_0^{r_1} + 2\left(\frac{a}{4}\right)^2 \int_0^{r_1} e^{-4r_2/a} dr_2$$
  
$$= -\left(\frac{a}{4}\right) r_1^2 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^2 r_1 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^3 e^{-4r_1/a} + 2\left(\frac{a}{4}\right)^3$$
  
$$I_{blue} = -\left(\frac{a}{4}\right) r_2 e^{-4r_2/a} \Big|_{r_1}^{\infty}$$

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The  $r_2$  integral is thus split into two parts

$$I_{2} = 4\pi \left( \frac{1}{r_{1}} \int_{0}^{r_{1}} r_{2}^{2} e^{-4r_{2}/a} dr_{2} + \int_{r_{1}}^{\infty} r_{2} e^{-4r_{2}/a} dr_{2} \right)$$

these can be solved using integration by parts with  $dv = e^{-4r_2/a} dr_2$ 

$$I_{red} = -\left(\frac{a}{4}\right) r_2^2 e^{-4r_2/a} \Big|_0^{r_1} + 2\left(\frac{a}{4}\right) \int_0^{r_1} r_2 e^{-4r_2/a} dr_2$$
  
$$= -\left(\frac{a}{4}\right) r_1^2 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^2 r_2 e^{-4r_2/a} \Big|_0^{r_1} + 2\left(\frac{a}{4}\right)^2 \int_0^{r_1} e^{-4r_2/a} dr_2$$
  
$$= -\left(\frac{a}{4}\right) r_1^2 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^2 r_1 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^3 e^{-4r_1/a} + 2\left(\frac{a}{4}\right)^3$$
  
$$I_{blue} = -\left(\frac{a}{4}\right) r_2 e^{-4r_2/a} \Big|_{r_1}^{\infty} + \left(\frac{a}{4}\right) \int_{r_1}^{\infty} e^{-4r_2/a} dr_2$$

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The  $r_2$  integral is thus split into two parts

$$I_{2} = 4\pi \left( \frac{1}{r_{1}} \int_{0}^{r_{1}} r_{2}^{2} e^{-4r_{2}/a} dr_{2} + \int_{r_{1}}^{\infty} r_{2} e^{-4r_{2}/a} dr_{2} \right)$$

these can be solved using integration by parts with  $dv = e^{-4r_2/a} dr_2$ 

$$I_{red} = -\left(\frac{a}{4}\right) r_2^2 e^{-4r_2/a} \Big|_0^{r_1} + 2\left(\frac{a}{4}\right) \int_0^{r_1} r_2 e^{-4r_2/a} dr_2$$
  
$$= -\left(\frac{a}{4}\right) r_1^2 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^2 r_2 e^{-4r_2/a} \Big|_0^{r_1} + 2\left(\frac{a}{4}\right)^2 \int_0^{r_1} e^{-4r_2/a} dr_2$$
  
$$= -\left(\frac{a}{4}\right) r_1^2 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^2 r_1 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^3 e^{-4r_1/a} + 2\left(\frac{a}{4}\right)^3$$
  
$$I_{blue} = -\left(\frac{a}{4}\right) r_2 e^{-4r_2/a} \Big|_{r_1}^{\infty} + \left(\frac{a}{4}\right) \int_{r_1}^{\infty} e^{-4r_2/a} dr_2 = + \left(\frac{a}{4}\right) r_1 e^{-4r_1/a}$$

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The  $r_2$  integral is thus split into two parts

$$I_{2} = 4\pi \left( \frac{1}{r_{1}} \int_{0}^{r_{1}} r_{2}^{2} e^{-4r_{2}/a} dr_{2} + \int_{r_{1}}^{\infty} r_{2} e^{-4r_{2}/a} dr_{2} \right)$$

these can be solved using integration by parts with  $dv = e^{-4r_2/a} dr_2$ 

$$\begin{split} I_{red} &= -\left(\frac{a}{4}\right) r_2^2 e^{-4r_2/a} \Big|_0^{r_1} + 2\left(\frac{a}{4}\right) \int_0^{r_1} r_2 e^{-4r_2/a} \, dr_2 \\ &= -\left(\frac{a}{4}\right) r_1^2 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^2 r_2 e^{-4r_2/a} \Big|_0^{r_1} + 2\left(\frac{a}{4}\right)^2 \int_0^{r_1} e^{-4r_2/a} \, dr_2 \\ &= -\left(\frac{a}{4}\right) r_1^2 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^2 r_1 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^3 e^{-4r_1/a} + 2\left(\frac{a}{4}\right)^3 \\ I_{blue} &= -\left(\frac{a}{4}\right) r_2 e^{-4r_2/a} \Big|_{r_1}^{\infty} + \left(\frac{a}{4}\right) \int_{r_1}^{\infty} e^{-4r_2/a} \, dr_2 \\ &= + \left(\frac{a}{4}\right) r_1 e^{-4r_1/a} + \left(\frac{a}{4}\right)^2 e^{-4r_1/a} - 2\left(\frac{a}{4}\right)^2 r_1 e^{-4r_1/a} + \left(\frac{a}{4}\right)^2 e^{-4r_1/a} \end{split}$$

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 $I_{2} = 4\pi \left[ \frac{1}{r_{1}} \int_{0}^{r_{1}} r_{2}^{2} e^{-4r_{2}/a} dr_{2} + \int_{r_{1}}^{\infty} r_{2} e^{-4r_{2}/a} dr_{2} \right]$ 



$$I_{2} = 4\pi \left[ \frac{1}{r_{1}} \int_{0}^{r_{1}} r_{2}^{2} e^{-4r_{2}/a} dr_{2} + \int_{r_{1}}^{\infty} r_{2} e^{-4r_{2}/a} dr_{2} \right]$$
  
=  $4\pi \left[ \frac{1}{r_{1}} \left( -\left(\frac{a}{4}\right) r_{1}^{2} e^{-4r_{1}/a} - 2\left(\frac{a}{4}\right)^{2} r_{1} e^{-4r_{1}/a} - 2\left(\frac{a}{4}\right)^{3} e^{-4r_{1}/a} + 2\left(\frac{a}{4}\right)^{3} \right) + \left(\frac{a}{4}\right) r_{1} e^{-4r_{1}/a} + \left(\frac{a}{4}\right)^{2} e^{-4r_{1}/a} \right]$ 

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$$I_{2} = 4\pi \left[ \frac{1}{r_{1}} \int_{0}^{r_{1}} r_{2}^{2} e^{-4r_{2}/a} dr_{2} + \int_{r_{1}}^{\infty} r_{2} e^{-4r_{2}/a} dr_{2} \right]$$
  
=  $4\pi \left[ \frac{1}{r_{1}} \left( -\left(\frac{a}{4}\right) r_{1}^{2} e^{-4r_{1}/a} - 2\left(\frac{a}{4}\right)^{2} r_{1} e^{-4r_{1}/a} - 2\left(\frac{a}{4}\right)^{3} e^{-4r_{1}/a} + 2\left(\frac{a}{4}\right)^{3} \right) + \left(\frac{a}{4}\right) r_{1} e^{-4r_{1}/a} + \left(\frac{a}{4}\right)^{2} e^{-4r_{1}/a} \right]$ 

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$$I_{2} = 4\pi \left[ \frac{1}{r_{1}} \int_{0}^{r_{1}} r_{2}^{2} e^{-4r_{2}/a} dr_{2} + \int_{r_{1}}^{\infty} r_{2} e^{-4r_{2}/a} dr_{2} \right]$$
  
=  $4\pi \left[ \frac{1}{r_{1}} \left( -\left(\frac{a}{4}\right) r_{1}^{2} e^{-4r_{1}/a} - 2\left(\frac{a}{4}\right)^{2} r_{1} e^{-4r_{1}/a} - 2\left(\frac{a}{4}\right)^{3} e^{-4r_{1}/a} + 2\left(\frac{a}{4}\right)^{3} \right) + \left(\frac{a}{4}\right) r_{1} e^{-4r_{1}/a} + \left(\frac{a}{4}\right)^{2} e^{-4r_{1}/a} \right]$ 

$$I_{2} = 4\pi \left[ -\left(\frac{a}{4}\right)^{2} e^{-4r_{1}/a} - \frac{2}{r_{1}} \left(\frac{a}{4}\right)^{3} e^{-4r_{1}/a} + \frac{2}{r_{1}} \left(\frac{a}{4}\right)^{3} \right]$$

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$$I_{2} = 4\pi \left[ \frac{1}{r_{1}} \int_{0}^{r_{1}} r_{2}^{2} e^{-4r_{2}/a} dr_{2} + \int_{r_{1}}^{\infty} r_{2} e^{-4r_{2}/a} dr_{2} \right]$$
  
=  $4\pi \left[ \frac{1}{r_{1}} \left( -\left(\frac{a}{4}\right) r_{1}^{2} e^{-4r_{1}/a} - 2\left(\frac{a}{4}\right)^{2} r_{1} e^{-4r_{1}/a} - 2\left(\frac{a}{4}\right)^{3} e^{-4r_{1}/a} + 2\left(\frac{a}{4}\right)^{3} \right) + \left(\frac{a}{4}\right) r_{1} e^{-4r_{1}/a} + \left(\frac{a}{4}\right)^{2} e^{-4r_{1}/a} \right]$ 

$$I_{2} = 4\pi \left[ -\left(\frac{a}{4}\right)^{2} e^{-4r_{1}/a} - \frac{2}{r_{1}} \left(\frac{a}{4}\right)^{3} e^{-4r_{1}/a} + \frac{2}{r_{1}} \left(\frac{a}{4}\right)^{3} \right]$$
$$= 4\pi \frac{2a^{3}}{64r_{1}} \left[ -\frac{r_{1}}{2} \left(\frac{4}{a}\right) e^{-4r_{1}/a} - e^{-4r_{1}/a} + 1 \right]$$

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$$I_{2} = 4\pi \left[ \frac{1}{r_{1}} \int_{0}^{r_{1}} r_{2}^{2} e^{-4r_{2}/a} dr_{2} + \int_{r_{1}}^{\infty} r_{2} e^{-4r_{2}/a} dr_{2} \right]$$
  
=  $4\pi \left[ \frac{1}{r_{1}} \left( -\left(\frac{a}{4}\right) r_{1}^{2} e^{-4r_{1}/a} - 2\left(\frac{a}{4}\right)^{2} r_{1} e^{-4r_{1}/a} - 2\left(\frac{a}{4}\right)^{3} e^{-4r_{1}/a} + 2\left(\frac{a}{4}\right)^{3} \right) + \left(\frac{a}{4}\right) r_{1} e^{-4r_{1}/a} + \left(\frac{a}{4}\right)^{2} e^{-4r_{1}/a} \right]$ 

$$\begin{split} l_2 &= 4\pi \left[ -\left(\frac{a}{4}\right)^2 e^{-4r_1/a} - \frac{2}{r_1} \left(\frac{a}{4}\right)^3 e^{-4r_1/a} + \frac{2}{r_1} \left(\frac{a}{4}\right)^3 \right] \\ &= 4\pi \frac{2a^3}{64r_1} \left[ -\frac{r_1}{2} \left(\frac{4}{a}\right) e^{-4r_1/a} - e^{-4r_1/a} + 1 \right] = \frac{\pi a^3}{8r_1} \left[ 1 - \left(1 + \frac{2r_1}{a}\right) e^{-4r_1/a} \right] \end{split}$$

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The expression for  $\langle V_{ee} \rangle$  becomes



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the variational principle, with no adjustable parameters, gives an upper bound for the ground state energy

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the variational principle, with no adjustable parameters, gives an upper bound for the ground state energy of  $E_0 = -109 + 34 = -75$  eV

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• Helium atom review

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- Helium atom review
- Improved trial wavefunction



- Helium atom review
- Improved trial wavefunction
- A much better energy!

### The helium atom so far...



The helium atom is an excellent "real world" example of the use of the variational method to compute the ground state energy of an analytically insoluble quantum system


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$$H = -\frac{\hbar^2}{2m} (\nabla_1^2 + \nabla_2^2) - \frac{e^2}{4\pi\epsilon_0} \left(\frac{2}{r_1} + \frac{2}{r_2} - \frac{1}{|\vec{r_1} - \vec{r_2}|}\right)$$



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the simplest model is to ignore the electronelectron interaction with both electrons in the hydrogenic ground state



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with the variational parameter being the effective charge Z which takes into account screening of the nucleus by the "other" electron



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Starting with the full Hamiltonian

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The helium atom



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$$\begin{aligned} H &= -\frac{\hbar^2}{2m} (\nabla_1^2 + \nabla_2^2) - \frac{e^2}{4\pi\epsilon_0} \left(\frac{2}{r_1} + \frac{2}{r_2}\right) + \frac{e^2}{4\pi\epsilon_0} \frac{1}{|\vec{r_1} - \vec{r_2}|} + \frac{e^2}{4\pi\epsilon_0} \left(\frac{Z}{r_1} - \frac{Z}{r_1} + \frac{Z}{r_2} - \frac{Z}{r_2}\right) \\ &= -\frac{\hbar^2}{2m} (\nabla_1^2 + \nabla_2^2) - \frac{e^2}{4\pi\epsilon_0} \left(\frac{Z}{r_1} + \frac{Z}{r_2}\right) \end{aligned}$$

this can be rearranged to be the sum of a non-interaction Hamiltonian matching the new trial function

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V

So,  $\langle H \rangle = -109 \text{ eV} + 34 \text{ eV} = -75 \text{ eV}$ , but it is possible to do better still with a more suitable trial function which has a variational parameter which will minimize the energy

$$\psi(\vec{r}_1, \vec{r}_2) = \frac{8}{\pi a^3} e^{-2r_1/a} e^{-2r_2/a} \qquad \qquad \psi_1(\vec{r}_1, \vec{r}_2) = \frac{Z^3}{\pi a^3} e^{-Zr_1/a} e^{-Zr_2/a}$$

with the variational parameter being the effective charge Z which takes into account screening of the nucleus by the "other" electron

Starting with the full Hamiltonian then adding and subtracting terms with Z replacing 2 in the Coulomb terms

$$\begin{aligned} \mathcal{H} &= -\frac{\hbar^2}{2m} (\nabla_1^2 + \nabla_2^2) - \frac{e^2}{4\pi\epsilon_0} \left(\frac{2}{r_1} + \frac{2}{r_2}\right) + \frac{e^2}{4\pi\epsilon_0} \frac{1}{|\vec{r_1} - \vec{r_2}|} + \frac{e^2}{4\pi\epsilon_0} \left(\frac{Z}{r_1} - \frac{Z}{r_1} + \frac{Z}{r_2} - \frac{Z}{r_2}\right) \\ &= -\frac{\hbar^2}{2m} (\nabla_1^2 + \nabla_2^2) - \frac{e^2}{4\pi\epsilon_0} \left(\frac{Z}{r_1} + \frac{Z}{r_2}\right) + \frac{e^2}{4\pi\epsilon_0} \left(\frac{Z-2}{r_1} + \frac{Z-2}{r_2} + \frac{1}{|\vec{r_1} - \vec{r_2}|}\right) \end{aligned}$$

this can be rearranged to be the sum of a non-interaction Hamiltonian matching the new trial function and an effective "interaction" term

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$$\langle H \rangle = 2Z^2 E_1 + 2(Z-2) \left(\frac{e^2}{4\pi\epsilon_0}\right) \left\langle \frac{1}{r} \right\rangle$$



Taking the expectation value of the Hamiltonian, we have:

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the final result is thus

$$\langle H \rangle = \left[ 2Z^2 - 4Z(Z-2) - \frac{5}{4}Z \right] E_1 = \left[ -2Z^2 + \frac{27}{4}Z \right] E_1$$

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# Improved helium energy

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This energy can be minimized with respect to the effective charge, Z

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V

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$$\langle H \rangle = \left[ 2 \left( \frac{27}{16} \right)^2 - 4 \left( \frac{27}{16} \right) \left( \frac{27}{16} - 2 \right) - \frac{6}{4} \frac{27}{16} \right] E_1$$



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$$= \left[ -2 \left( \frac{27}{16} \right)^2 + \frac{27}{4} \frac{27}{16} \right] E_1 = 2 \left( \frac{27}{16} \right)^2 E_1$$



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The improved energy estimate becomes

$$\begin{aligned} H \rangle &= \left[ 2 \left( \frac{27}{16} \right)^2 - 4 \left( \frac{27}{16} \right) \left( \frac{27}{16} - 2 \right) - \frac{6}{4} \frac{27}{16} \right] E_1 \\ &= \left[ -2 \left( \frac{27}{16} \right)^2 + \frac{27}{4} \frac{27}{16} \right] E_1 = 2 \left( \frac{27}{16} \right)^2 E_1 = \frac{1}{2} \left( \frac{3}{2} \right)^6 E_1 = -77.5 \, \text{eV} \end{aligned}$$

comparing the different solutions with the experimental value



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comparing the different solutions with the experimental value

experimental helium energy  $\longrightarrow$  -78.975 eV

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comparing the different solutions with the experimental value

 $\begin{array}{rcl} \mbox{experimental helium energy} & \longrightarrow & -78.975 \mbox{ eV} \\ \mbox{ignoring e-e interaction} & \longrightarrow & -109 \mbox{ eV} & \sim 38\% \mbox{ error} \end{array}$ 



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The helium atom